

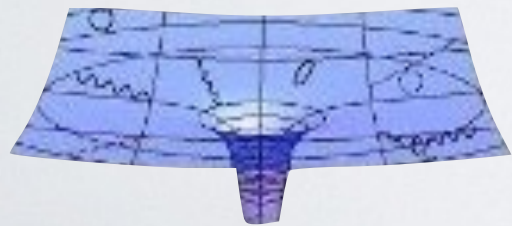
PHYSICAL PROPERTIES OF PERTURBED BLACK-HOLE INITIAL DATA

(ON HYPERBOLOIDAL SLICES)

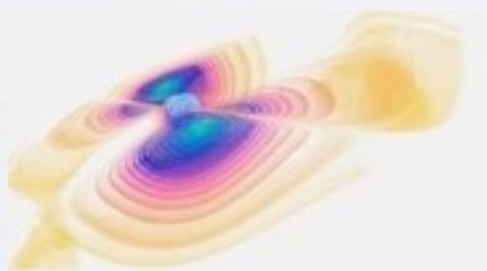
David Schinkel
Rodrigo P. Macedo
Marcus Ansorg

Class. Quantum Grav. 31 075017 (2014)

Class. Quantum Grav. 31 165001 (2014)



GRK 1523
Quantum and
Gravitational Fields



SFB/Transregio 7



Friedrich-Schiller-Universität Jena



OUTLINE

1. INITIAL DATA
2. (ON HYPERBOLOIDAL SLICES)
3. PERTURBED BLACK-HOLE
4. PHYSICAL PROPERTIES OF

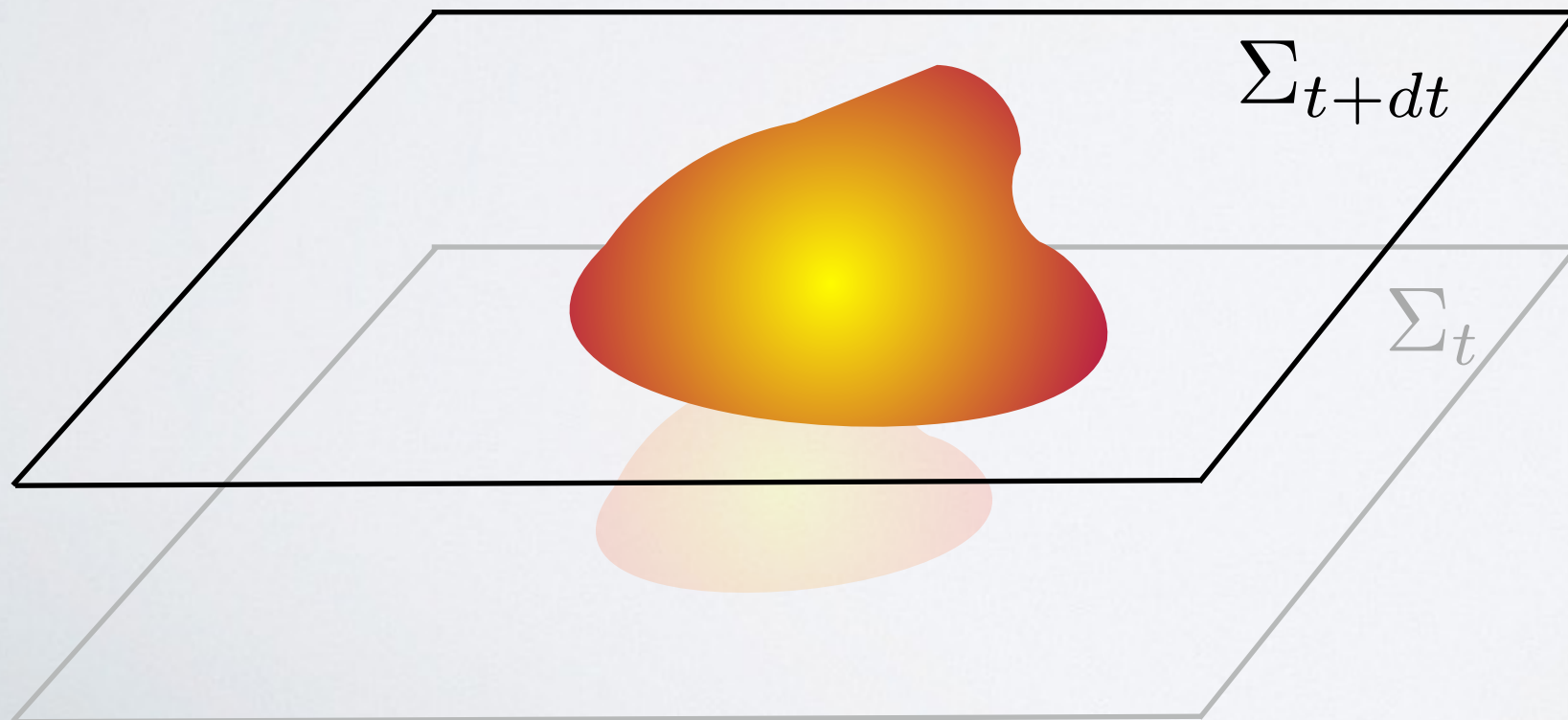
INITIAL-DATA PROBLEM

3+1 Decomposition : $\{x^0; x^i\} \quad i = 1 \dots 3$



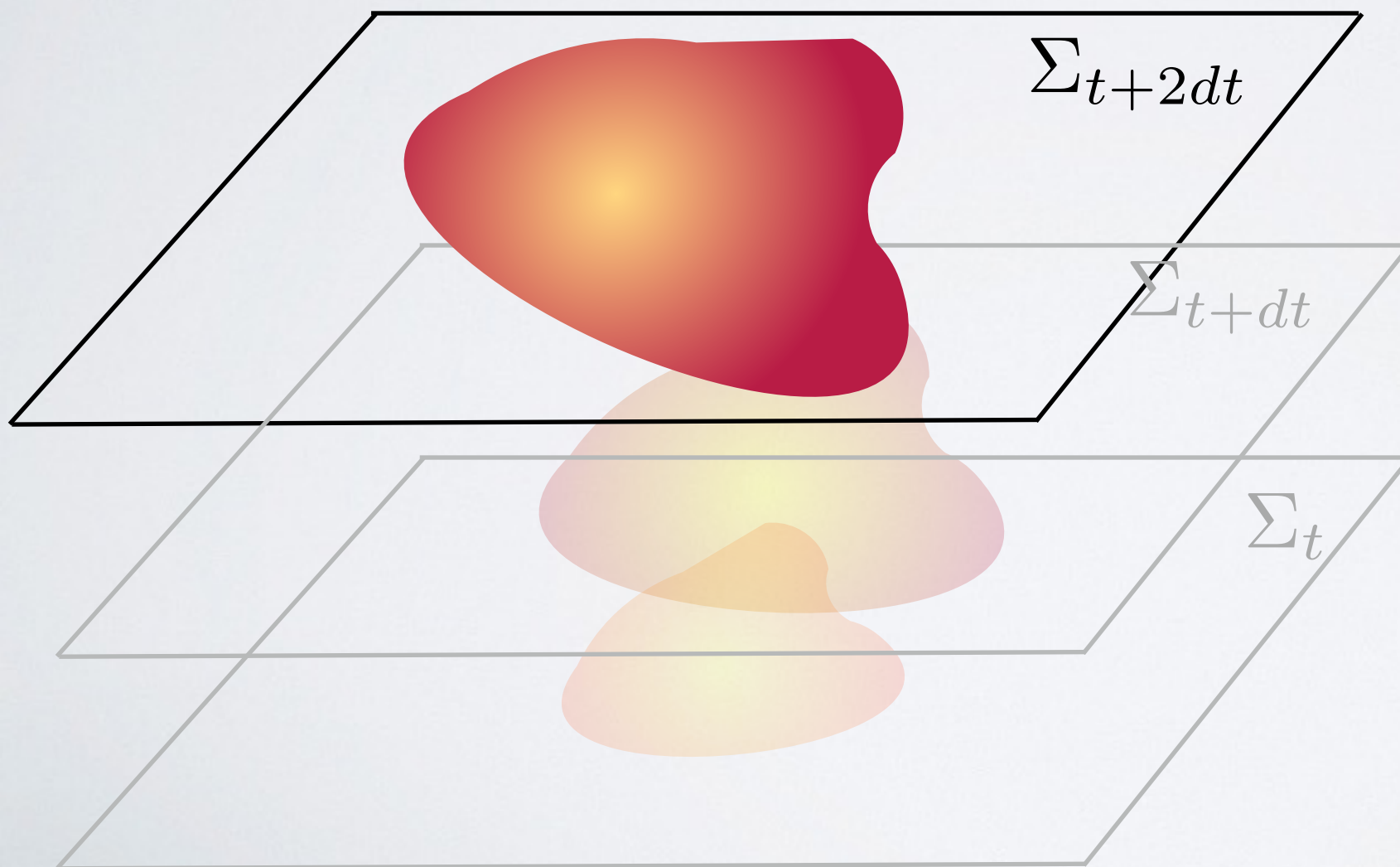
INITIAL-DATA PROBLEM

3+1 Decomposition : $\{x^0; x^i\} \quad i = 1 \dots 3$

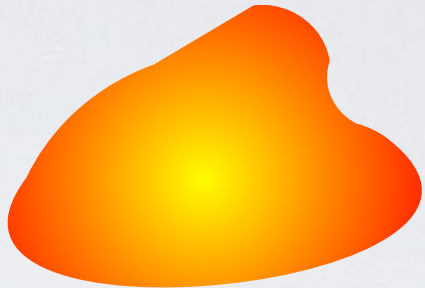


INITIAL-DATA PROBLEM

3+1 Decomposition : $\{x^0; x^i\} \quad i = 1 \dots 3$



INITIAL-DATA PROBLEM

How do we obtain  ?

What are the geometrical quantities describing it?

How do we give a physical meaning to it?



INITIAL-DATA PROBLEM

geometrical quantities

- Spatial metric at Σ_t : γ_{ij}

$$ds^2 = g_{00}(dx^0)^2 + 2g_{0i}dx^0dx^i + \gamma_{ij}dx^i dx^j$$



INITIAL-DATA PROBLEM

geometrical quantities

- Spatial metric at Σ_t : γ_{ij}

$$ds^2|_{\Sigma_t} = \underbrace{g_{00}(dx^0)^2 + 2g_{0i}dx^0dx^i}_{\text{Gauge Freedom}} + \gamma_{ij}dx^i dx^j$$

Gauge Freedom



INITIAL-DATA PROBLEM

geometrical quantities

- Spatial metric at Σ_t : γ_{ij}
- Time derivative of γ_{ij} : $\left. \frac{\partial \gamma_{ij}}{\partial x^0} \right|_{\Sigma_t}$



INITIAL-DATA PROBLEM

geometrical quantities

- Spatial metric at Σ_t : γ_{ij}
- Time derivative of γ_{ij} : ~~$\frac{\partial \gamma_{ij}}{\partial x^0} \Big|_{\Sigma_t}$~~ **NOT A TENSOR**

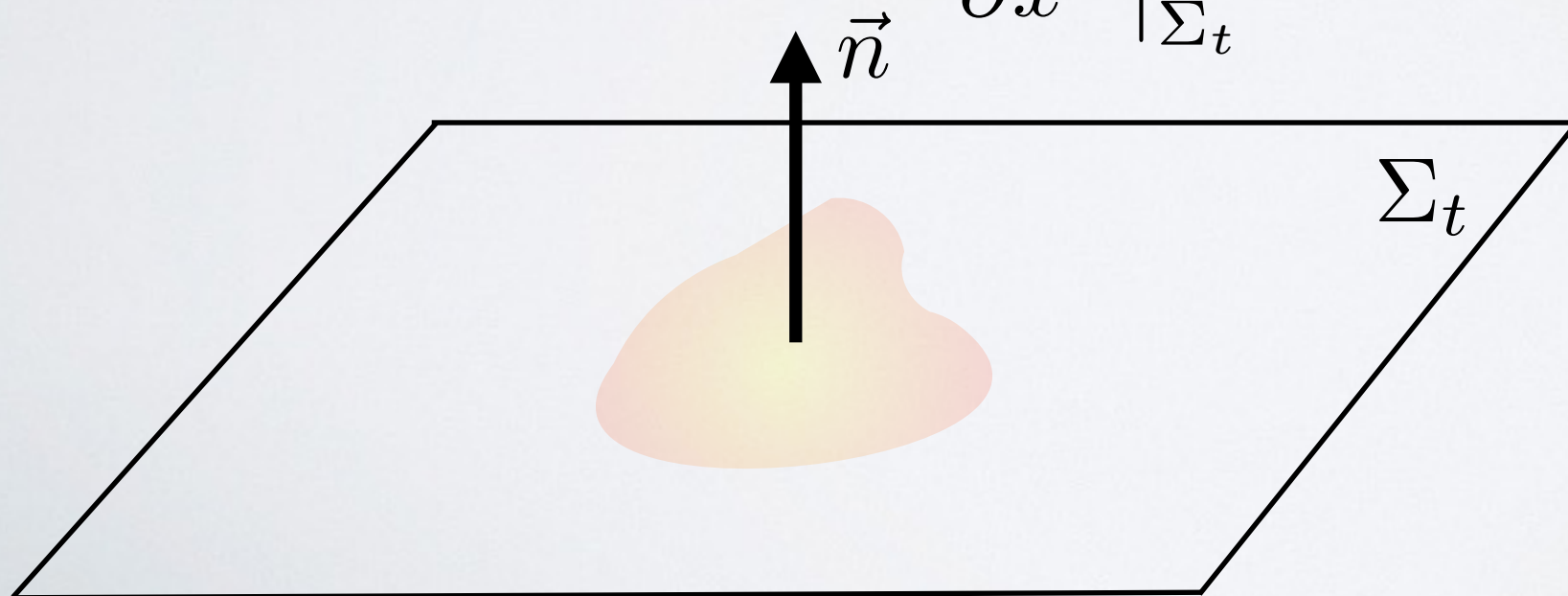


INITIAL-DATA PROBLEM

geometrical quantities

- Spatial metric at Σ_t : γ_{ij}
- Time derivative of γ_{ij} : ~~$\frac{\partial \gamma_{ij}}{\partial x^0} \Big|_{\Sigma_t}$~~ **NOT A TENSOR**
- Extrinsic Curvature : $K_{ij} = \mathcal{L}_{\vec{n}} \gamma_{ij}$

$$K_{ij} \propto \frac{\partial \gamma_{ij}}{\partial x^0} \Big|_{\Sigma_t} + (\text{gauge freedom})_{ij}$$



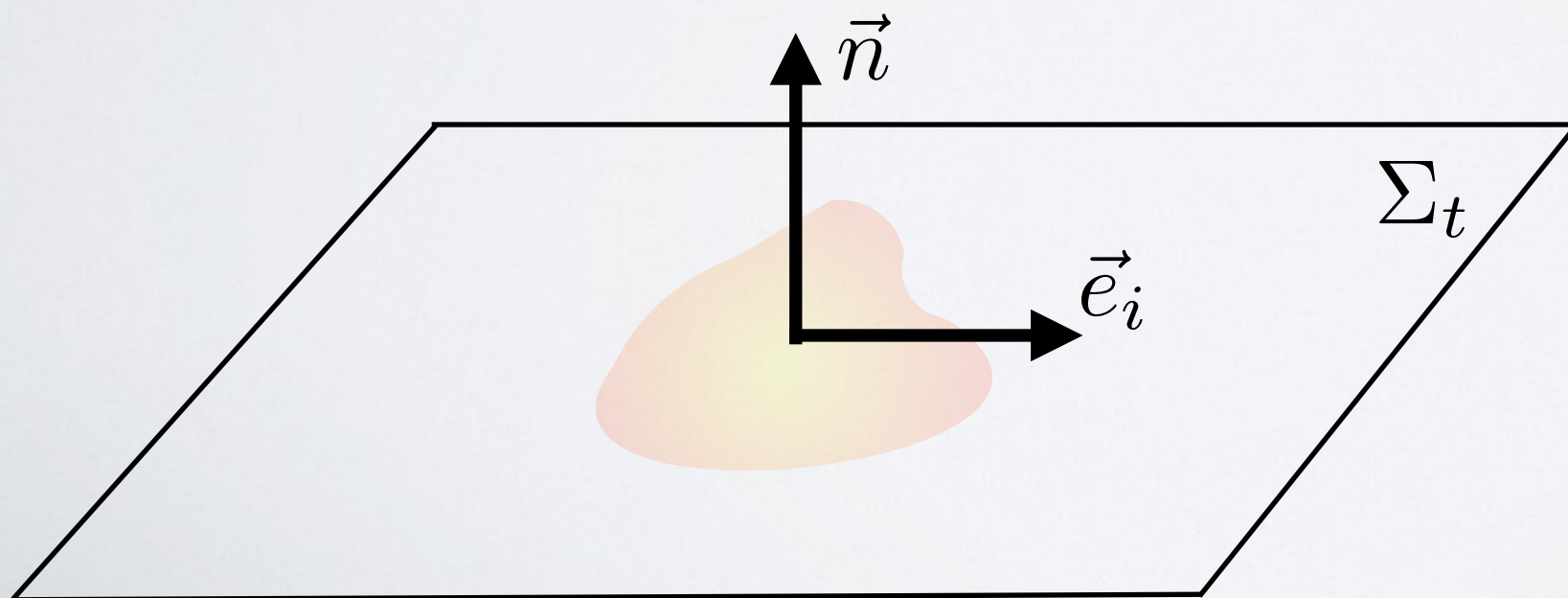
INITIAL-DATA PROBLEM

Physical meaning

- Variables: γ_{ij}, K_{ij}
- (Vacuum) Einstein's equations : Constraint Equations

$$\mathbb{H} := n^a n^b G_{ab} = {}^{(3)}R + K^2 - K_{ij} K^{ij} = 0 \quad \text{(Hamiltonian Constraint)}$$

$$\mathbb{P}_i := n^a e_i^b G_{ab} = {}^{(3)}\nabla^j (K_{ij} - \gamma_{ij} K) = 0 \quad \text{(Momentum Constraint)}$$



INITIAL-DATA PROBLEM

(Vacuum) Maxwell Equations:

- Variables: γ_{ij}, K_{ij} Dynamical Equations:

- (Vacuum) Einstein's equations: Constraint Equations

$$\frac{\partial \vec{B}}{\partial t} + \nabla \times \vec{E} = 0 \quad \frac{\partial \vec{E}}{\partial t} - \nabla \times \vec{B} = 0$$

$$\mathbb{H} := n^a n^b G_{ab} = {}^{(3)}R + K^2 - K_{ij} K^{ij} = 0 \quad \text{(Hamiltonian Constraint)}$$

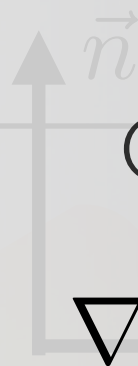
$$\mathbb{P}_i := n^a e_i^b G_{ab} = {}^{(3)}\nabla^j (K_{ij} - \gamma_{ij} K) = 0 \quad \text{(Momentum Constraint)}$$

Initial Data:

$$\vec{E}(0, \vec{x}), \quad \vec{B}(0, \vec{x})$$

Constraint Equations:

$$\nabla \cdot \vec{E} = 0 \quad \nabla \cdot \vec{B} = 0$$



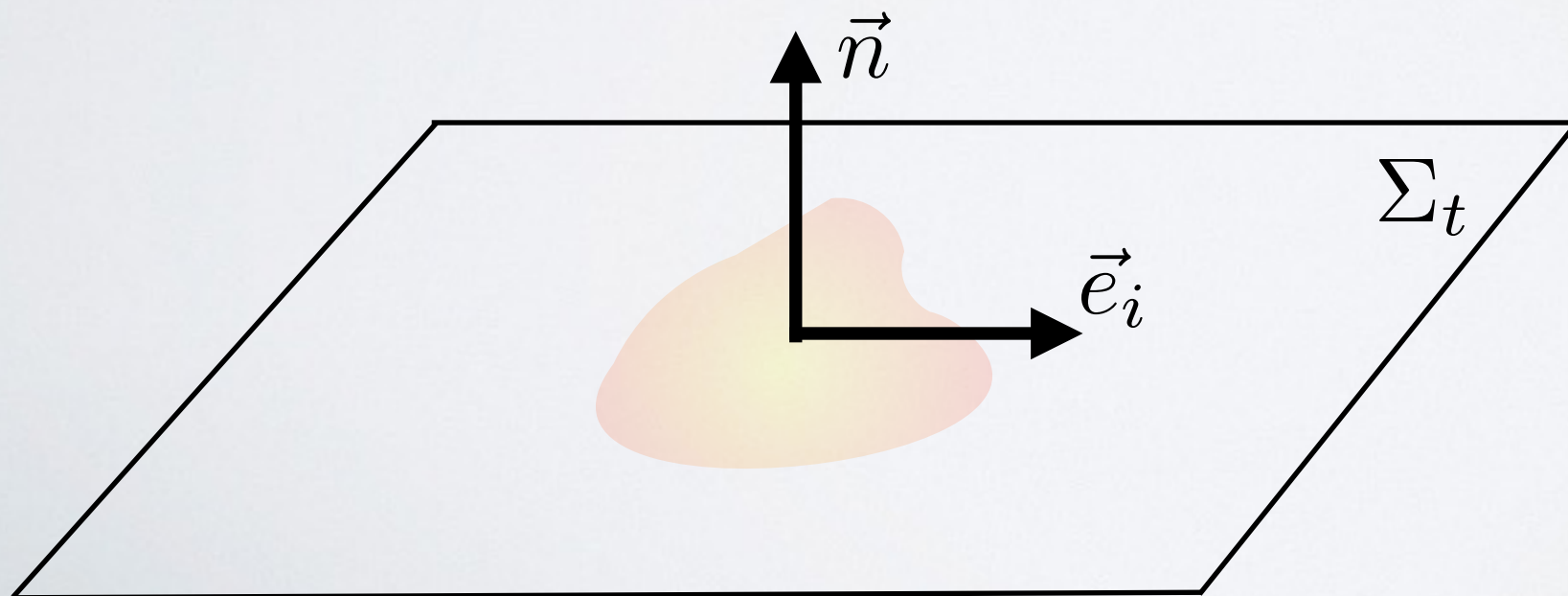
INITIAL-DATA PROBLEM

Physical meaning

- Variables: γ_{ij}, K_{ij}
- (Vacuum) Einstein's equations : Constraint Equations

$$\mathbb{H} := n^a n^b G_{ab} = {}^{(3)}R + K^2 - K_{ij} K^{ij} = 0 \quad \text{(Hamiltonian Constraint)}$$

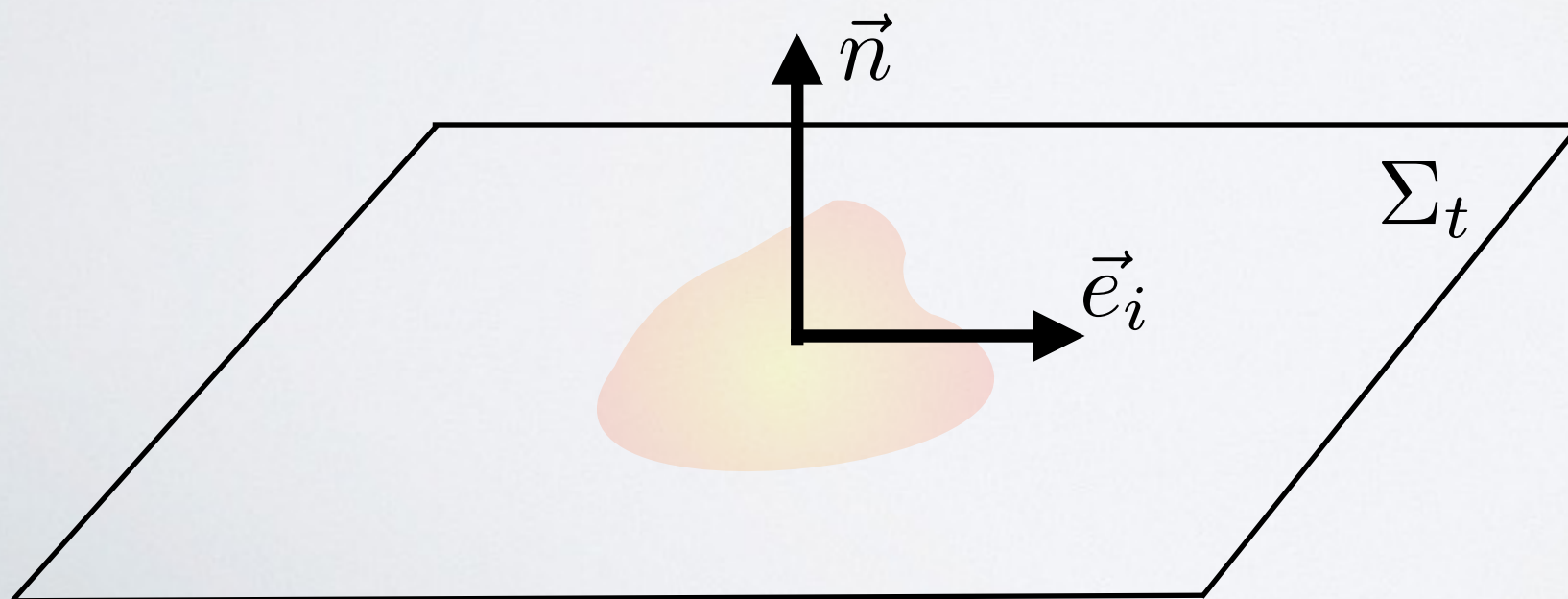
$$\mathbb{P}_i := n^a e_i^b G_{ab} = {}^{(3)}\nabla^j (K_{ij} - \gamma_{ij} K) = 0 \quad \text{(Momentum Constraint)}$$



INITIAL-DATA PROBLEM

Physical meaning

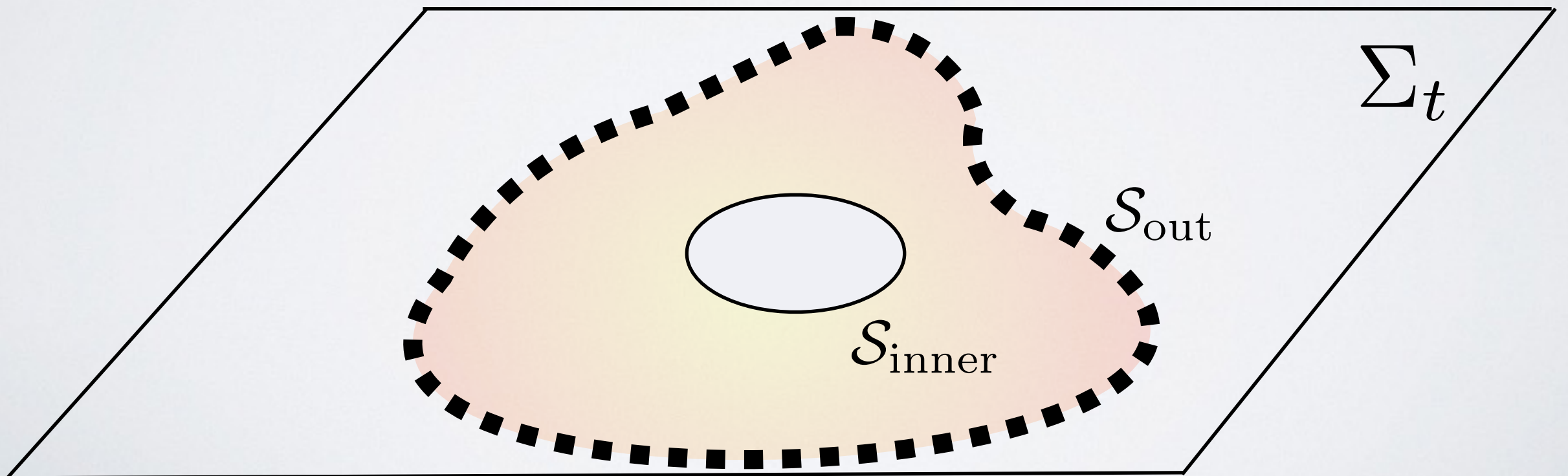
- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)
- **Elliptic Equations : Boundary value problem**



INITIAL-DATA PROBLEM

Physical meaning

- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)
- **Elliptic Equations : Boundary value problem**
 $\mathcal{S}_{\text{inner}}$: Black-hole horizon $\mathcal{S}_{\text{out}}$: infinity (“ $r \rightarrow \infty$ ”)



INITIAL-DATA PROBLEM

Physical meaning

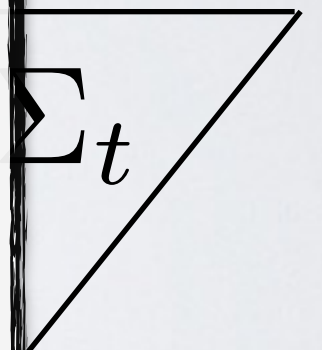
- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)

- **Elliptic Equations : Boundary value problem**

S_{inner} : Black-hole horizon : Outer Boundary : infinity (“ $r \rightarrow \infty$ ”)

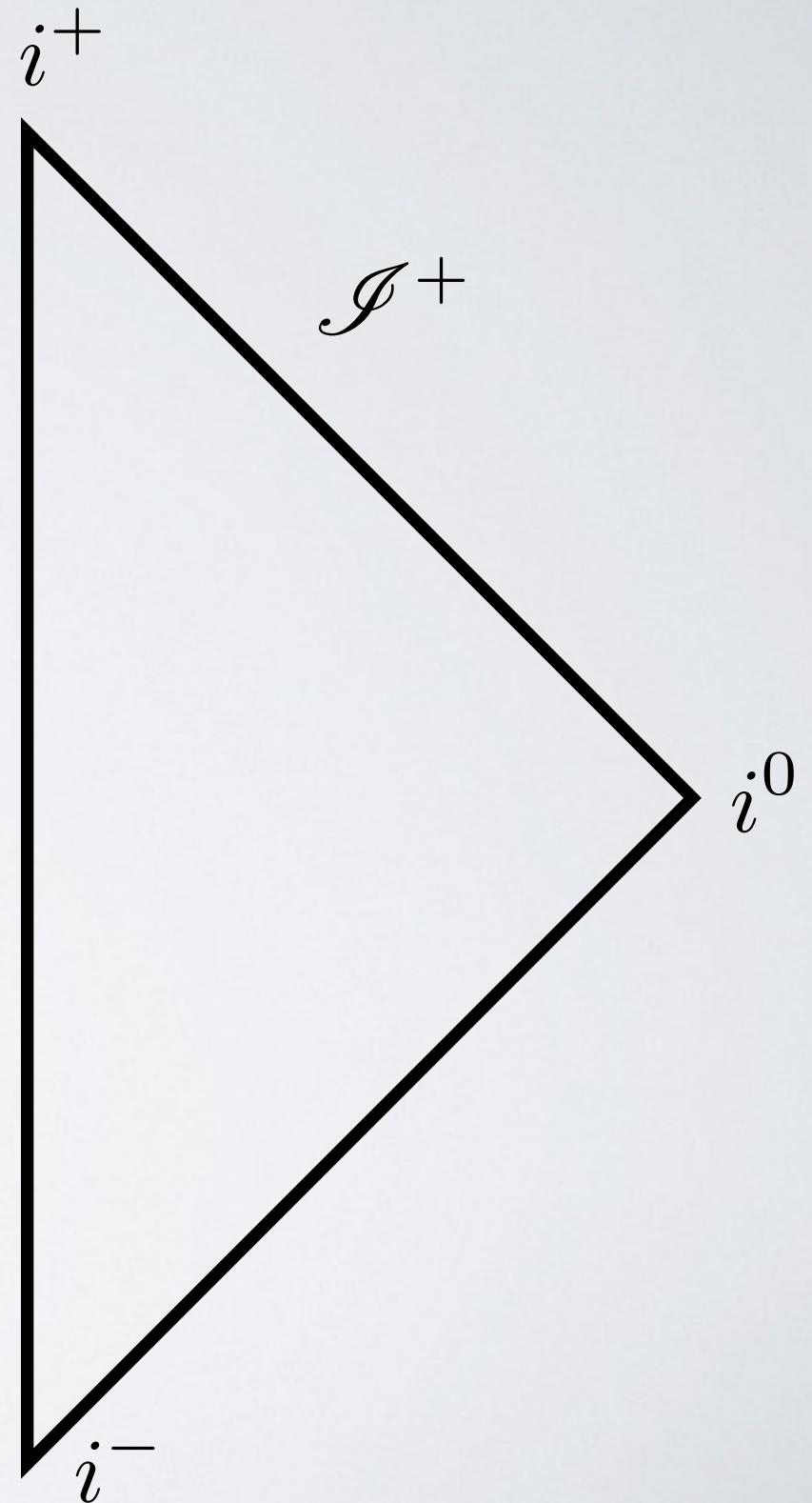
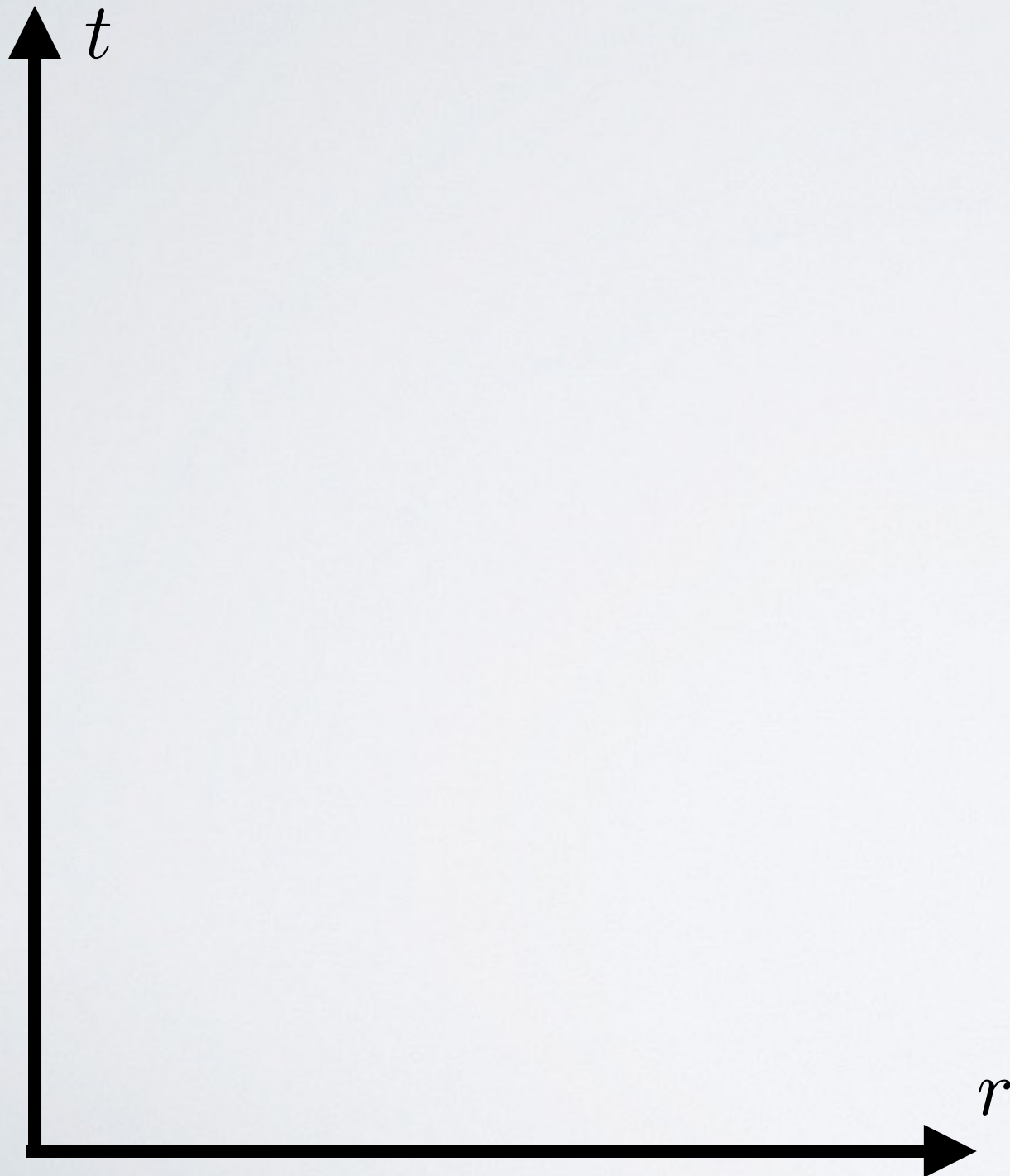
S_{out} : infinity (“ $r \rightarrow \infty$ ”)

- Which infinity ?
- How to perform the limit $r \rightarrow \infty$?



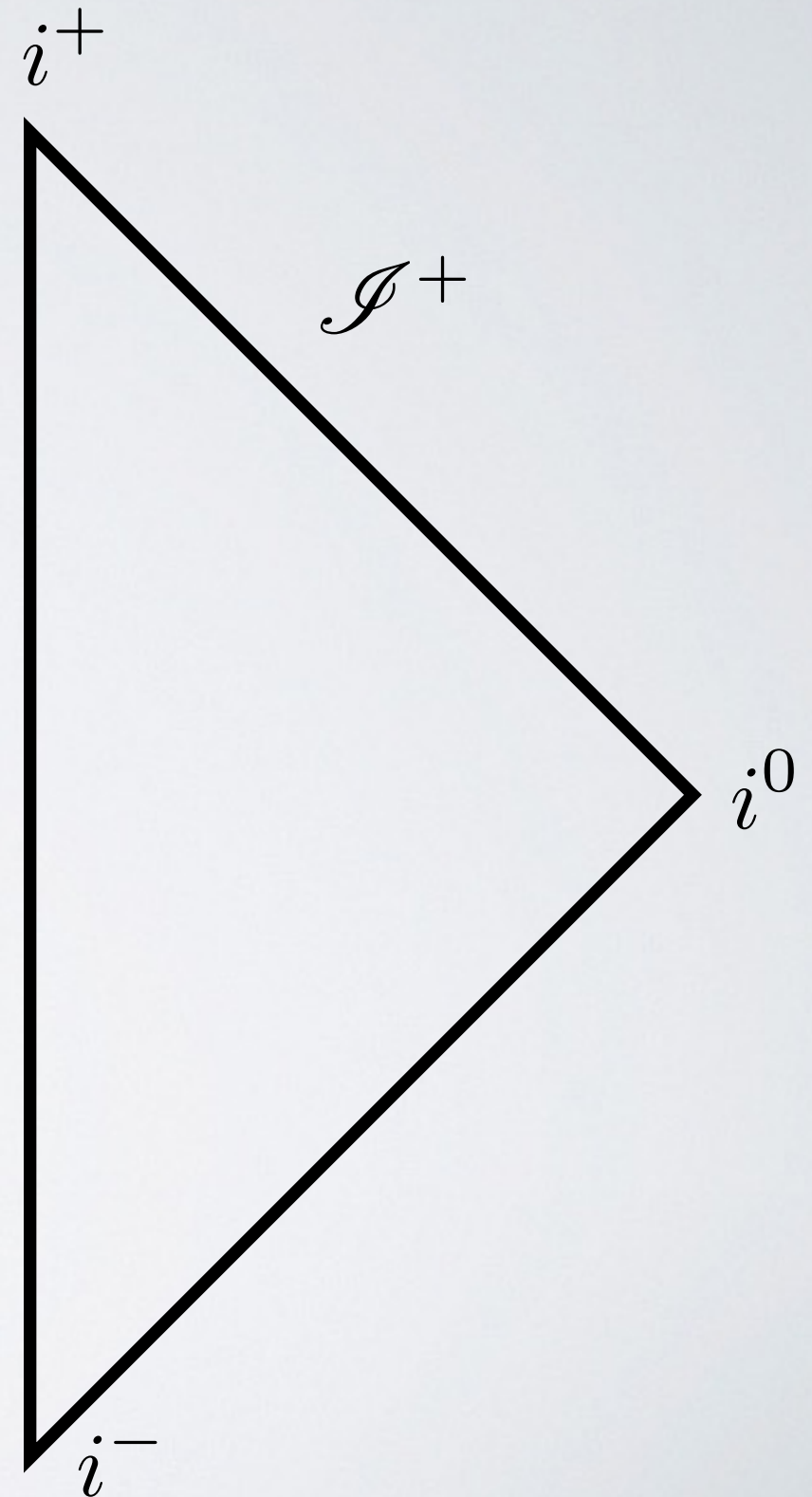
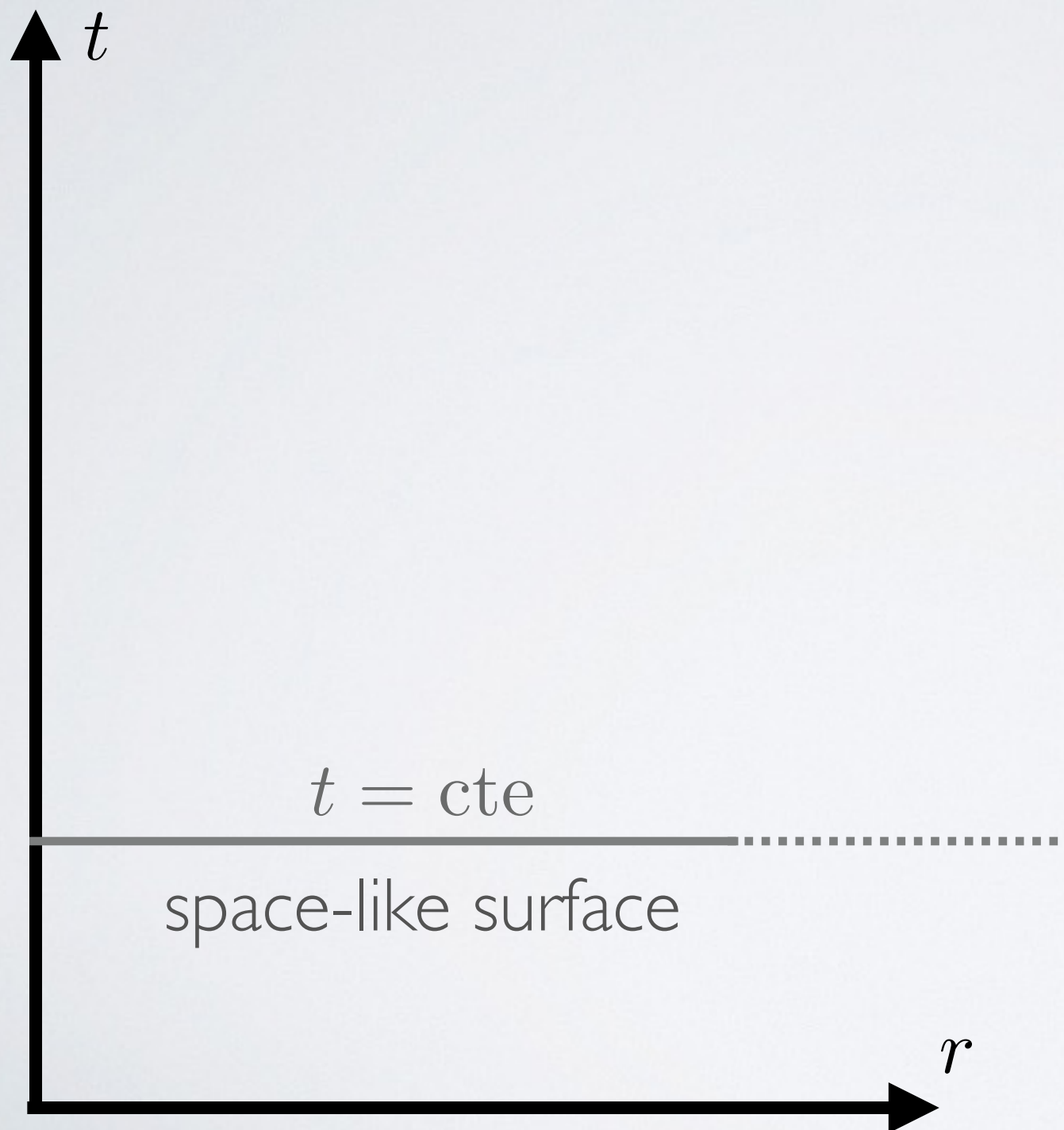
HYPERBOLOIDAL SLICES

- Minkowski Spacetime:



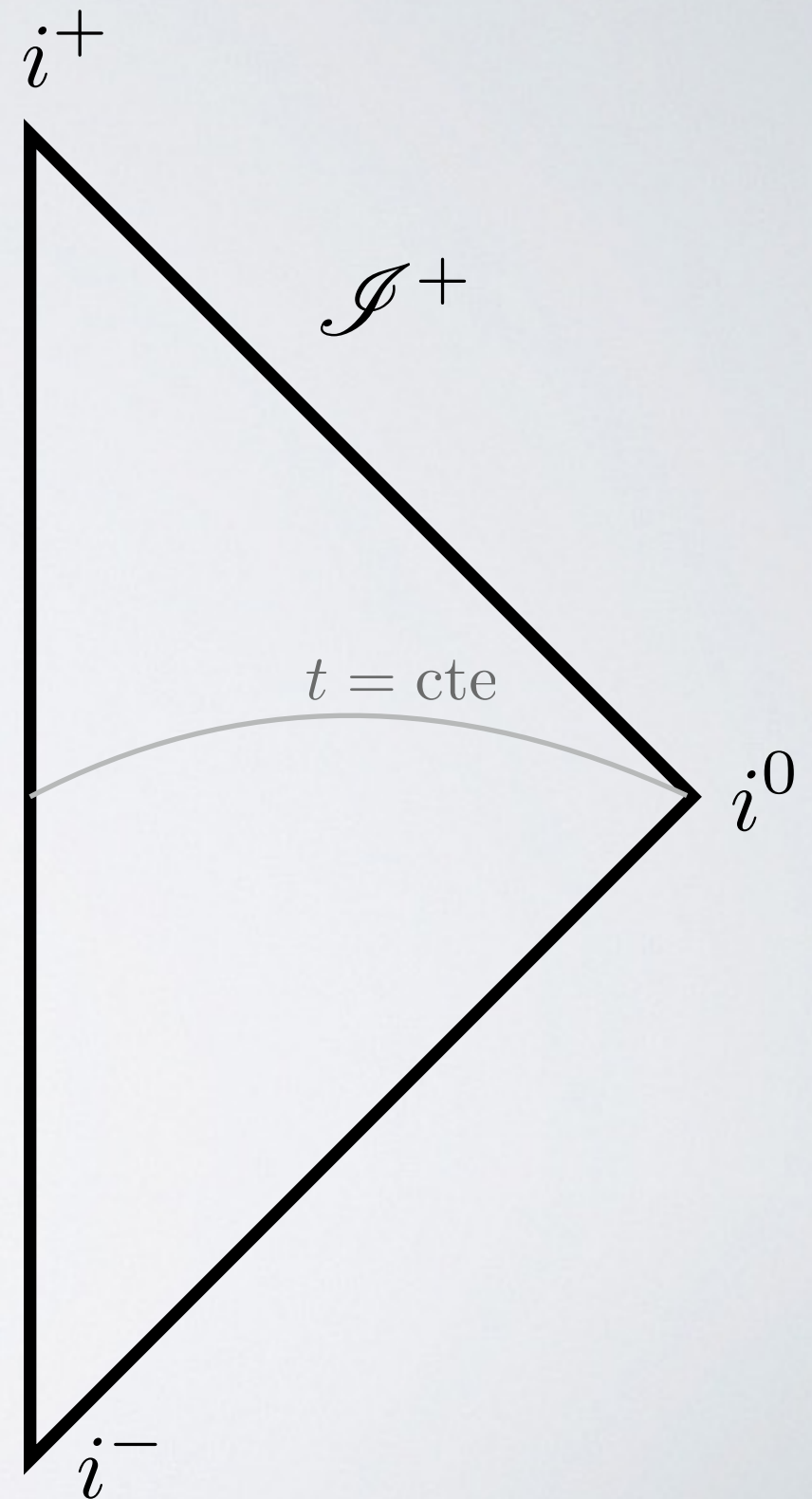
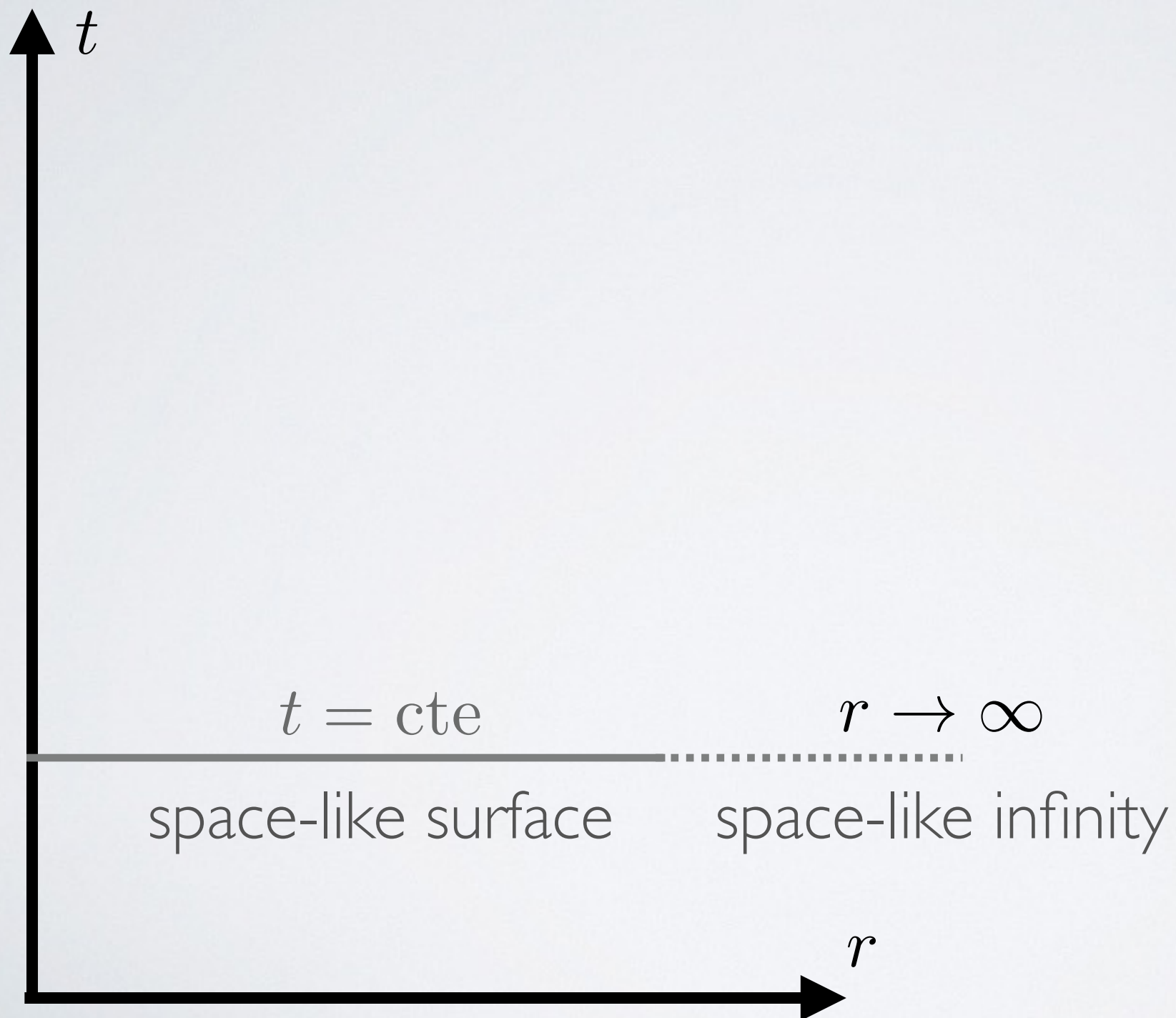
HYPERBOLOIDAL SLICES

- Minkowski Spacetime:



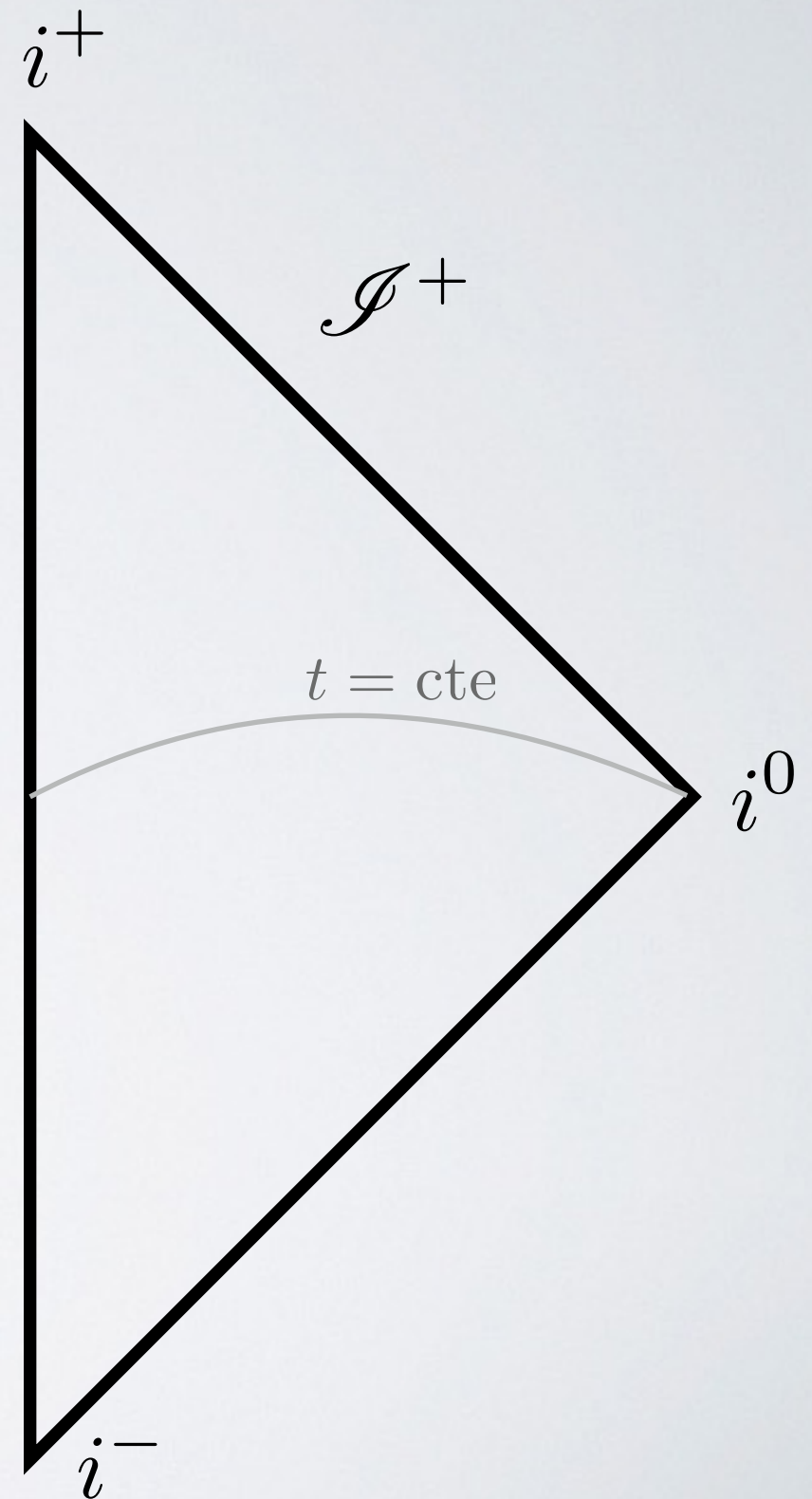
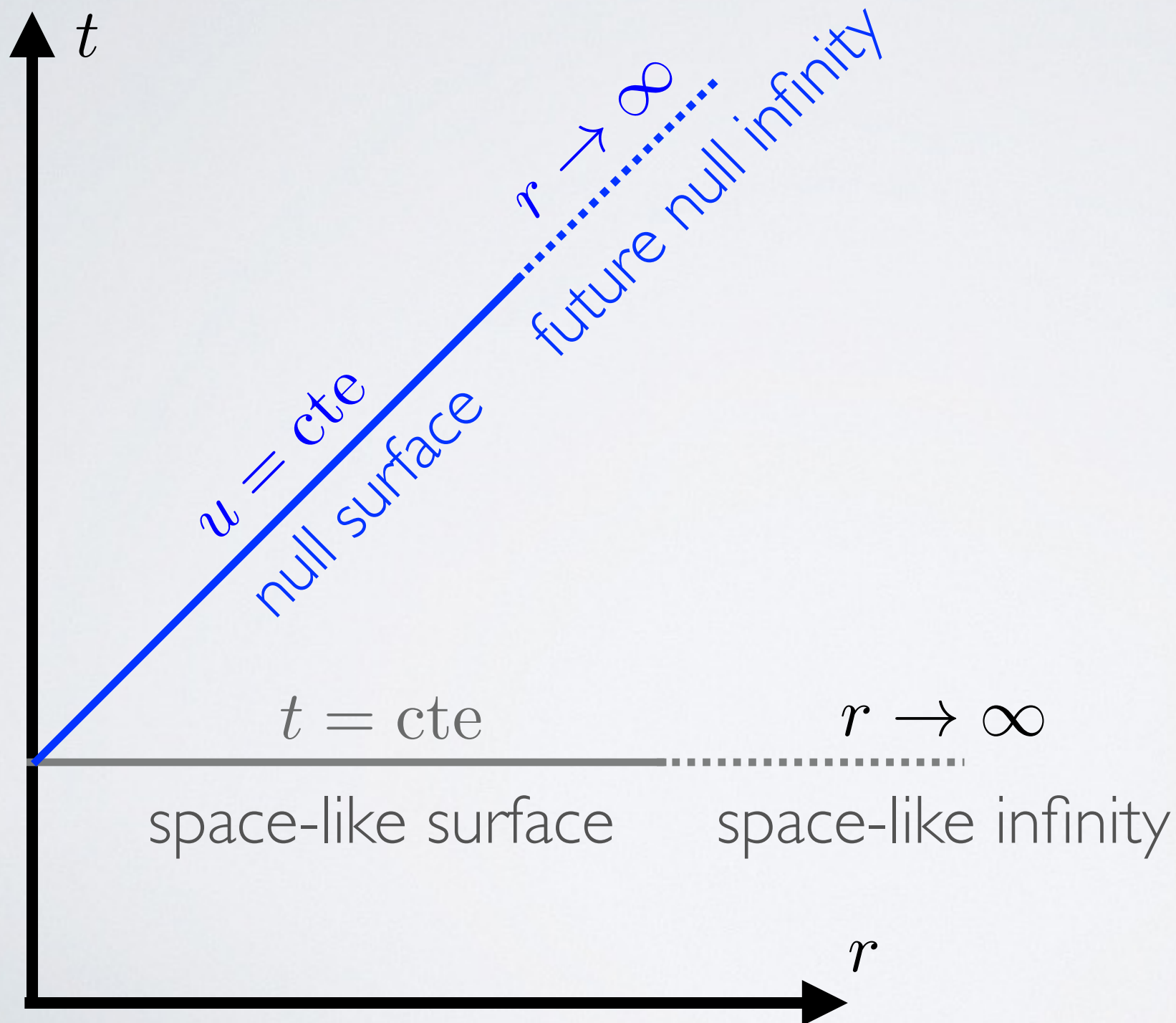
HYPERBOLOIDAL SLICES

- Minkowski Spacetime:



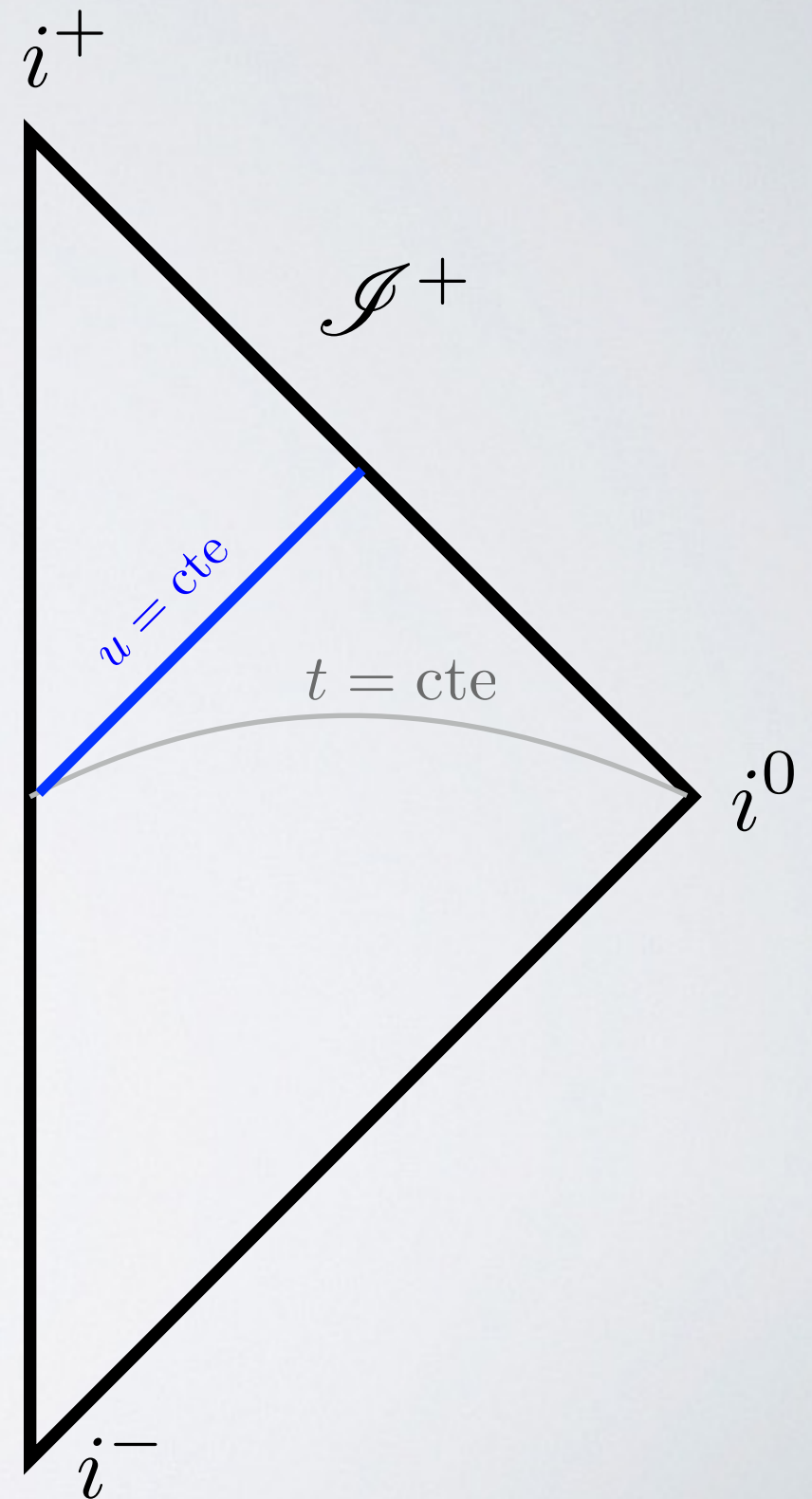
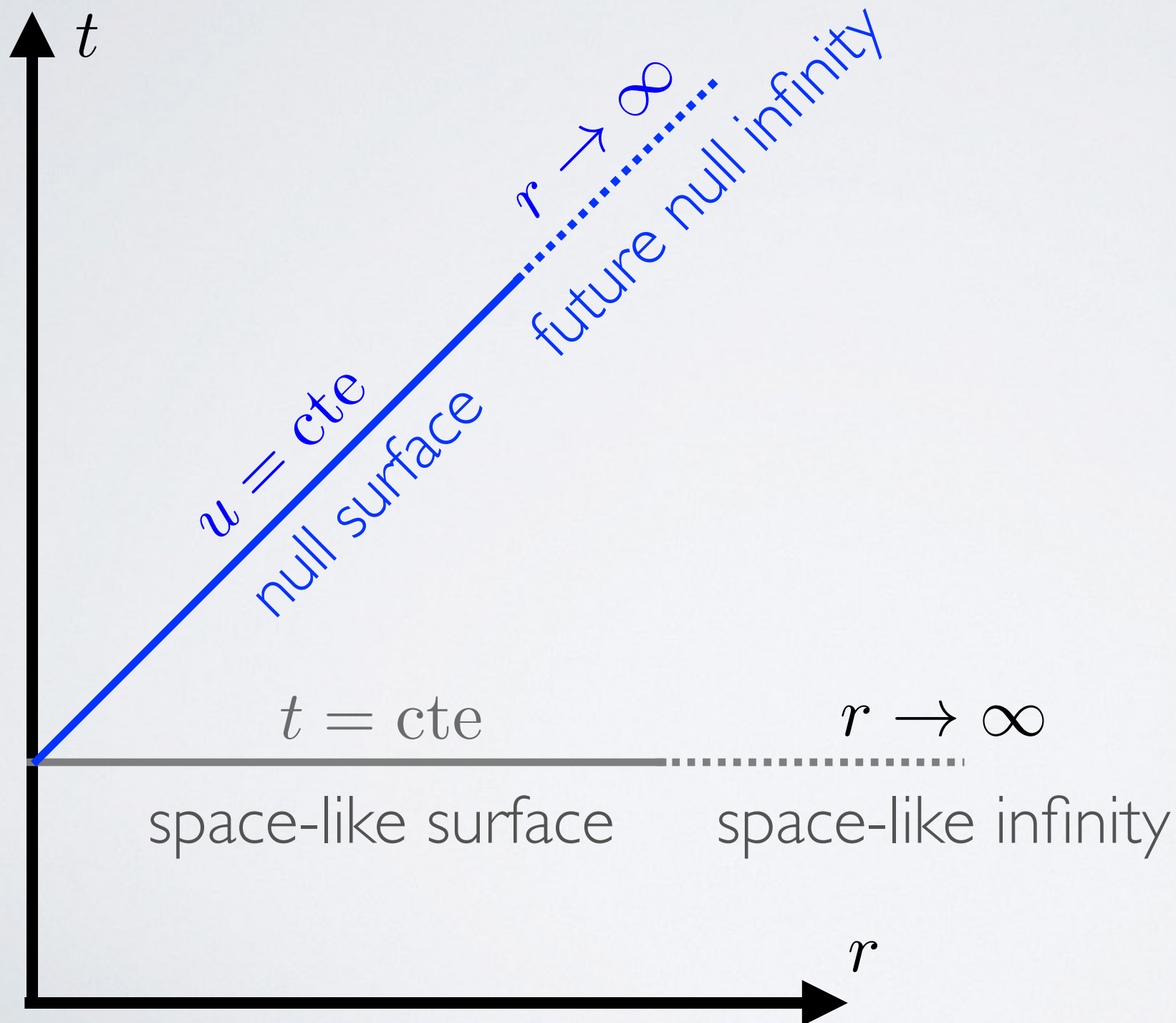
HYPERBOLOIDAL SLICES

- Minkowski Spacetime:



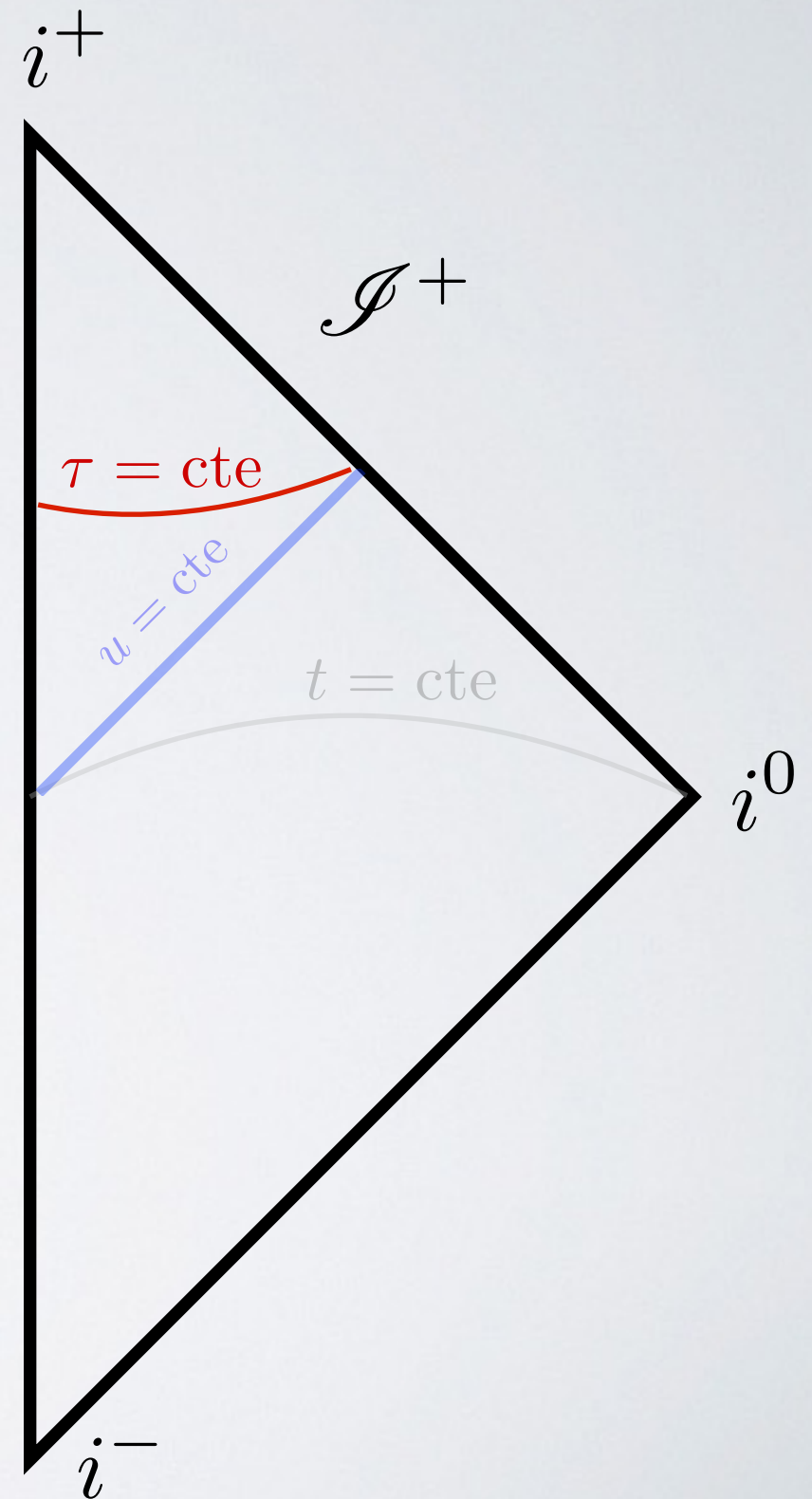
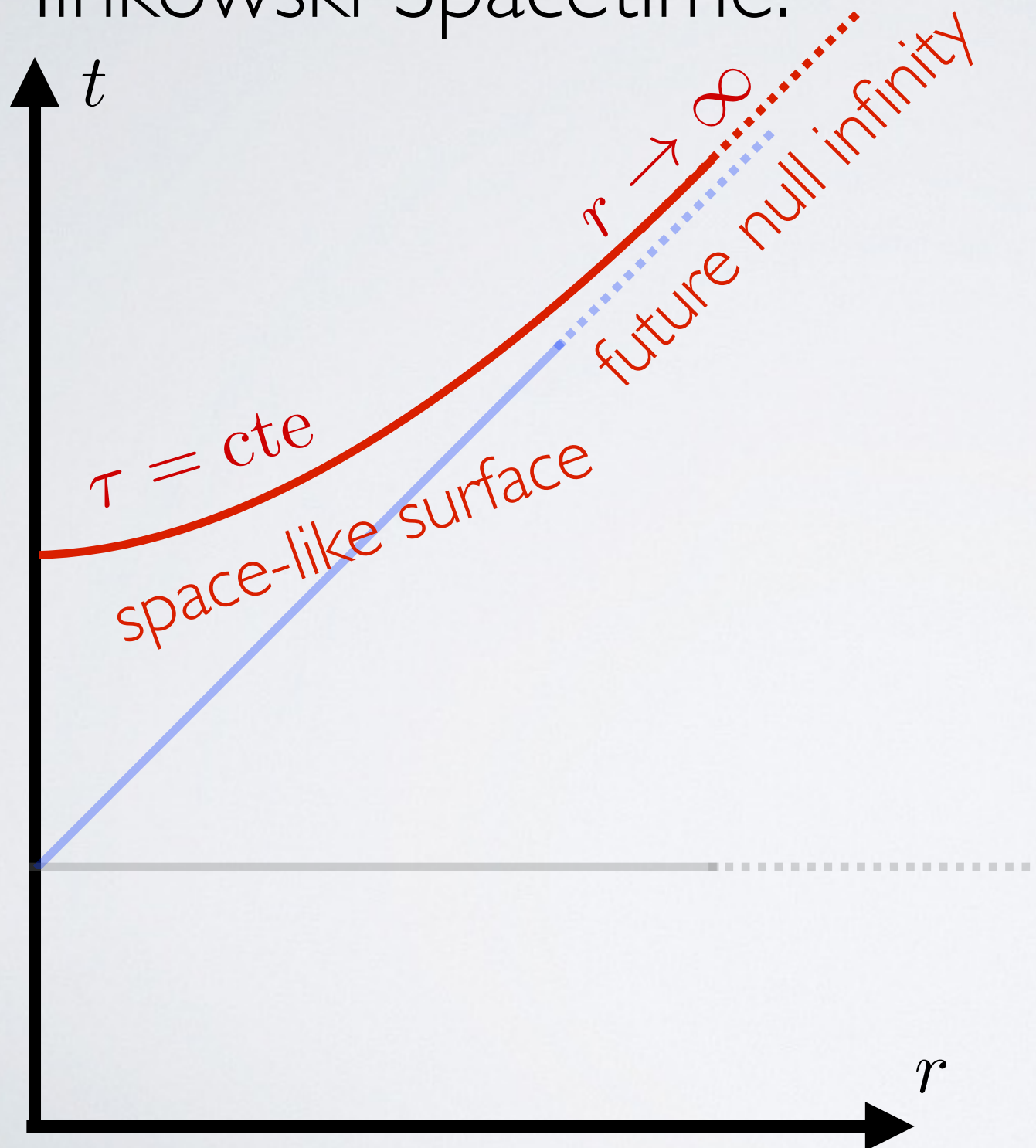
HYPERBOLOIDAL SLICES

- Minkowski Spacetime:



HYPERBOLOIDAL SLICES

- Minkowski Spacetime:



HYPERBOLOIDAL SLICES

- Schwarzschild Spacetime:

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + r^2 d\omega^2$$

$\{\tau, \sigma\}$

HYPERBOLOIDAL SLICES

- Schwarzschild Spacetime:

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + r^2 d\omega^2$$

- Introduce the new coordinates $\{\tau, \sigma\}$:

HYPERBOLOIDAL SLICES

- Schwarzschild Spacetime:

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + r^2 d\omega^2$$

- Introduce the new coordinates $\{\tau, \sigma\}$:

$$t = 2M \left\{ 2\tau + \frac{1}{\sigma} - \ln [\sigma(1 - \sigma)] \right\} \quad (\text{Hyperboloidal coordinate})$$

HYPERBOLOIDAL SLICES

- Schwarzschild Spacetime:

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + r^2 d\omega^2$$

- Introduce the new coordinates $\{\tau, \sigma\}$:

$$t = 2M \left\{ 2\tau + \frac{1}{\sigma} - \ln [\sigma(1 - \sigma)] \right\} \quad \begin{array}{l} \text{(Hyperboloidal} \\ \text{coordinate)} \end{array}$$

$$r = \frac{2M}{\sigma} \quad \text{(spatial compactification)}$$

HYPERBOLOIDAL SLICES

- Schwarzschild Spacetime:

$$ds^2 = - \left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + r^2 d\omega^2$$

- Introduce the new coordinates $\{\tau, \sigma\}$:

$$t = 2M \left\{ 2\tau + \frac{1}{\sigma} - \ln [\sigma(1 - \sigma)] \right\} \quad (\text{Hyperboloidal coordinate})$$

$$r = \frac{2M}{\sigma} \quad (\text{spatial compactification})$$

$$ds^2 = \Omega^{-2} \left[-\sigma^2(1 - \sigma)d\tau^2 + (1 + \sigma)d\sigma^2 + (1 - 2\sigma)d\tau d\sigma + \frac{1}{4}d\omega^2 \right]$$

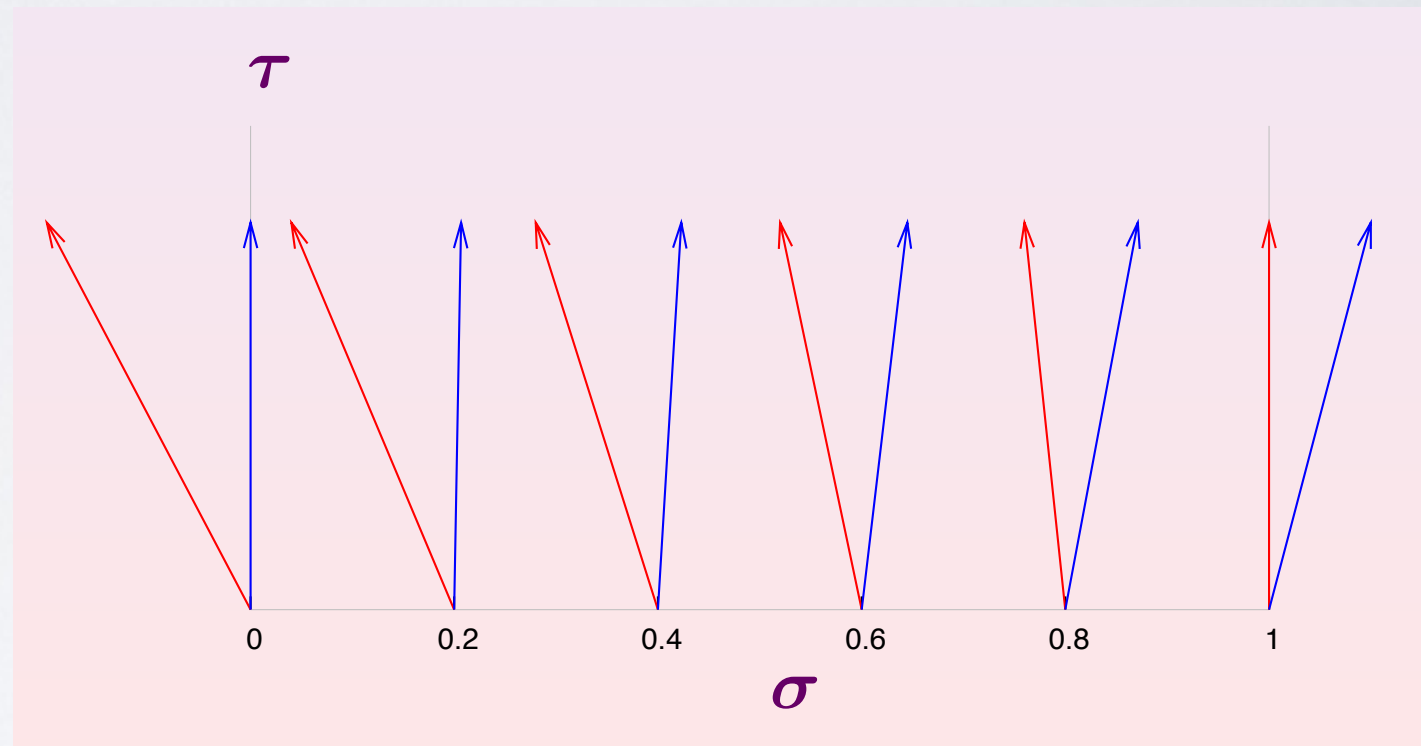
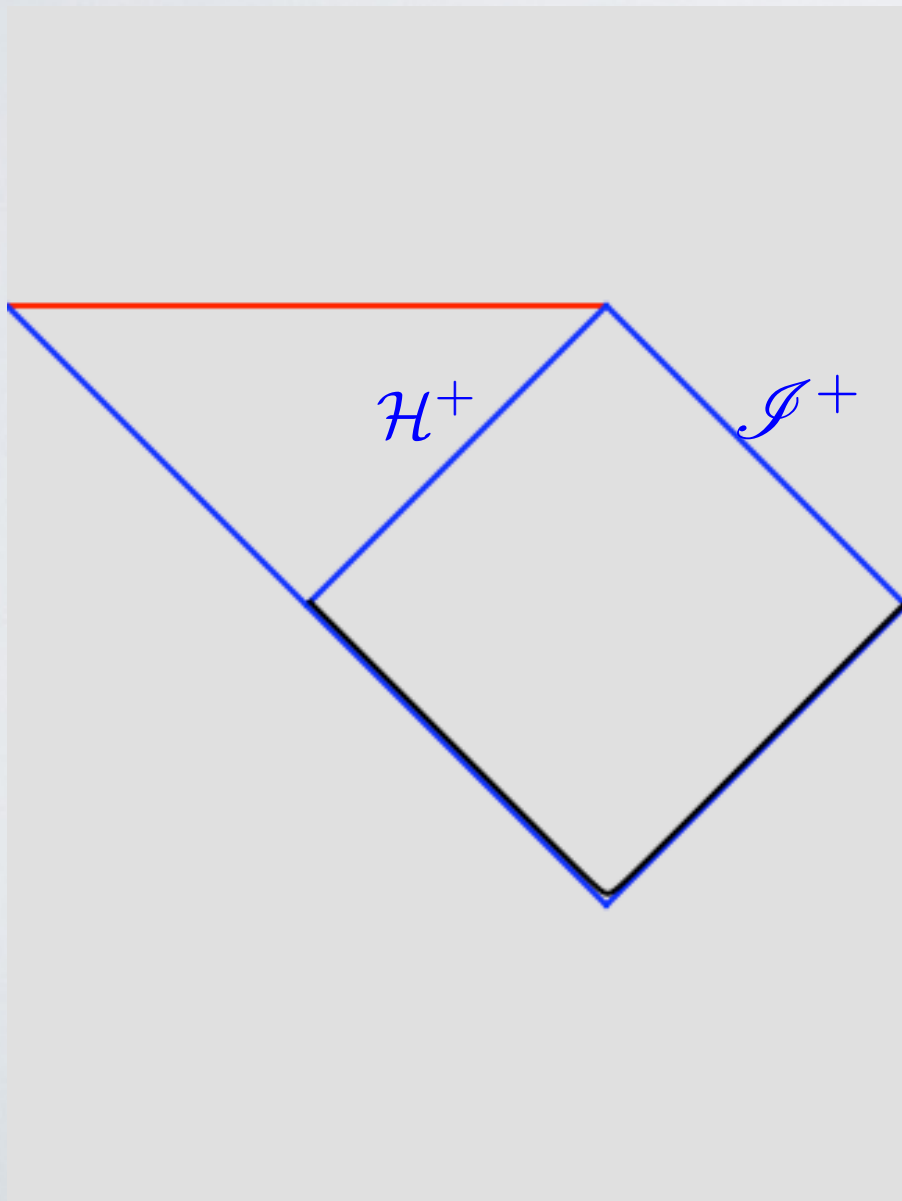
- Conformal factor: $\Omega = \frac{\sigma}{4M}$

HYPERBOLOIDAL SLICES

- Schwarzschild Spacetime:

$$\sigma = \frac{2M}{r}$$

- Location of $\mathcal{H}^+$: $\sigma = 1$
- Location of $\mathcal{I}^+$: $\sigma = 0$

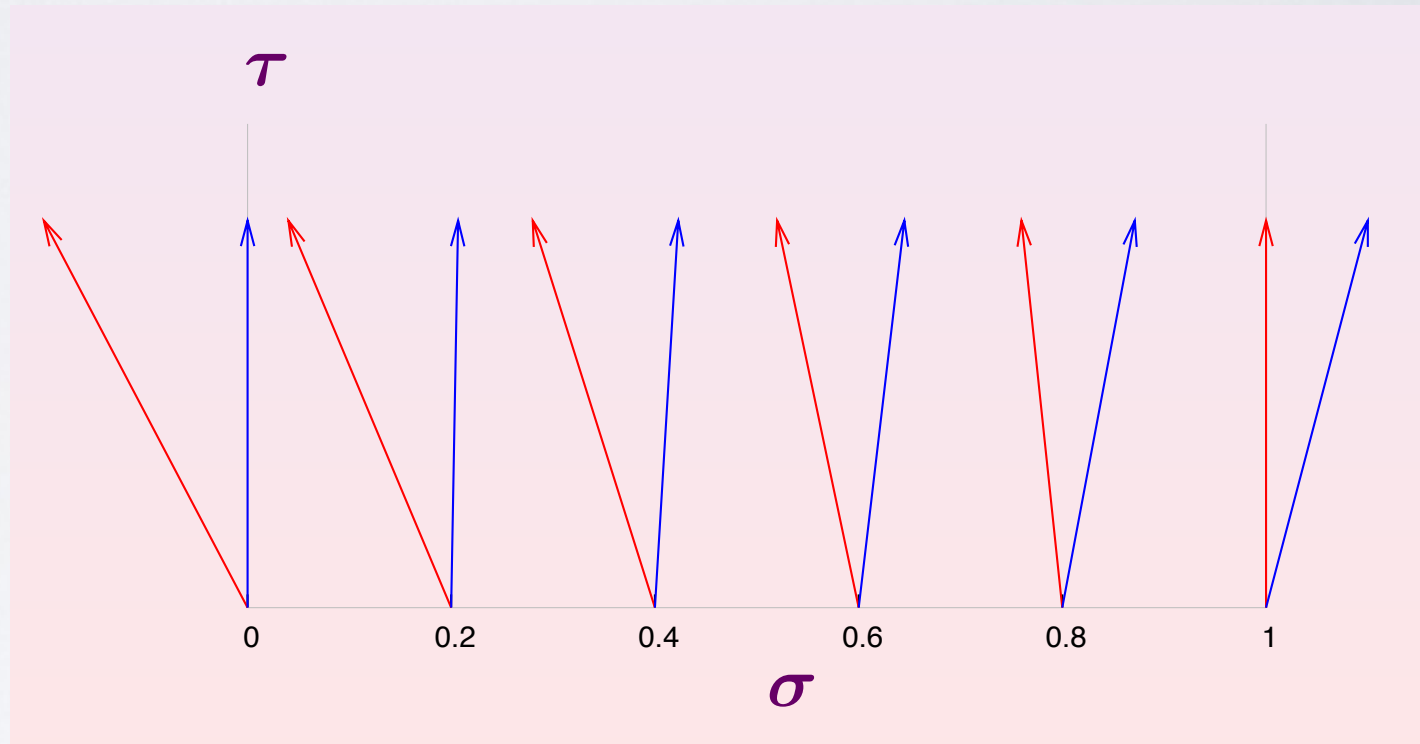
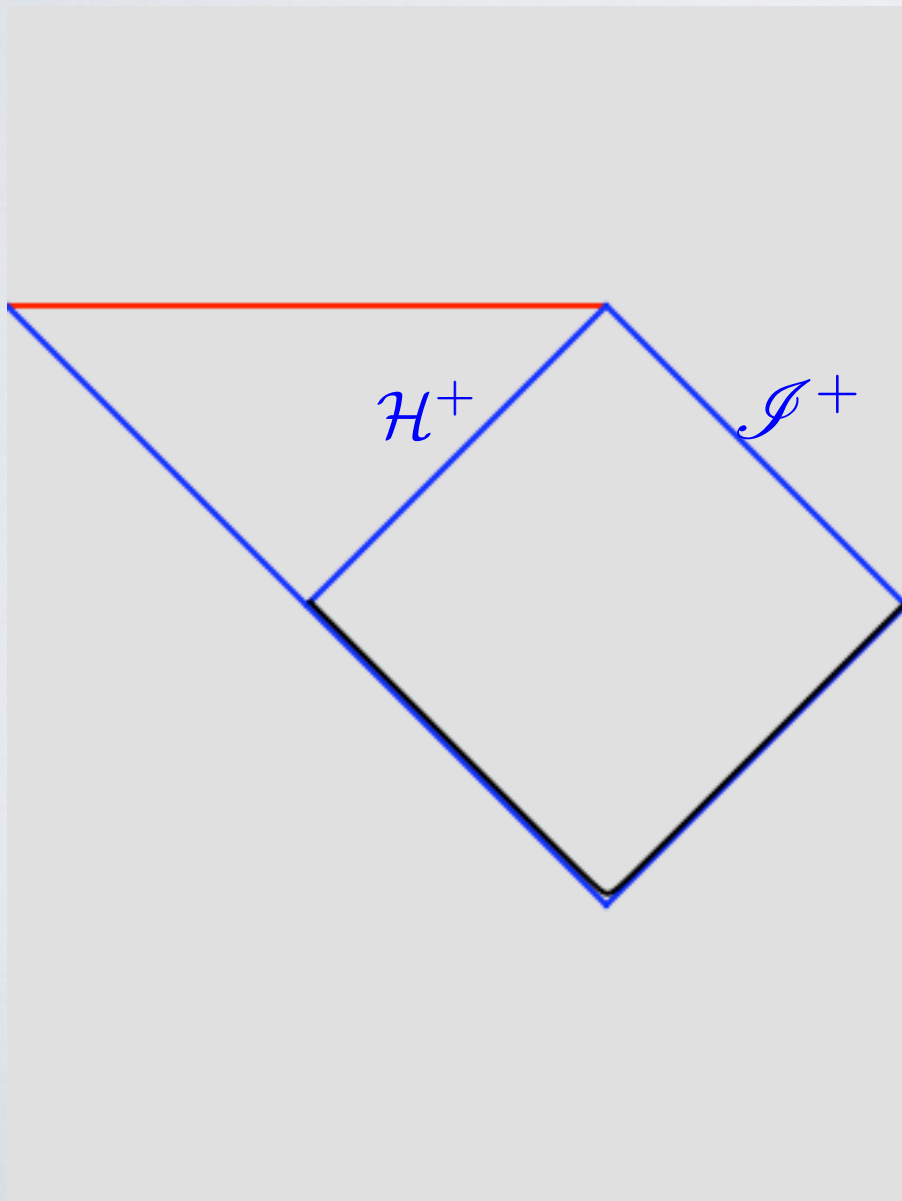


HYPERBOLOIDAL SLICES

- Schwarzschild Spacetime:

$$\sigma = \frac{2M}{r}$$

- Location of $\mathcal{H}^+$: $\sigma = 1$
- Location of $\mathcal{I}^+$: $\sigma = 0$



HYPERBOLOIDAL SLICES

- Generic Properties:

$$g_{ab} = \Omega^{-2} \tilde{g}_{ab}$$

HYPERBOLOIDAL SLICES

- Generic Properties:

$$g_{ab} = \Omega^{-2} \tilde{g}_{ab}$$

Physical metric
Irregular at $\mathcal{I}^+$

HYPERBOLOIDAL SLICES

- Generic Properties:

$$g_{ab} = \Omega^{-2} \tilde{g}_{ab}$$

Physical metric
Irregular at $\mathcal{I}^+$

“non-physical” metric
regular at $\mathcal{I}^+$

HYPERBOLOIDAL SLICES

- Generic Properties:

$$g_{ab} = \Omega^{-2} \tilde{g}_{ab}$$

Physical metric
Irregular at $\mathcal{I}^+$

“non-physical” metric
regular at $\mathcal{I}^+$

$\mathcal{I}^+$ located at $\Omega = 0$

The diagram illustrates the relationship between two metrics, g_{ab} and $\tilde{g}_{ab}$, in the context of hyperboloidal slices. The equation $g_{ab} = \Omega^{-2} \tilde{g}_{ab}$ is shown at the top. Below it, three arrows point from the terms in the equation to descriptive text. An arrow from g_{ab} points to the text 'Physical metric' and 'Irregular at $\mathcal{I}^+$ '. An arrow from $\tilde{g}_{ab}$ points to the text '“non-physical” metric' and 'regular at $\mathcal{I}^+$ '. A third arrow points from the Ω^{-2} term down to the text ' $\mathcal{I}^+$ located at $\Omega = 0$ '.

HYPERBOLOIDAL SLICES

- Generic Properties:

$$g_{ab} = \Omega^{-2} \tilde{g}_{ab}$$

Physical metric
Irregular at $\mathcal{I}^+$

“non-physical” metric
regular at $\mathcal{I}^+$

$\mathcal{I}^+$ located at $\Omega = 0$

✓ $\mathcal{I}^+$ is casually disconnected from numerical domain: no need of artificial boundary condition for time evolution

HYPERBOLOIDAL SLICES

- Generic Properties:

$$g_{ab} = \Omega^{-2} \tilde{g}_{ab}$$

Physical metric
Irregular at $\mathcal{I}^+$

“non-physical” metric
regular at $\mathcal{I}^+$

$\mathcal{I}^+$ located at $\Omega = 0$

- ✓ $\mathcal{I}^+$ is casually disconnected from numerical domain: no need of artificial boundary condition for time evolution
- ✓ $\mathcal{I}^+$ formal mathematical definition of the radiation zone: clean, well-defined notion of gravitational waves

HYPERBOLOIDAL SLICES

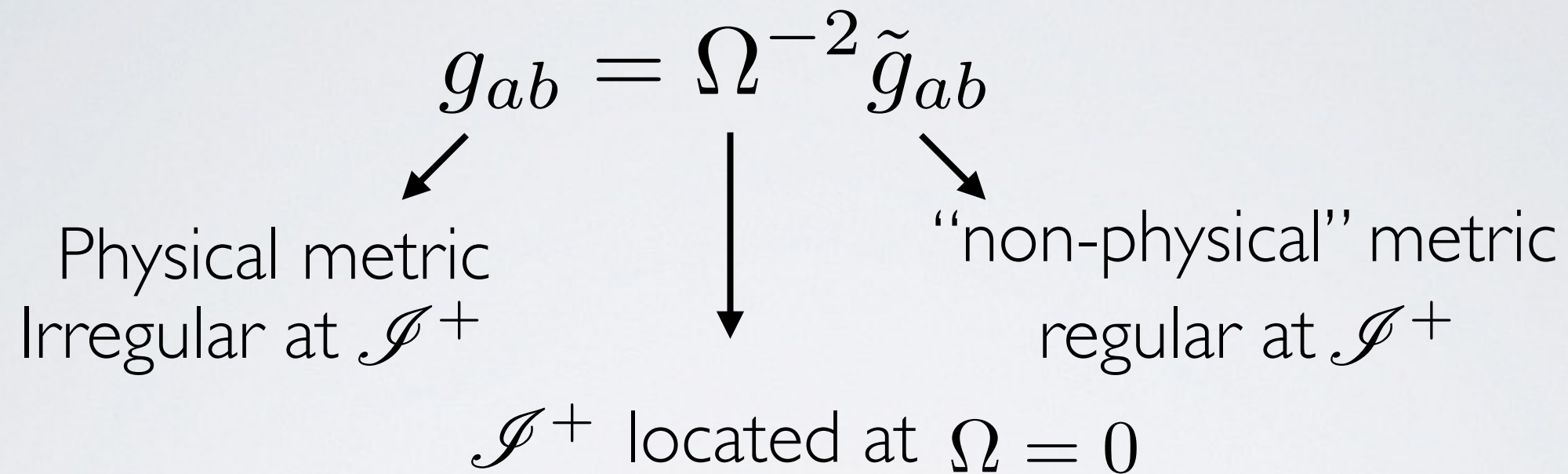
- Generic Properties:

$$g_{ab} = \Omega^{-2} \tilde{g}_{ab}$$

Physical metric
Irregular at $\mathcal{I}^+$

“non-physical” metric
regular at $\mathcal{I}^+$

$\mathcal{I}^+$ located at $\Omega = 0$



- * Einstein's equation pick up Ω^{-2} terms, thus irregular at $\mathcal{I}^+$
- * One needs robust numerical methods to deal with them

INITIAL-DATA PROBLEM

Physical meaning

- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)
- **Elliptic Equations : Boundary value problem**

$\mathcal{S}_{\text{out}}$: infinity (“ $r \rightarrow \infty$ ”)

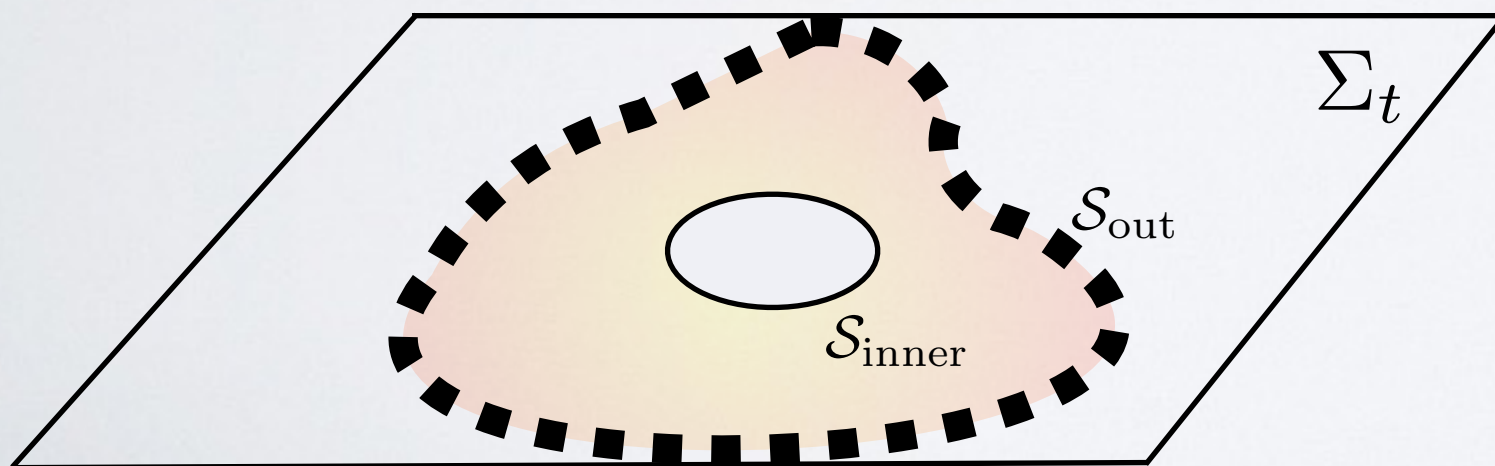
$$\gamma_{ij} = \Omega^{-2} \tilde{\gamma}_{ij}$$

$$\mathcal{S}_{\text{out}}: \Omega = 0$$

- **Boundary Conditions:**

Regularity conditions of

$$\mathbb{H} = 0 \quad \mathbb{P}_i = 0$$

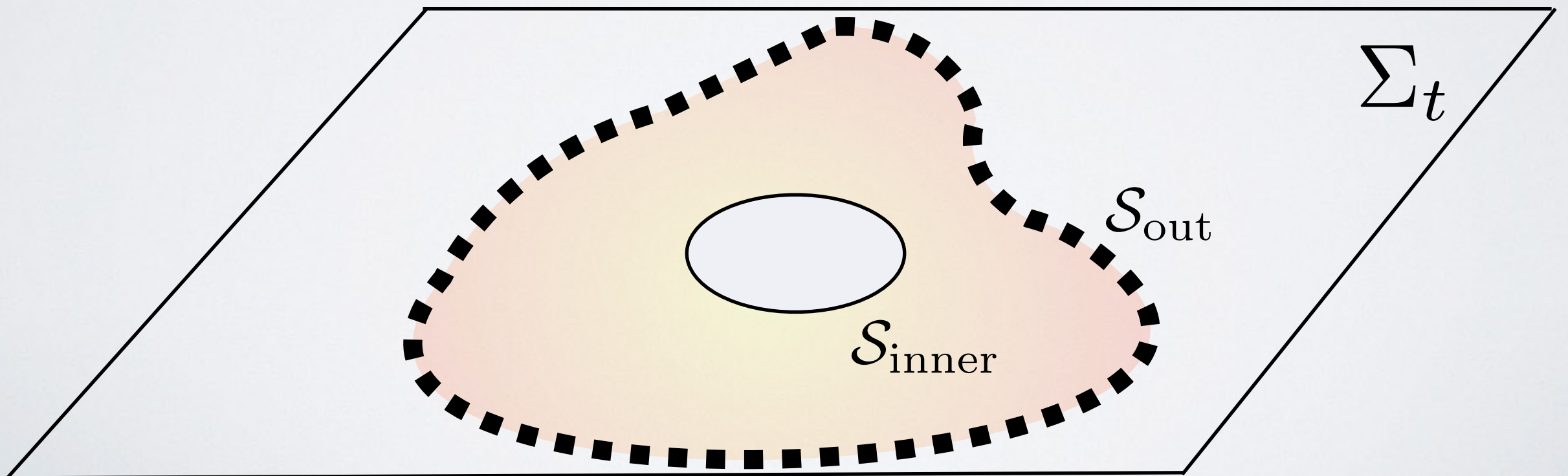


(L. Andersson, P. Chrusciel and H. Friedrich '92,
L. Andersson and P. Chrusciel '93 '94)

INITIAL-DATA PROBLEM

Physical meaning

- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)
- **Elliptic Equations : Boundary value problem**
 $\mathcal{S}_{\text{inner}}$: Black-hole horizon $\mathcal{S}_{\text{out}}$: infinity (“ $r \rightarrow \infty$ ”)



INITIAL-DATA PROBLEM

Physical meaning

- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)

- **Elliptic Equations : Boundary value problem**

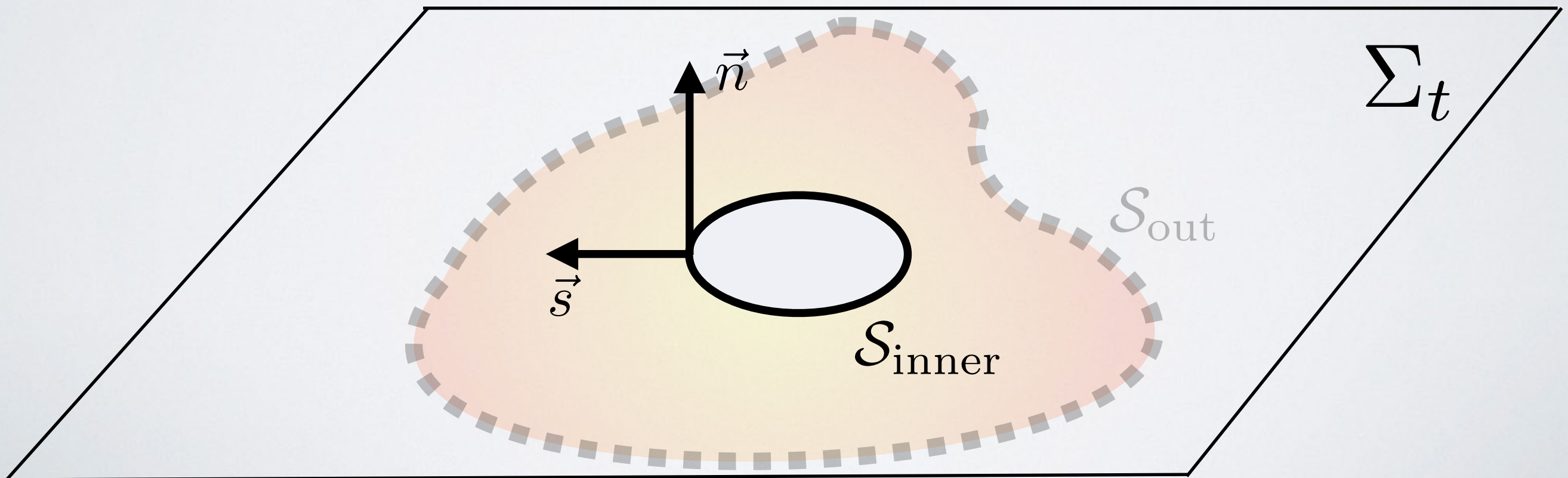
$\mathcal{S}_{\text{inner}}$: Inner Boundary (Black-hole horizon) $\mathcal{S}_{\text{outer}}$: Outer Boundary (Infinity (“ $r \rightarrow \infty$ ”))

$\mathcal{S}_{\text{inner}}$: Black-hole horizon

- Local definition of horizon?
- How to perturb the black hole?

$\mathcal{S}_{\text{inner}}$

PERTURBED BLACK HOLES

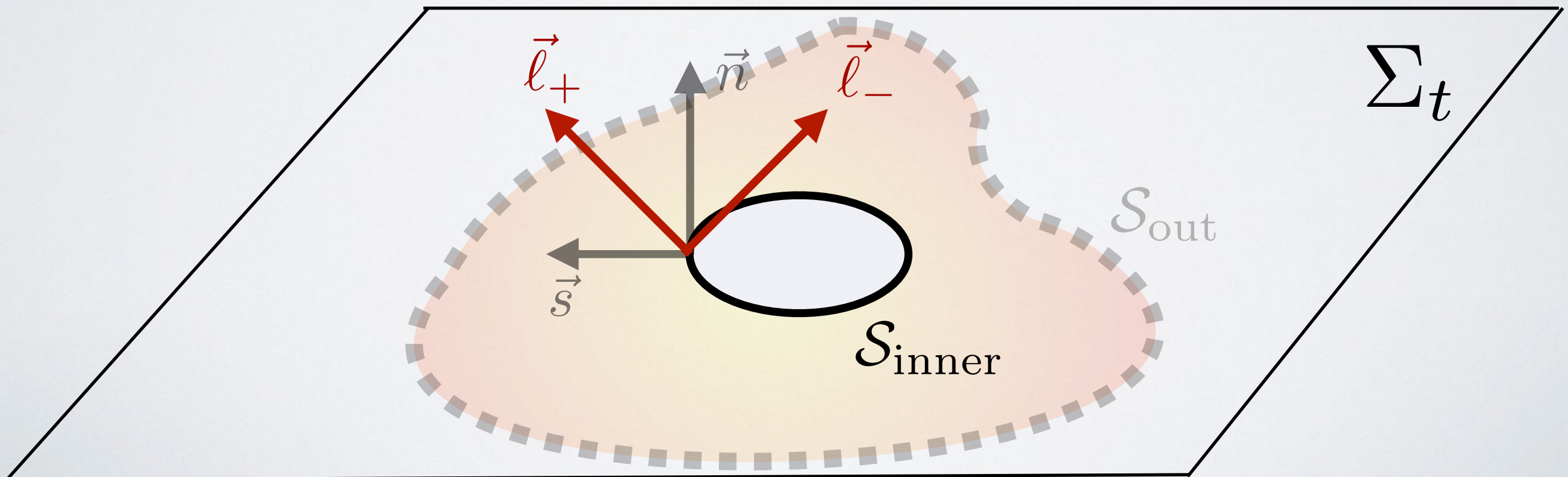


PERTURBED BLACK HOLES

- Outgoing / Ingoing light-rays directions: $\vec{\ell}_+ / \vec{\ell}_-$

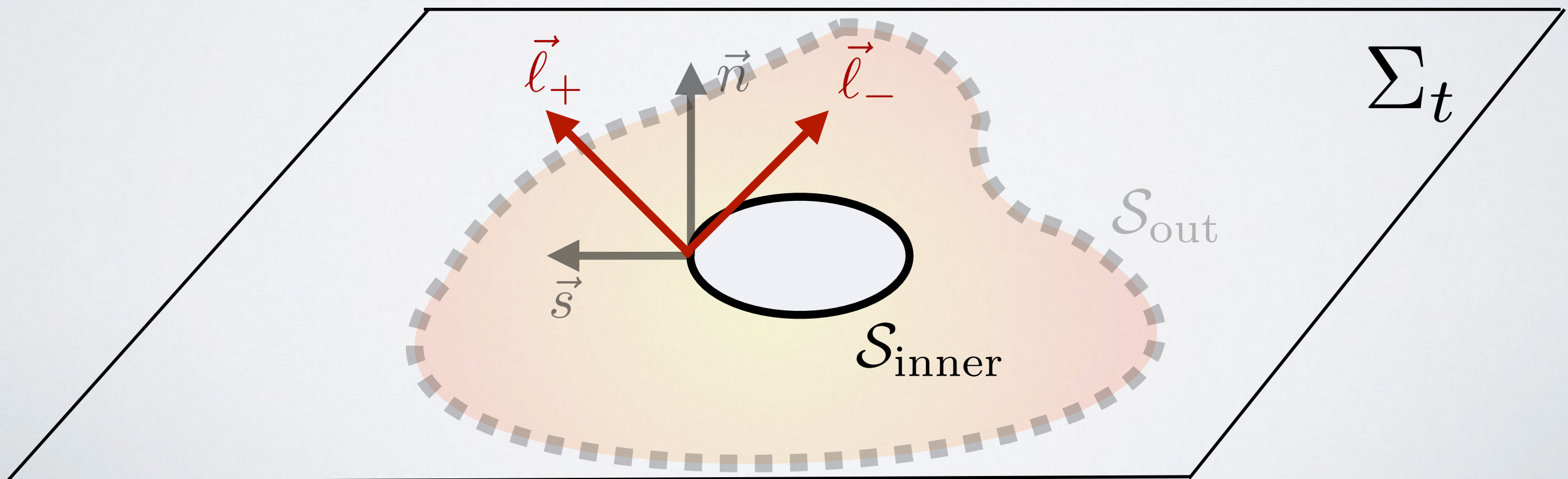
- Expansion of outgoing / ingoing light-rays:

$$\Theta_{\pm} = q_{ij} \nabla^i \ell_{\pm}^j \quad \text{with} \quad q_{ij} = \gamma_{ij}|_{\mathcal{S}_{\text{inner}}}$$



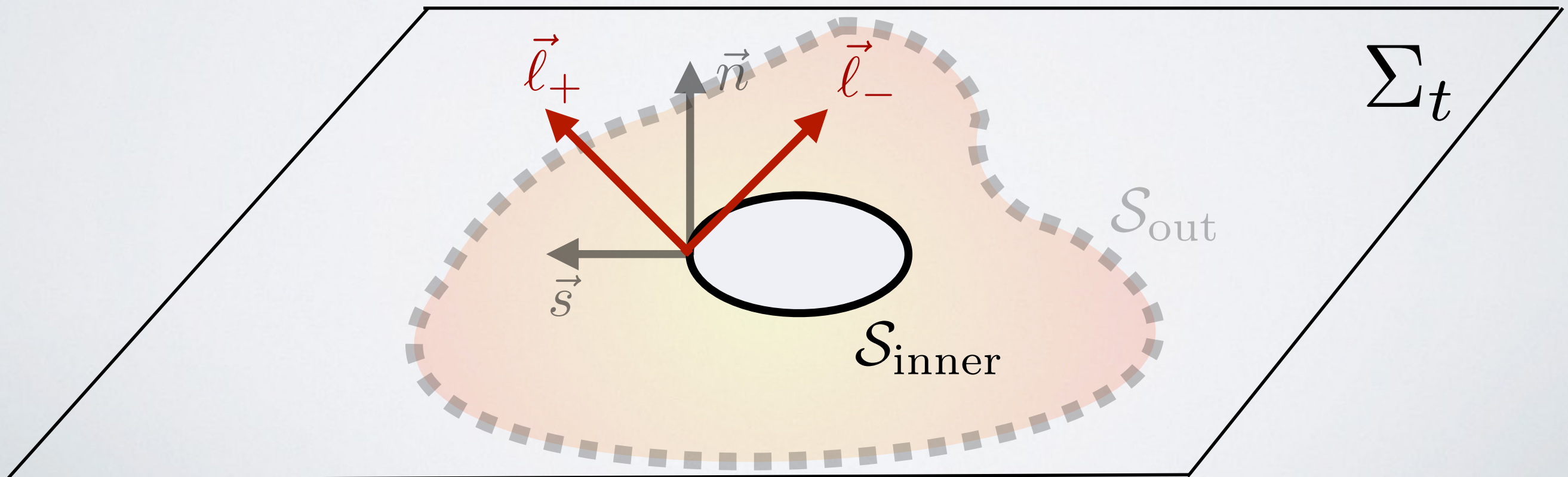
PERTURBED BLACK HOLES

- Outgoing / Ingoing light-rays directions: $\vec{\ell}_+ / \vec{\ell}_-$
- Expansion of outgoing / ingoing light-rays:
$$\Theta_{\pm} = K - s^i s^j \pm {}^{(3)}\nabla_i s^i$$
- **Black Hole apparent Horizont: $\Theta_+ = 0$**
(outgoing light-rays stop expanding outwards)



PERTURBED BLACK HOLES

- Outgoing / Ingoing light-rays directions: $\vec{\ell}_+ / \vec{\ell}_-$
- Expansion of outgoing / ingoing light-rays:
$$\Theta_{\pm} = K - s^i s^j \pm {}^{(3)}\nabla_i s^i$$
- **White Hole apparent Horizont: $\Theta_- = 0$**
(ingoing light-rays stop “expanding inwards” - contracting)



INITIAL-DATA PROBLEM

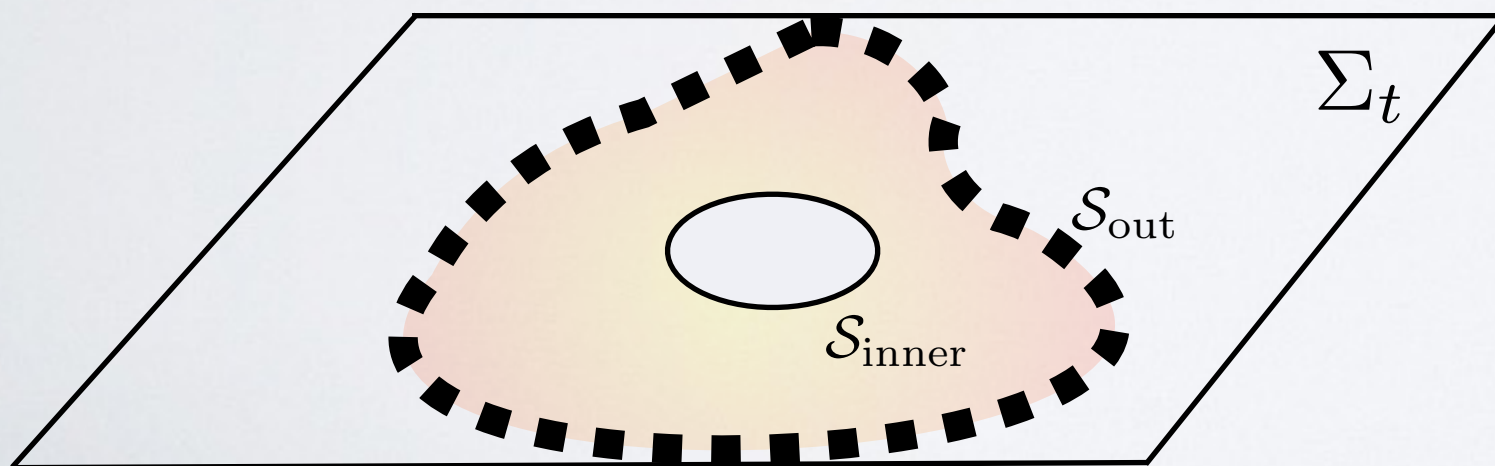
Physical meaning

- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)

- **Elliptic Equations : Boundary value problem**

$\mathcal{S}_{\text{inner}}$: Black-hole horizon

$$\mathcal{S}_{\text{inner}} : \Theta_+ = 0$$



- **Boundary Condition:**

$$K - s^i s^j \pm {}^{(3)}\nabla_i s^i = 0$$

PERTURBED BLACK HOLES

- Variables: γ_{ij} K_{ij} (2 , 3X3 symmetric tensors)

PERTURBED BLACK HOLES

- Variables: γ_{ij} K_{ij} (2 , 3X3 symmetric tensors)

➡ Total of: $2 \times 6 = 12$ unknown variables

PERTURBED BLACK HOLES

- Variables: γ_{ij} K_{ij} (2 , 3X3 symmetric tensors)
 - ➡ Total of: $2 \times 6 = 12$ unknown variables
- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)

PERTURBED BLACK HOLES

- Variables: γ_{ij} K_{ij} (2 , 3X3 symmetric tensors)
 - ➡ Total of: $2 \times 6 = 12$ unknown variables
- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)
 - ➡ Total of: 4 equations

PERTURBED BLACK HOLES

- Variables: γ_{ij} K_{ij} (2 , 3X3 symmetric tensors)
 - ➡ Total of: $2 \times 6 = 12$ unknown variables
- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)
 - ➡ Total of: 4 equations
- We can **freely** specify 8 quantities

PERTURBED BLACK HOLES

- Variables: γ_{ij} K_{ij} (2, 3X3 symmetric tensors)
 - ➡ Total of: $2 \times 6 = 12$ unknown variables
- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)
 - ➡ Total of: 4 equations
- We can **freely** specify 8 quantities
- Constraint equations (with boundary conditions) fix the other 4 variables

PERTURBED BLACK HOLES

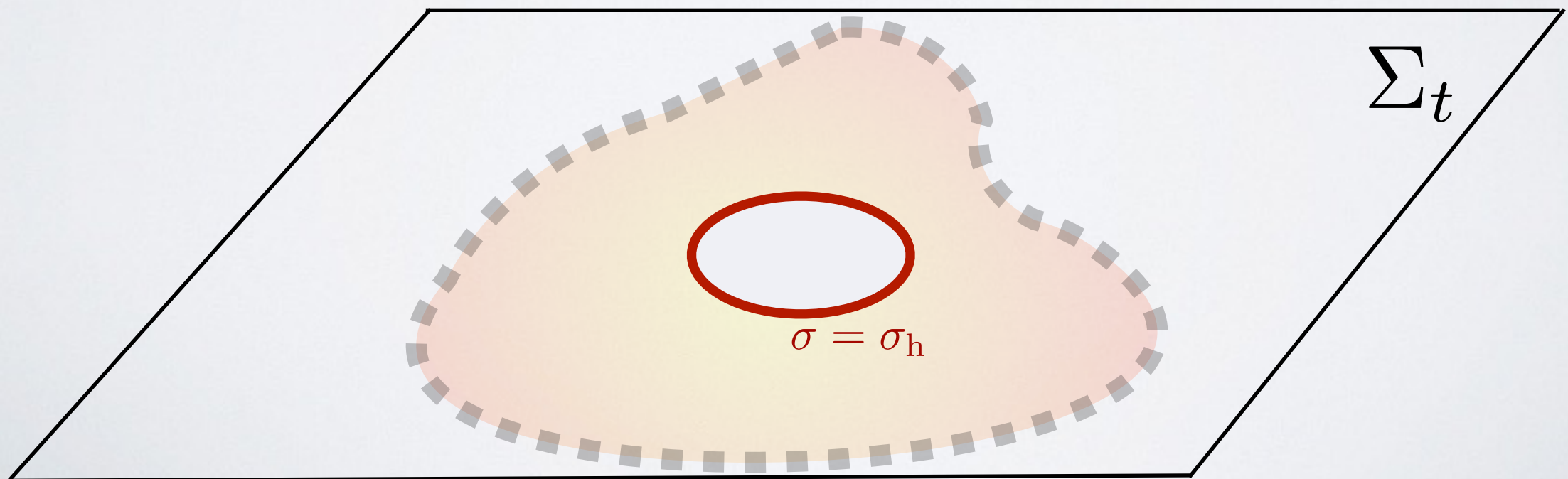
- Variables: γ_{ij} K_{ij} (2, 3X3 symmetric tensors)
 - ➡ Total of: $2 \times 6 = 12$ unknown variables
- Constraint Equations : $\mathbb{H} = 0$ (Hamiltonian Constraint)
 $\mathbb{P}_i = 0$ (Momentum Constraint)
 - ➡ Total of: 4 equations
- We can **freely** specify 8 quantities
- Constraint equations (with boundary conditions) fix the other 4 variables

“Conformal transverse traceless decomposition”

PERTURBED BLACK HOLES

- We can **freely** specify 8 quantities :

Read from Kerr space-time ($j = \frac{a}{M}$) Black Hole
in hyperboloidal coordinates $\Rightarrow \sigma = \sigma_h$



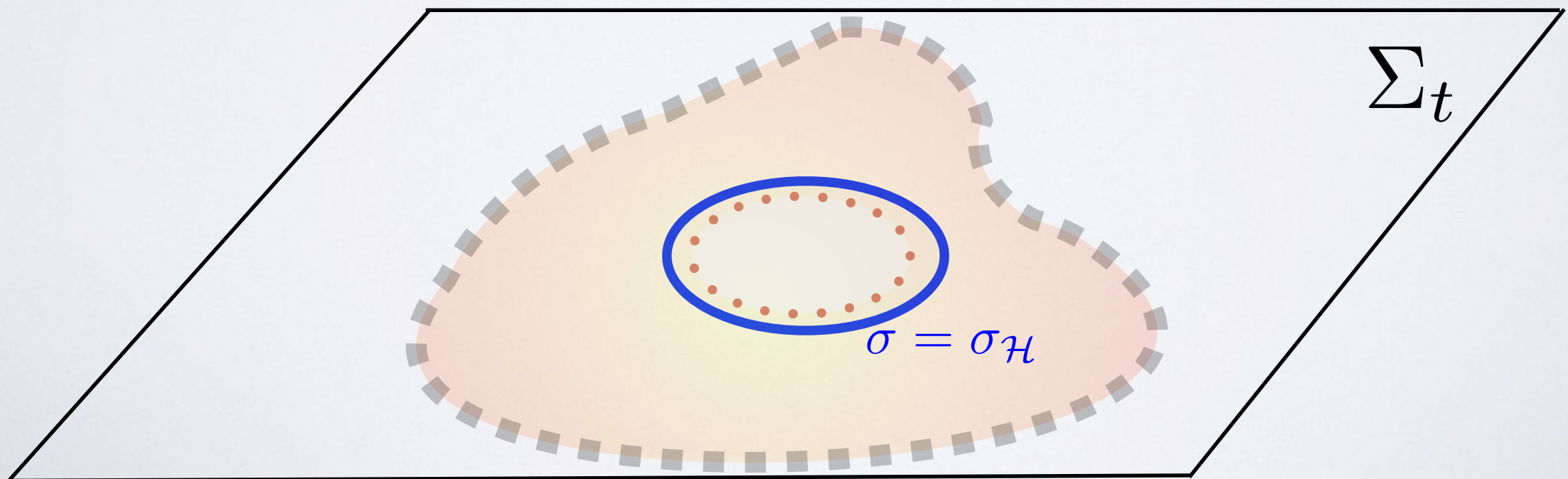
PERTURBED BLACK HOLES

- We can **freely** specify 8 quantities :

Read from Kerr space-time ($j = \frac{a}{M}$) Black Hole
in hyperboloidal coordinates $\Rightarrow \sigma = \sigma_h$

- Constraint equations (with boundary conditions) fix the other 4 variables

$$\mathcal{S}_{\text{inner}} : \Theta_+ = 0 \quad \text{at} \quad \sigma = \sigma_{\mathcal{H}} \neq \sigma_h$$



PERTURBED BLACK HOLES

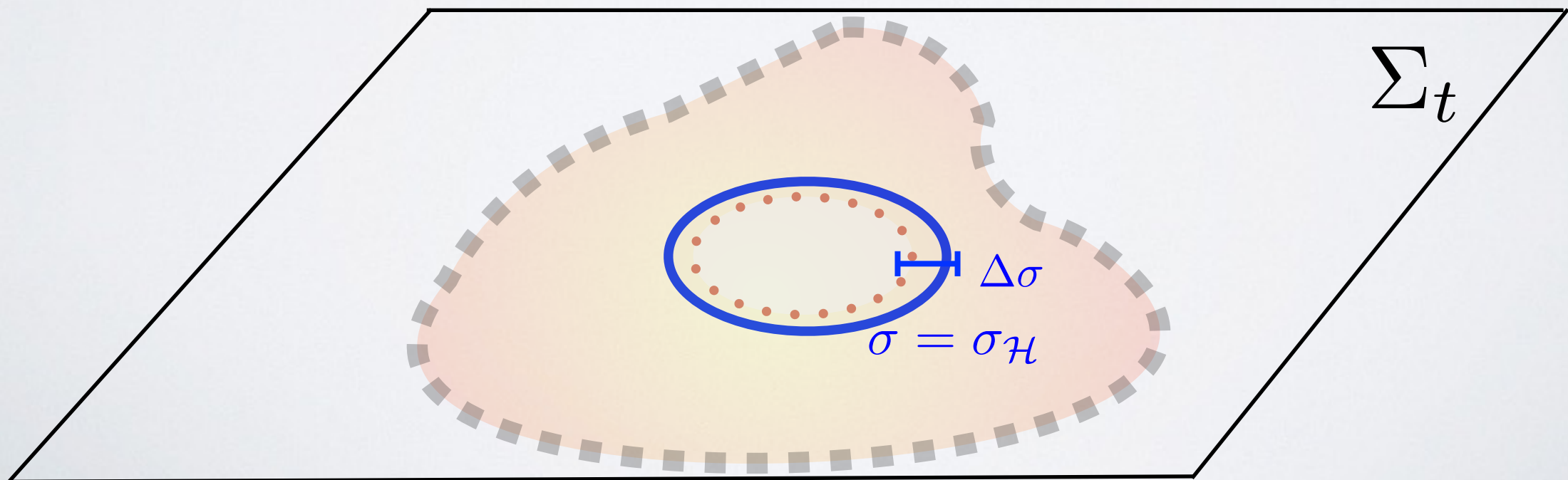
- We can **freely** specify 8 quantities :

Read from Kerr space-time ($j = \frac{a}{M}$) Black Hole
in hyperboloidal coordinates $\Rightarrow \sigma = \sigma_h$

- Constraint equations (with boundary conditions) fix the other 4 variables

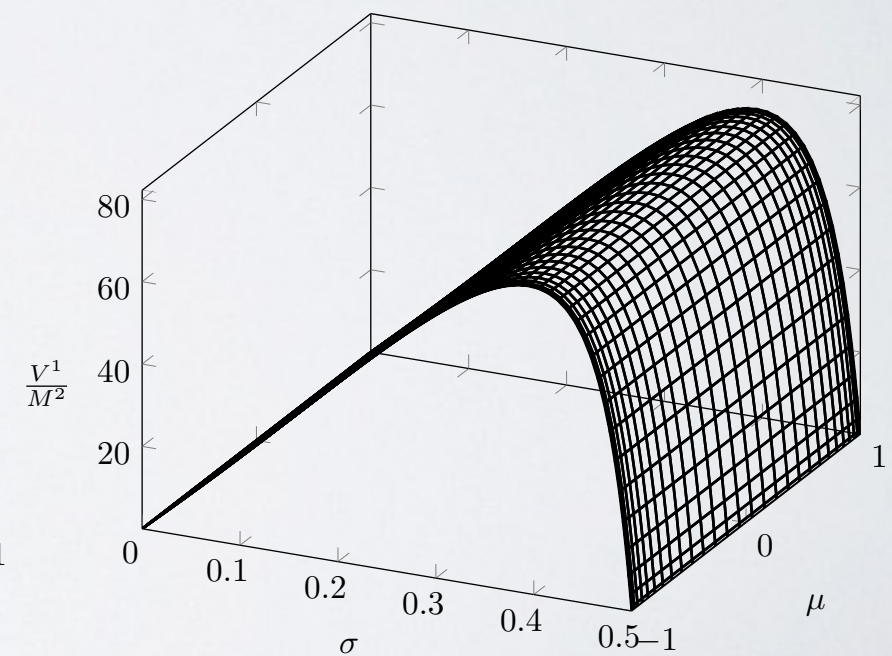
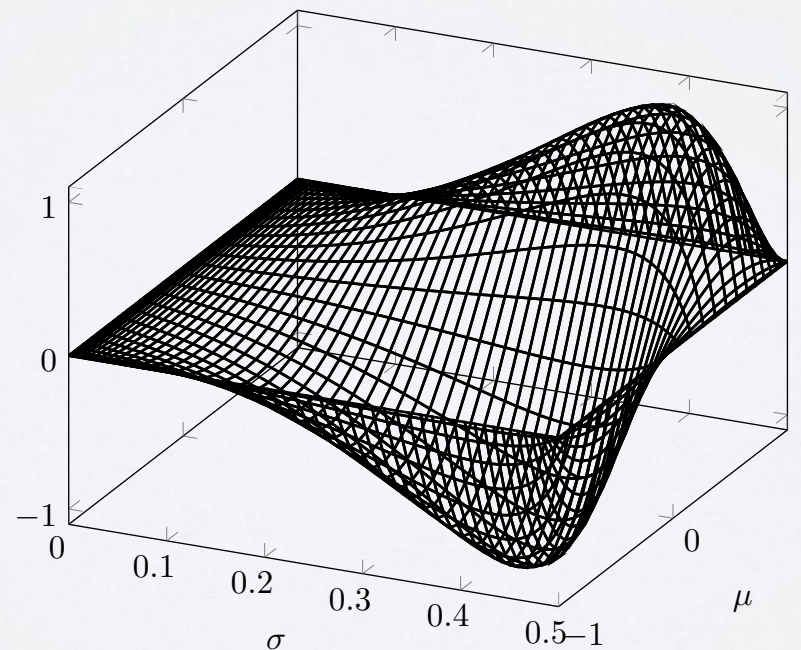
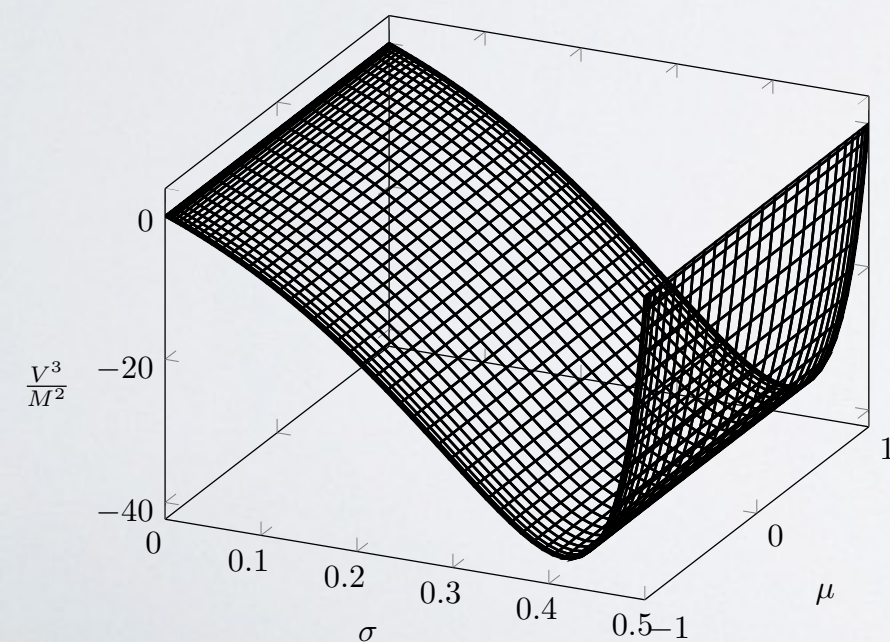
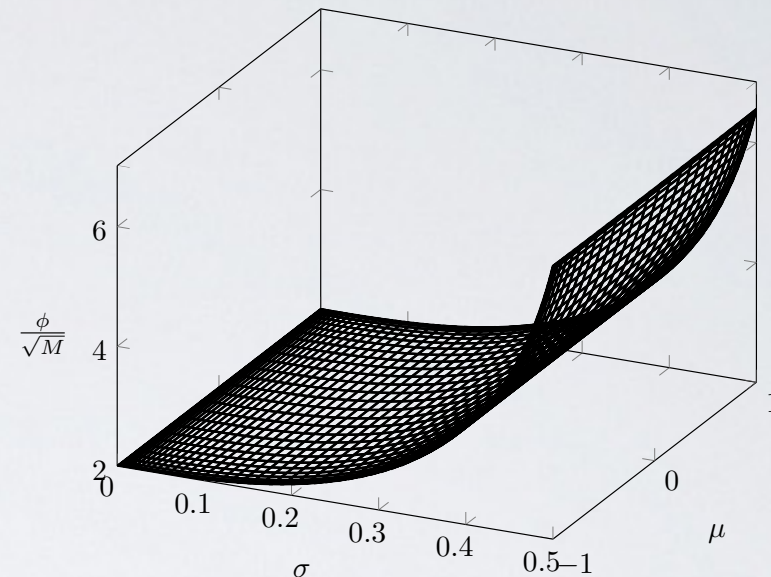
$$\mathcal{S}_{\text{inner}} : \Theta_+ = 0 \quad \text{at} \quad \sigma = \sigma_{\mathcal{H}} \neq \sigma_h$$

$$\text{Perturbation parameter: } \Delta\sigma = \sigma_h - \sigma_{\mathcal{H}}$$



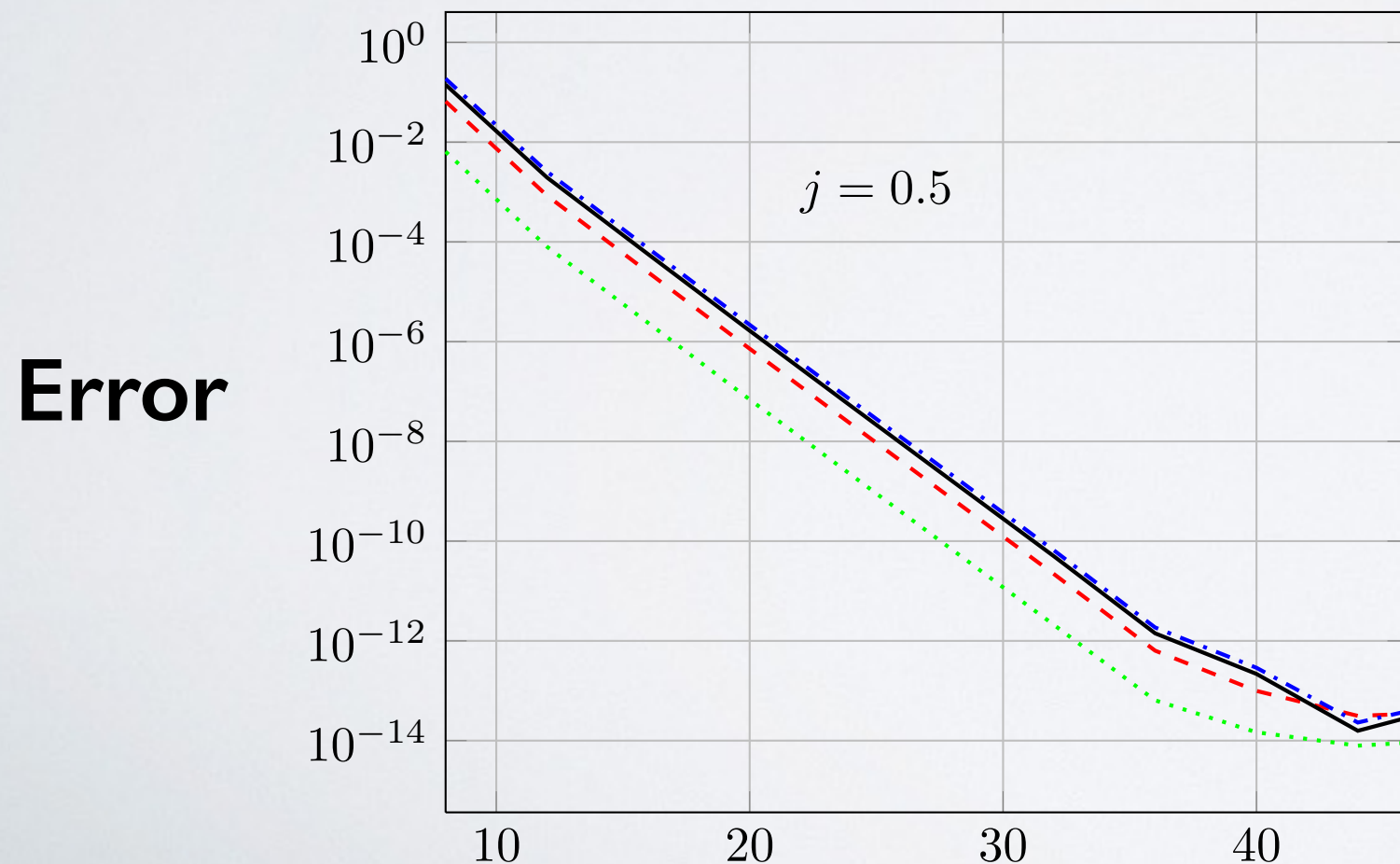
NUMERICAL SOLUTION

- Axisymmetry
- Pseudo-Spectral methods
- Gauß-Lobatto grid points
- Newton-Raphson method
- BicCGStab with finite-difference pre-conditioner



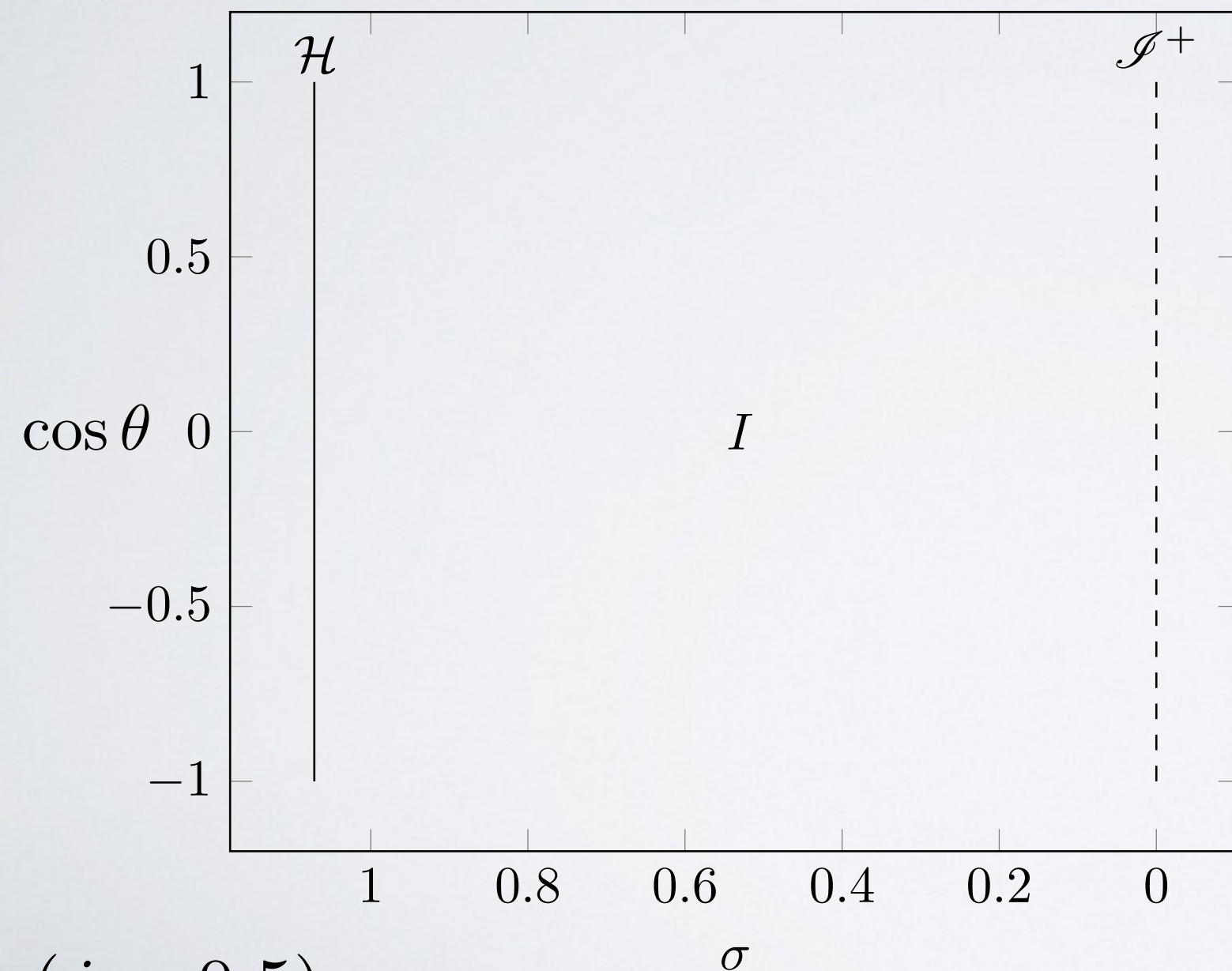
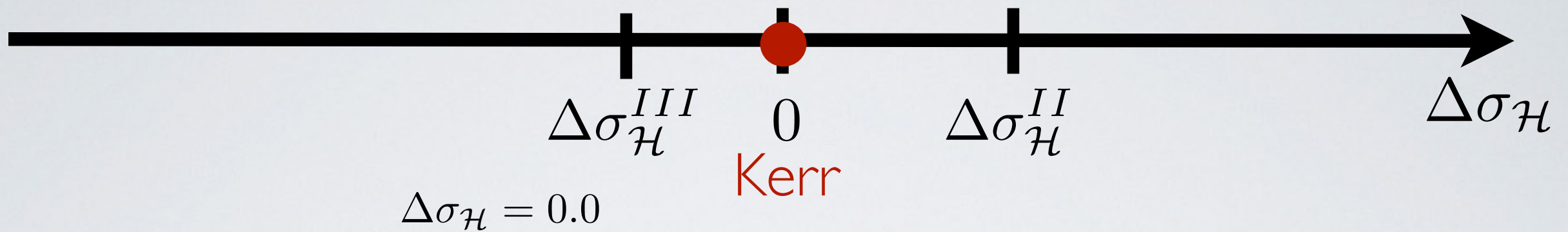
NUMERICAL SOLUTION

- Axisymmetry
- Pseudo-Spectral methods
- Gauß-Lobatto grid points
- Newton-Raphson method
- BicCGStab with finite-difference pre-conditioner



**Accuracy down to
machine precision:
11~12 decimal digits**

Trapped Regions



$(j = 0.5)$

- Numerical Boundaries

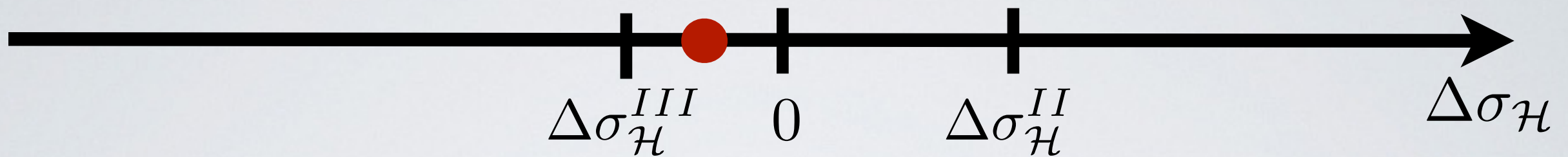
$$\mathcal{J}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- Region I:

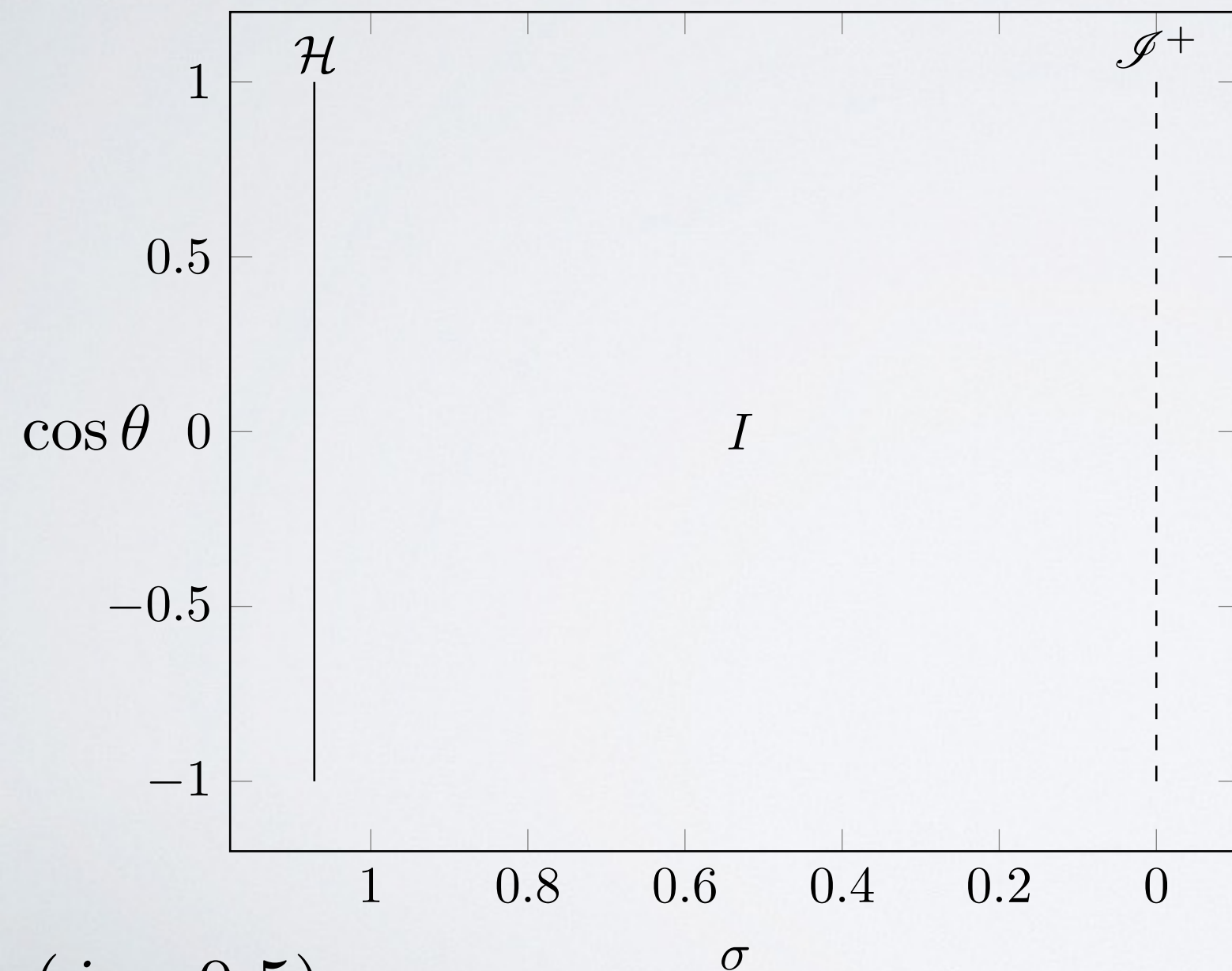
$$\Theta_+ > 0 \quad \Theta_- < 0$$

Exterior solution

Trapped Regions



$$\Delta\sigma_{\mathcal{H}} = 0.0$$



$$(j = 0.5)$$

- Numerical Boundaries

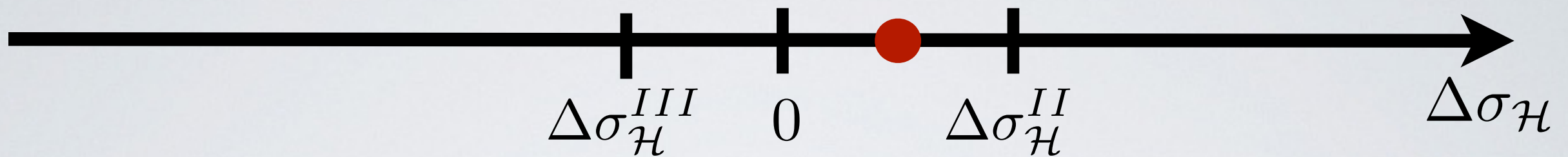
$$\mathcal{J}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- Region I:

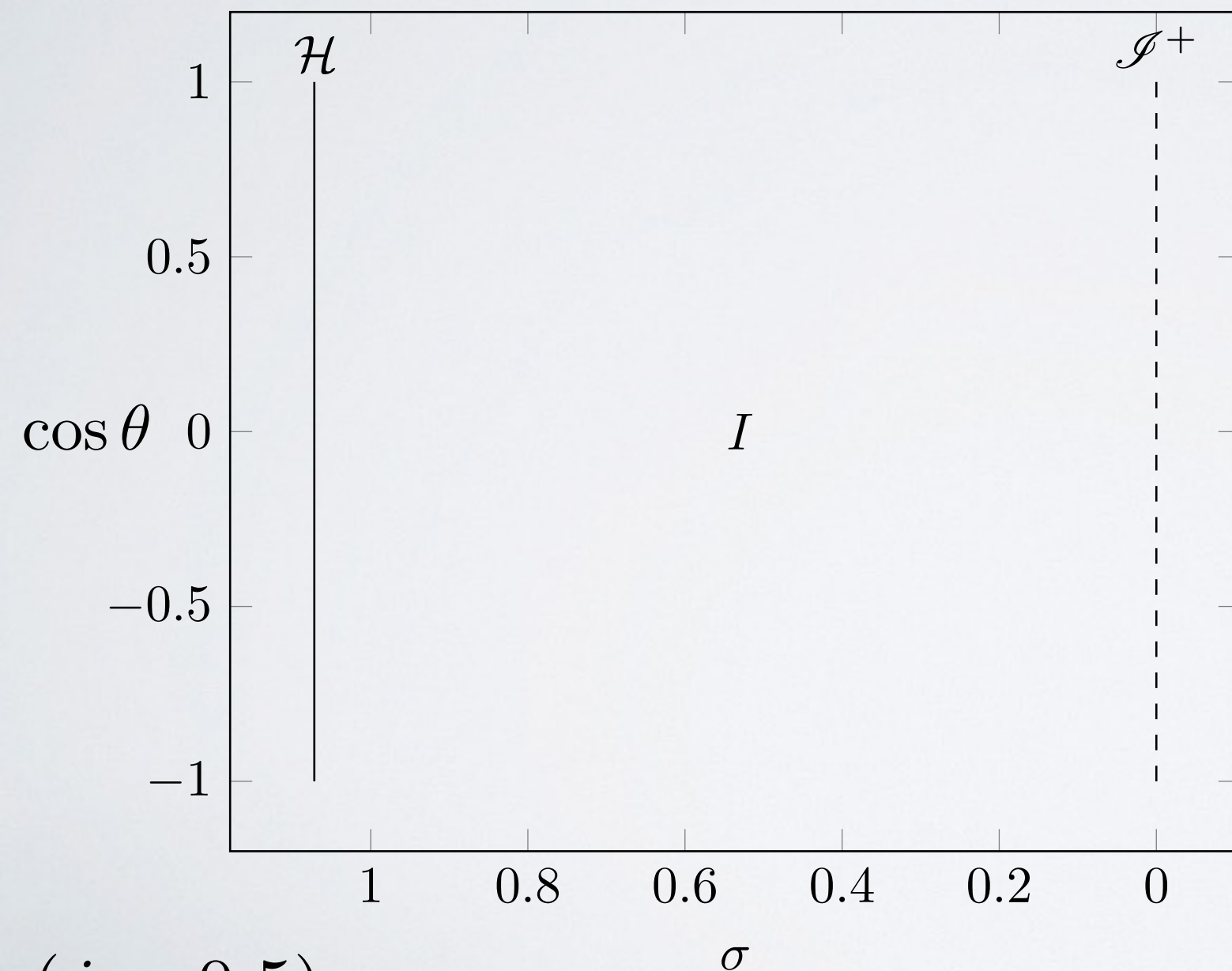
$$\Theta_+ > 0 \quad \Theta_- < 0$$

Exterior solution

Trapped Regions



$$\Delta\sigma_{\mathcal{H}} = 0.0$$



$$(j = 0.5)$$

- Numerical Boundaries

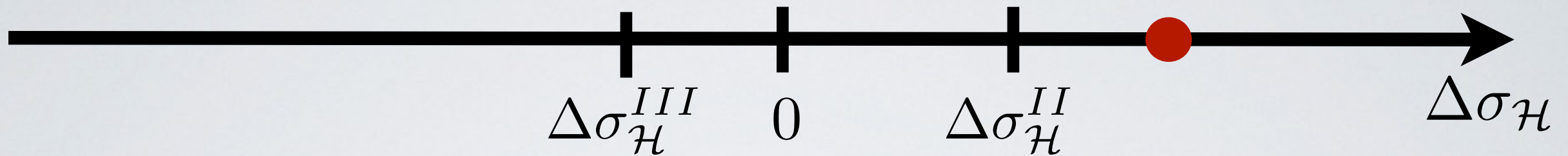
$$\mathcal{J}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- Region I:

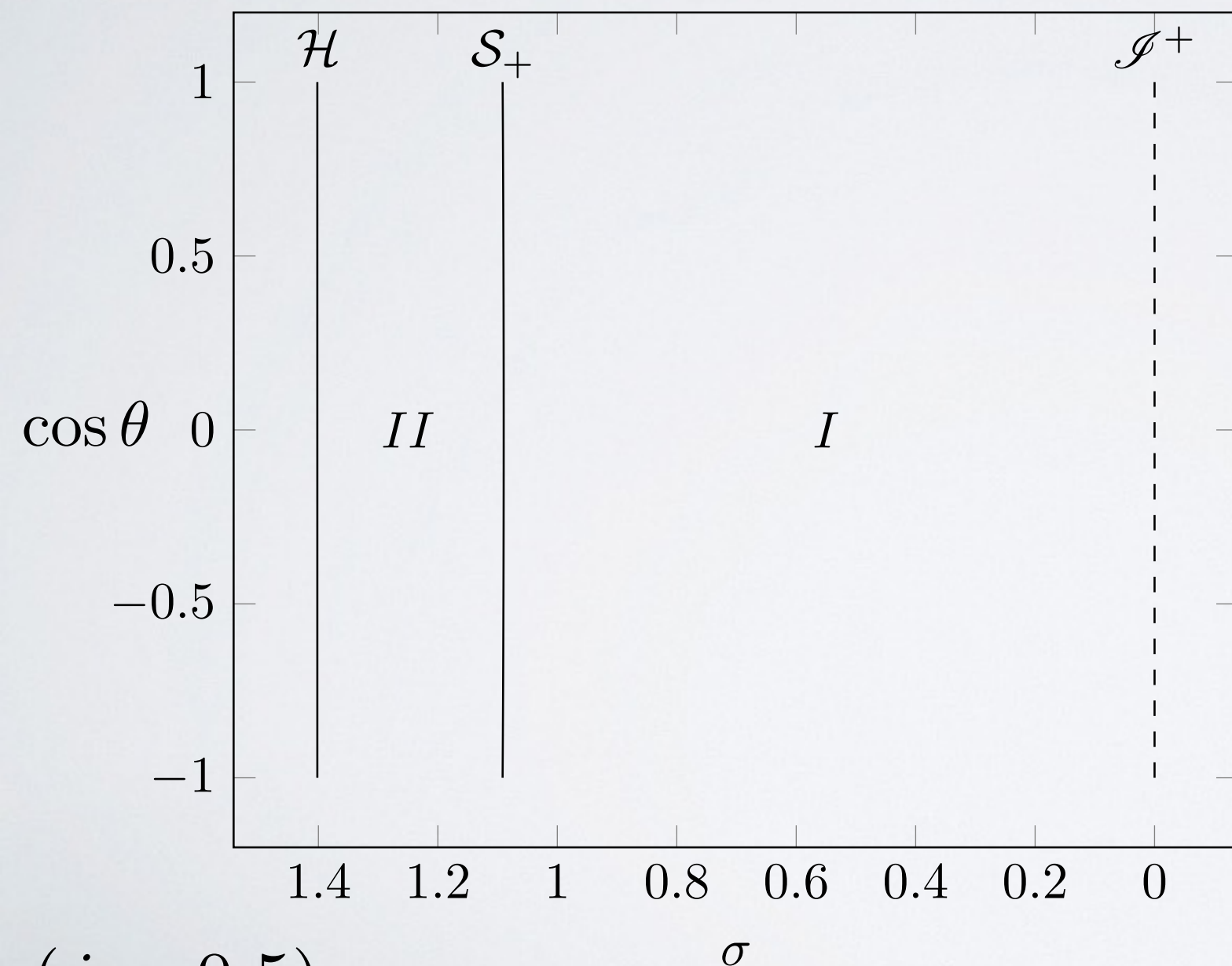
$$\Theta_+ > 0 \quad \Theta_- < 0$$

Exterior solution

Trapped Regions



$$\Delta\sigma_{\mathcal{H}} = 0.33$$



($j = 0.5$)

- Numerical Boundaries

$$\mathcal{J}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- New surface $\mathcal{S}_+$

$$\Theta_+ = 0$$

- Region I:

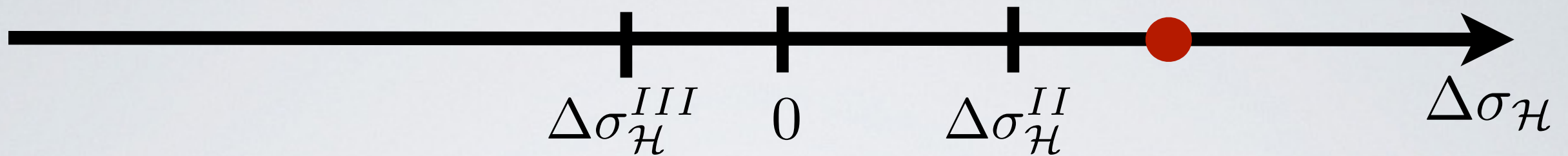
$$\Theta_+ > 0 \quad \Theta_- < 0$$

Exterior solution

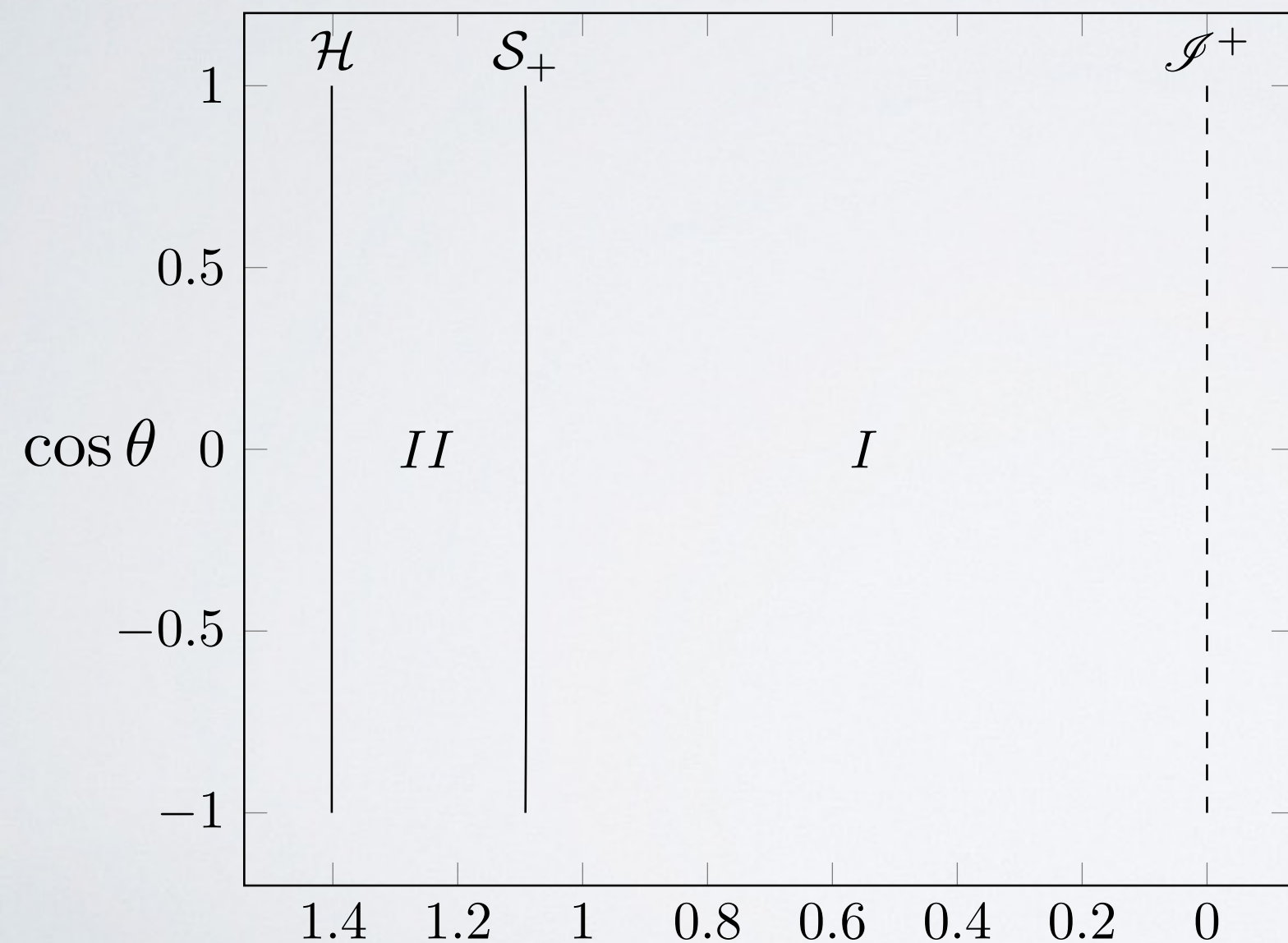
- Region II:

$$\Theta_+ < 0 \quad \Theta_- < 0$$

Trapped Regions



$$\Delta\sigma_{\mathcal{H}} = 0.33$$



$$(j = 0.5)$$

- Numerical Boundaries

$$\mathcal{J}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- New surface $\mathcal{S}_+$

$$\Theta_+ = 0$$

- Region I:

$$\Theta_+ > 0 \quad \Theta_- < 0$$

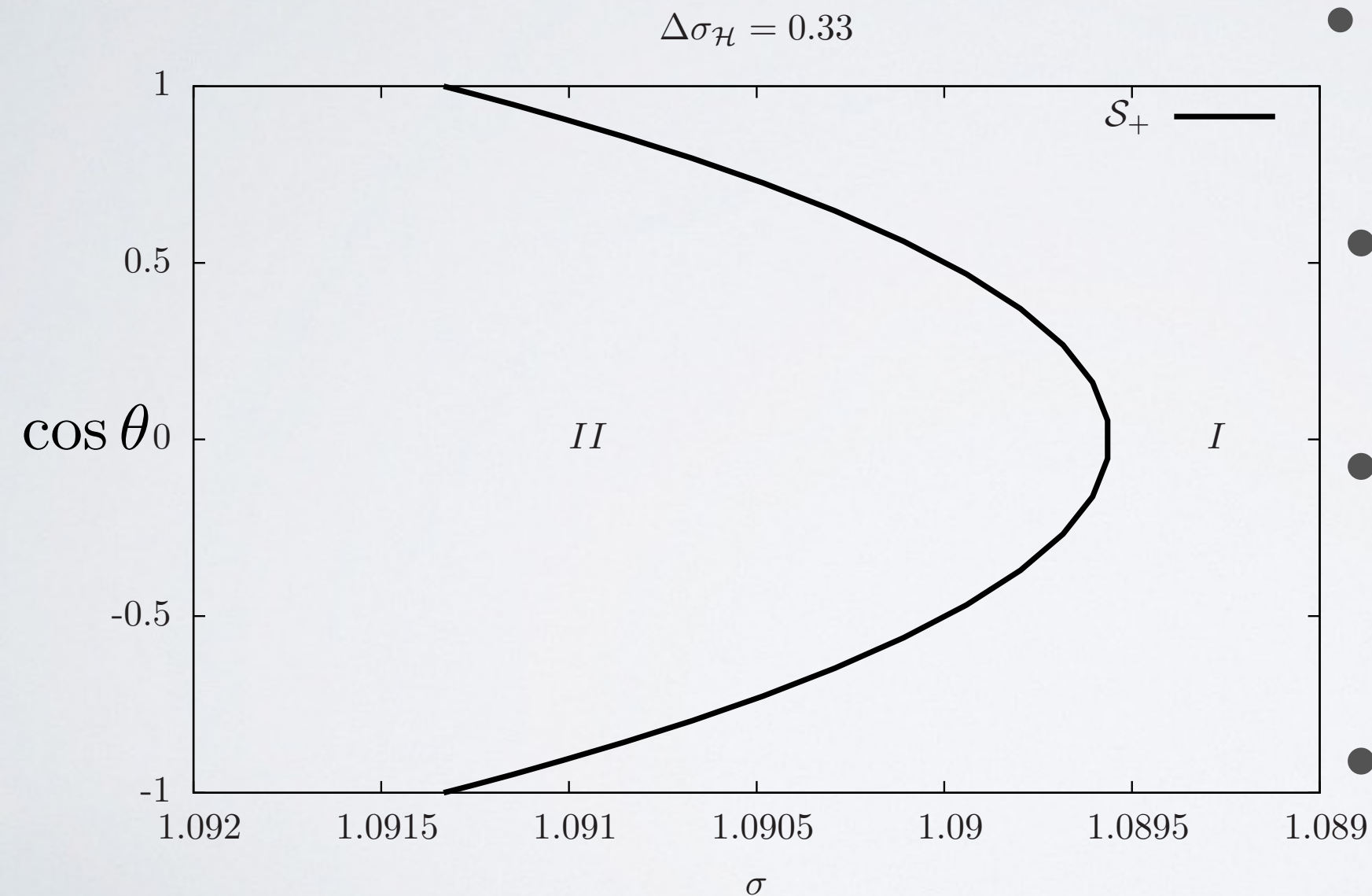
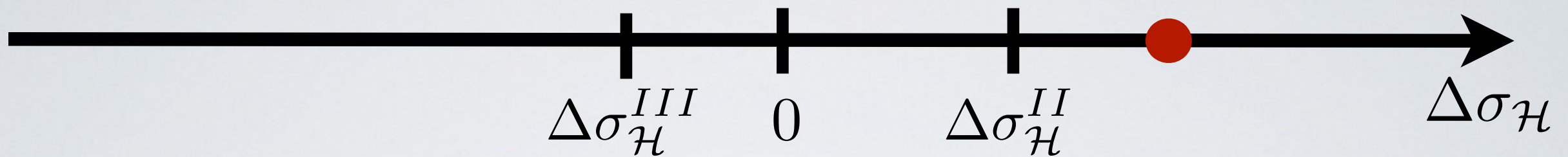
Exterior solution

- Region II:

$$\Theta_+ < 0 \quad \Theta_- < 0$$

Black Hole

Trapped Regions



- Numerical Boundaries

$$\mathcal{I}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- New surface $\mathcal{S}_+$

$$\Theta_+ = 0$$

- Region I:

$$\Theta_+ > 0 \quad \Theta_- < 0$$

Exterior solution

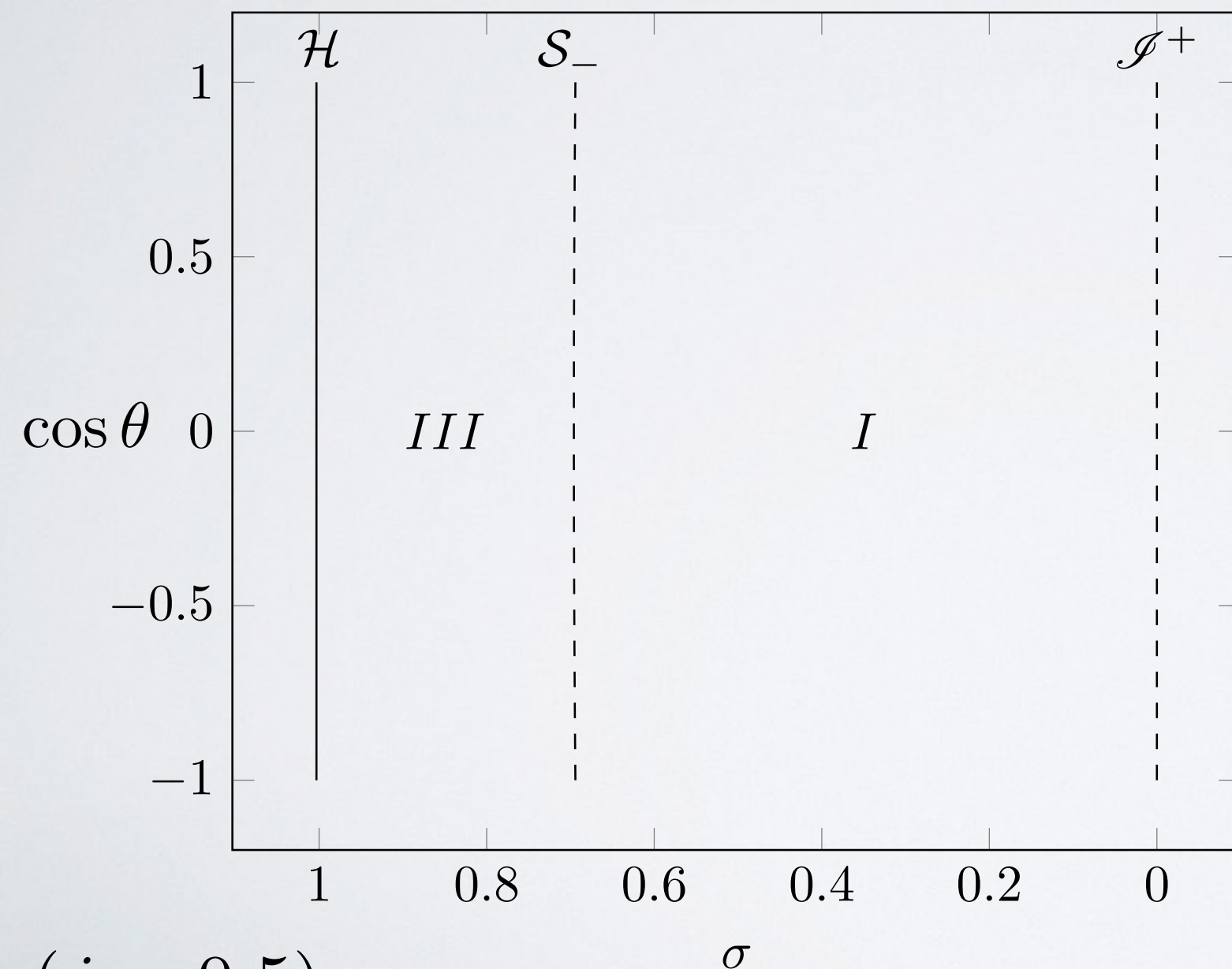
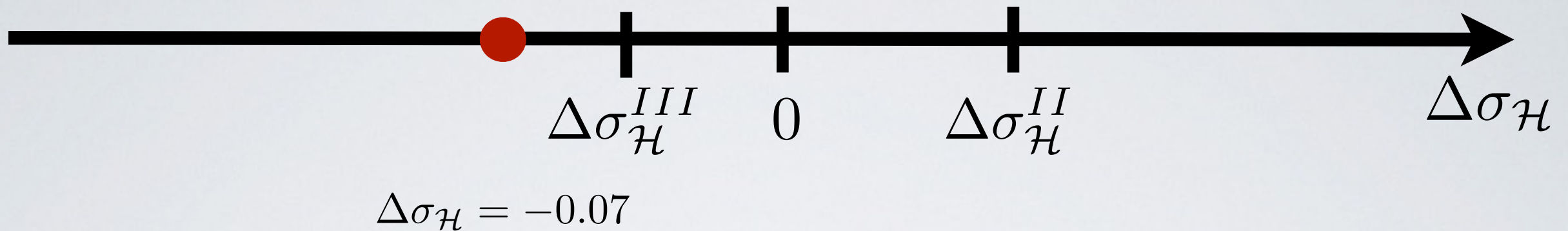
- Region II:

$$\Theta_+ < 0 \quad \Theta_- < 0$$

Black Hole

$$(j = 0.5)$$

Trapped Regions



$(j = 0.5)$

- Numerical Boundaries

$$\mathcal{J}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- New surface $\mathcal{S}_-$

$$\Theta_- = 0$$

- Region I:

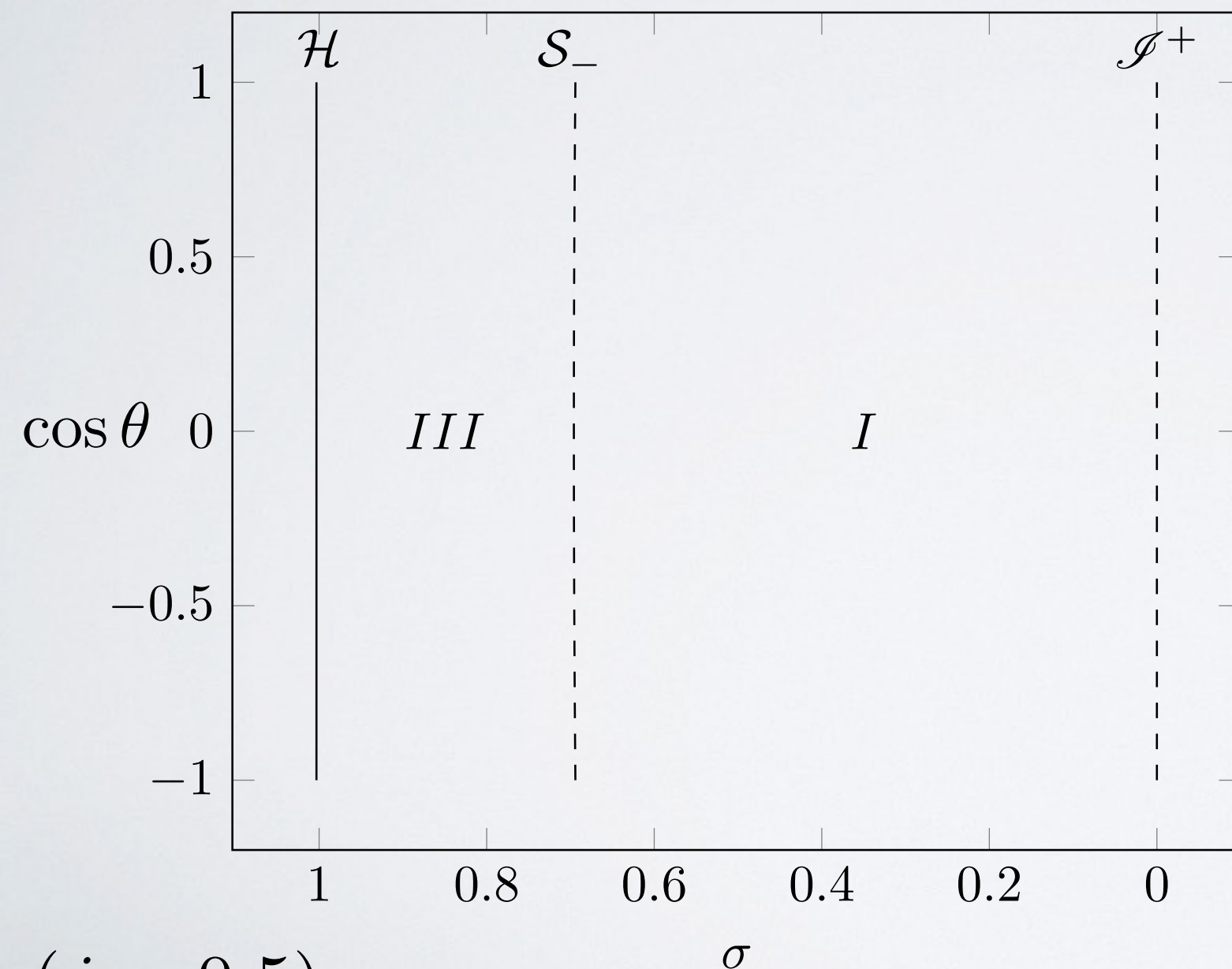
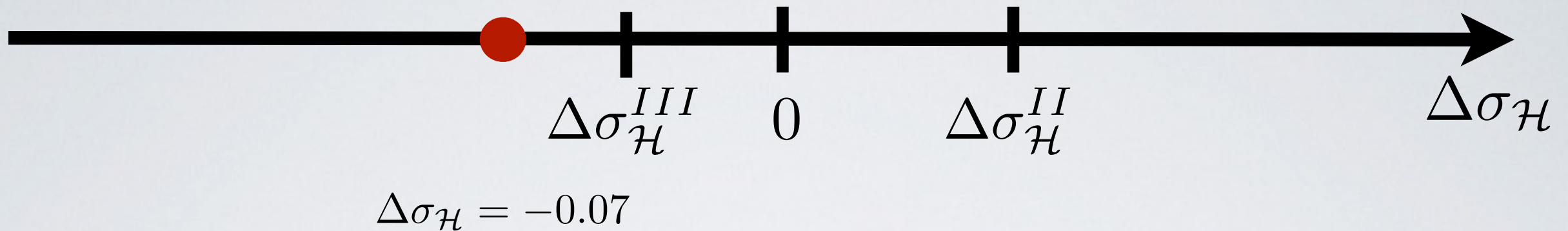
$$\Theta_+ > 0 \quad \Theta_- < 0$$

Exterior solution

- Region III:

$$\Theta_+ > 0 \quad \Theta_- > 0$$

Trapped Regions



$(j = 0.5)$

- Numerical Boundaries

$$\mathcal{J}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- New surface $\mathcal{S}_-$

$$\Theta_- = 0$$

- Region I:

$$\Theta_+ > 0 \quad \Theta_- < 0$$

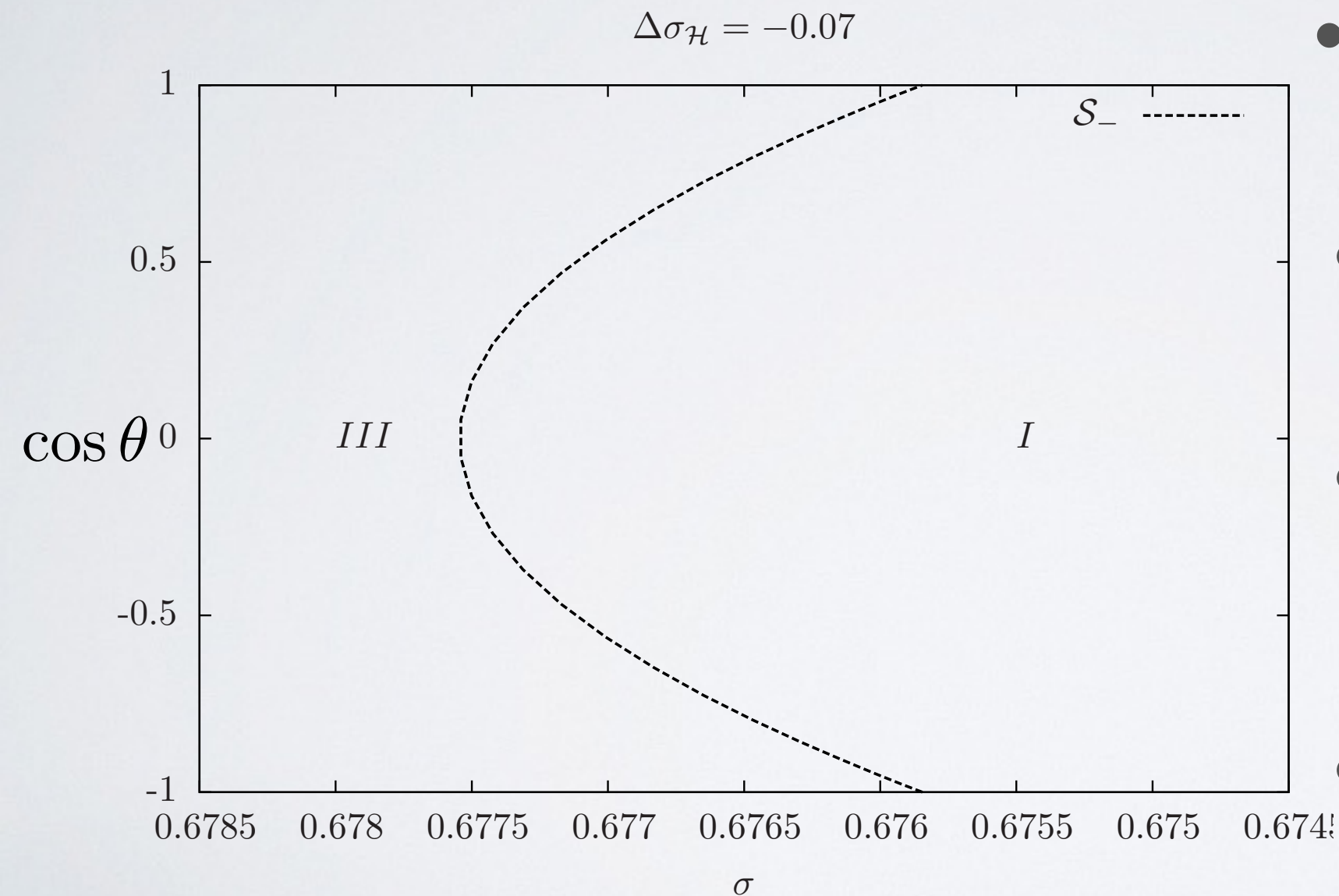
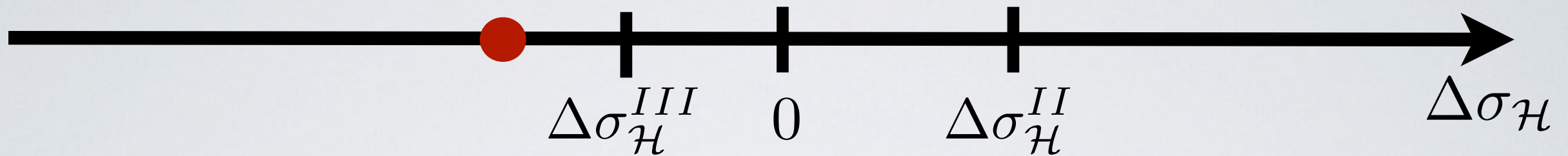
Exterior solution

- Region III:

$$\Theta_+ > 0 \quad \Theta_- > 0$$

White Hole

Trapped Regions



- Numerical Boundaries

$$\mathcal{I}^+ \quad \sigma = \sigma_{\mathcal{H}}$$

- New surface $\mathcal{S}_-$

$$\Theta_- = 0$$

- Region I:

$$\Theta_+ > 0 \quad \Theta_- < 0$$

Exterior solution

- Region III:

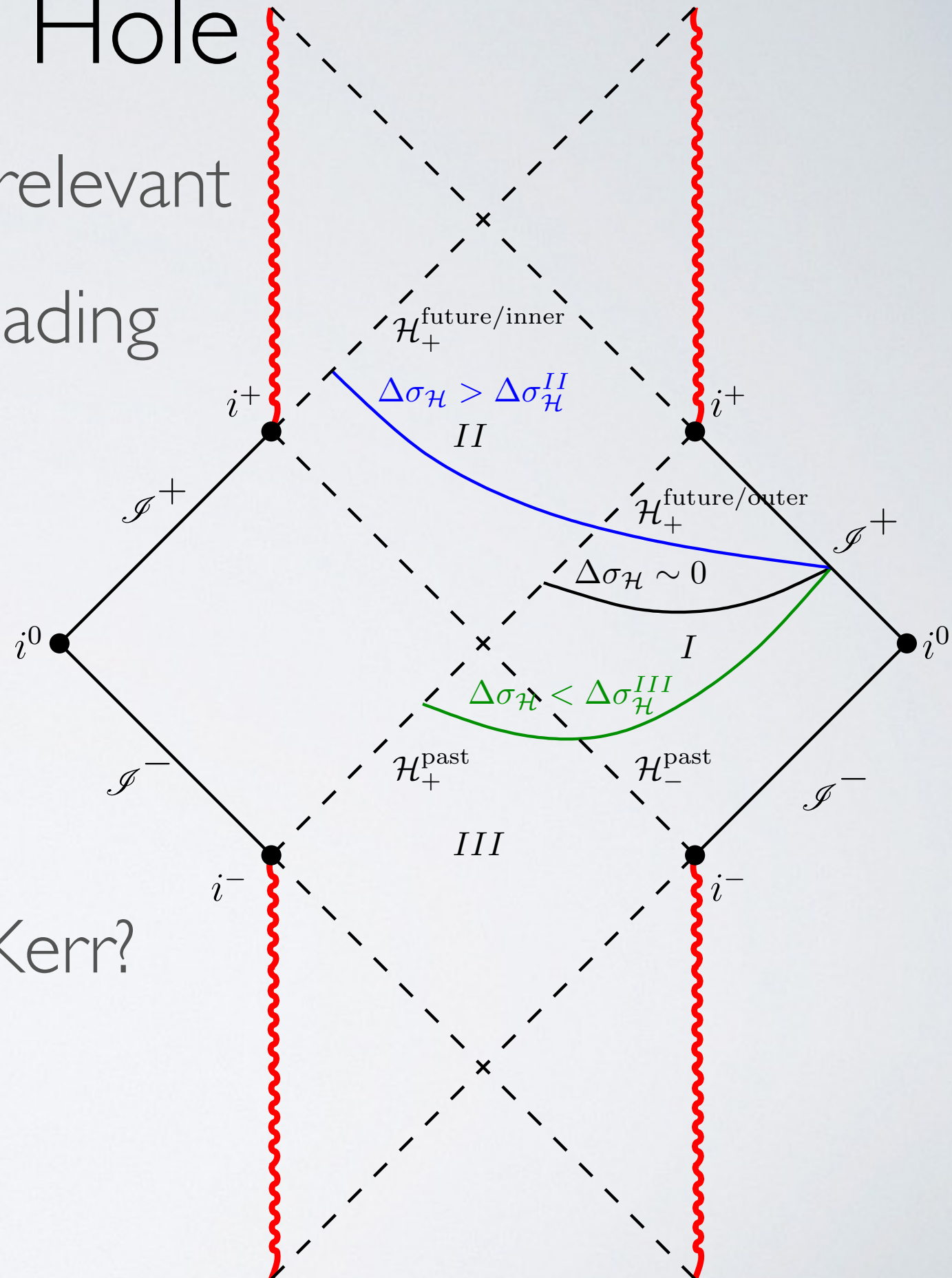
$$\Theta_+ > 0 \quad \Theta_- > 0$$

White Hole

$$(j = 0.5)$$

White Hole

- Astrophysical: solution not relevant
 - No scenario in nature leading to white holes
- Mathematic perspective:
 - Any role in stability of Kerr?
 - Bifurcation point?



Bondi Mass and Angular Momentum

- Bondi mass: Hawking mass evaluated at $\mathcal{I}^+$

$$M_B = \lim_{S \rightarrow \mathcal{I}^+} \left\{ \sqrt{\frac{\mathcal{A}_S}{16\pi}} \left[1 + \frac{1}{8\pi} \oint_S \Theta_+ \Theta_- dA \right] \right\},$$

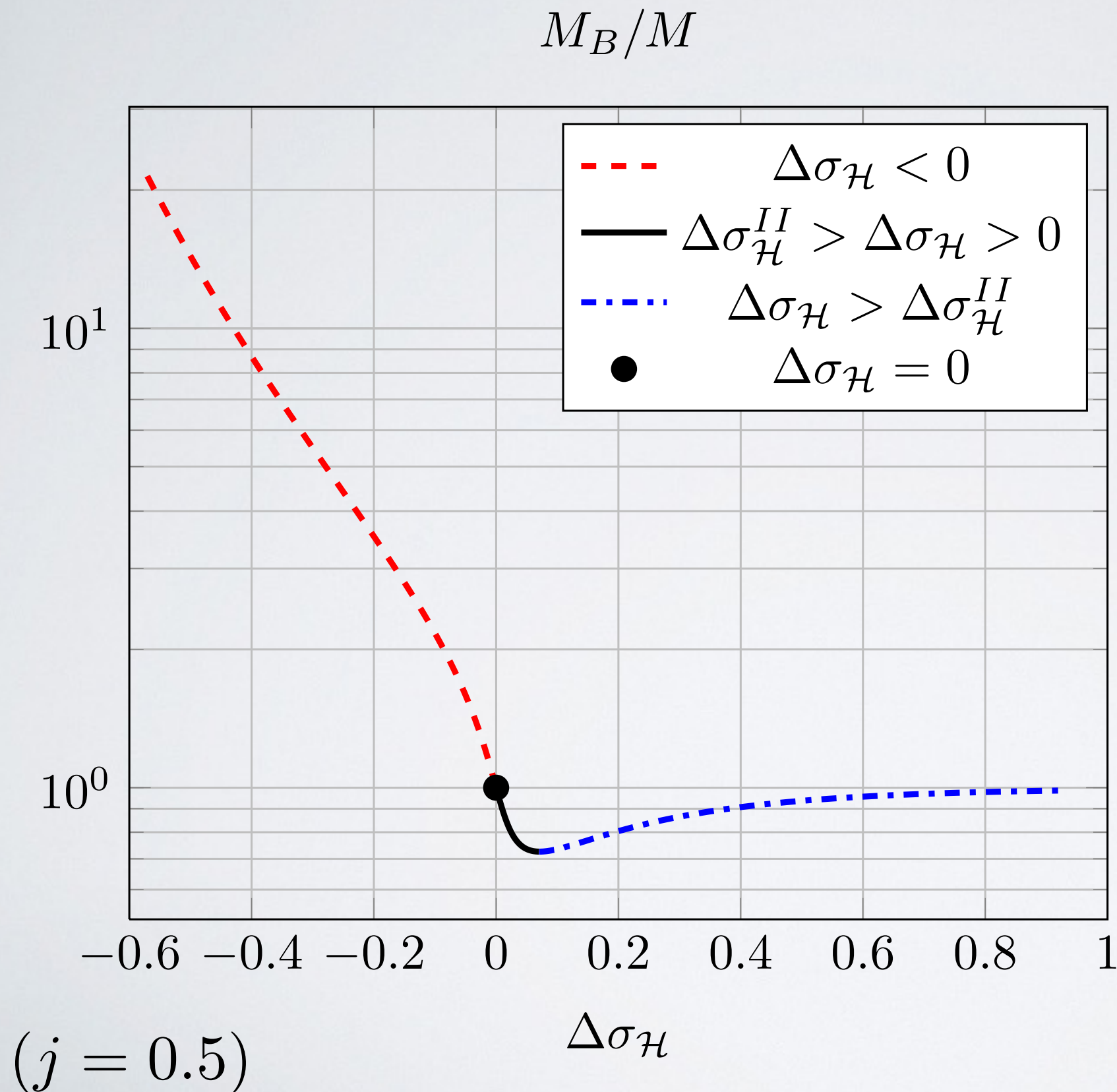
- Angular momentum: Komar angular momentum
(or Asthekar's first angular multipole)

$$J = \frac{1}{8\pi} \oint_{\mathcal{H}} K_{ij} \phi^i s^j dA$$

- Dimensionless angular momentum:

$$\dot{j}_B = \frac{J}{M_B^2}$$

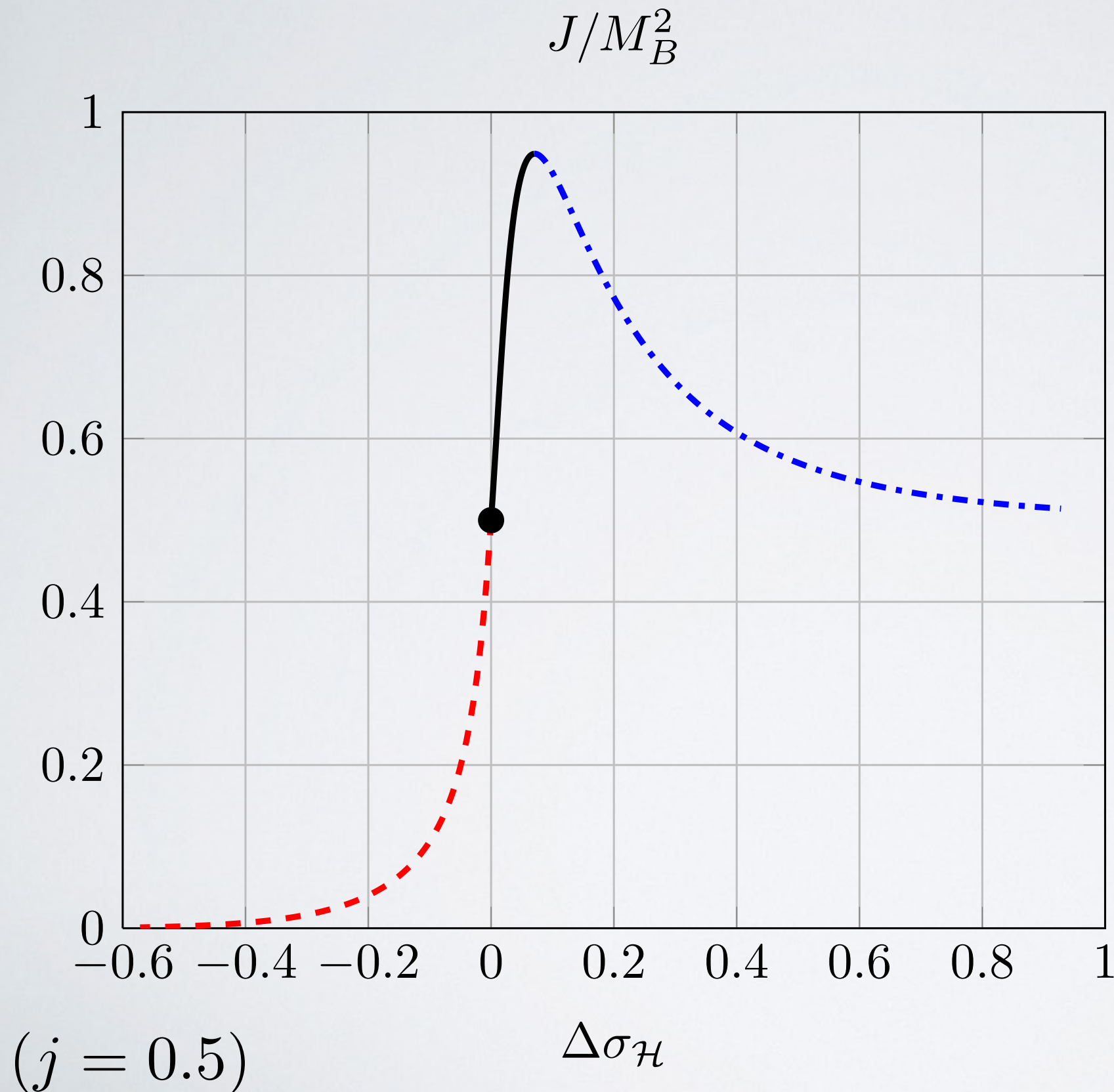
Bondi Mass and Angular Momentum



- For $\Delta\sigma_{\mathcal{H}} \gg \Delta\sigma_{\mathcal{H}}^{\text{II}}$

$$M_B \rightarrow M$$

Bondi Mass and Angular Momentum



- Maximum:

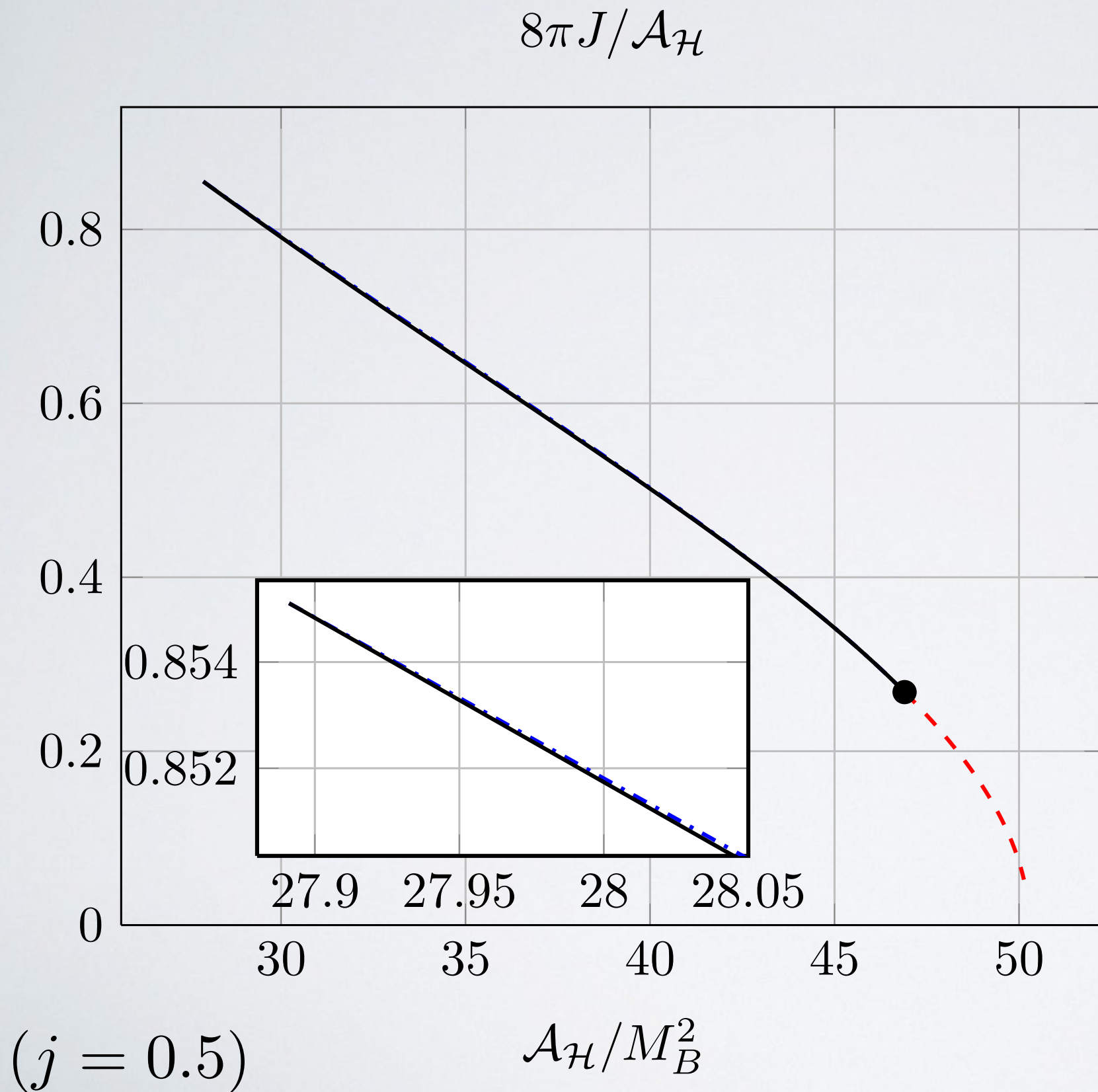
$$\dot{j}_B \approx 0.95$$

- For $\Delta\sigma_{\mathcal{H}} \gg \Delta\sigma_{\mathcal{H}}^{II}$

$$M_B \rightarrow M$$

$$\dot{j}_B \rightarrow 0.5$$

Bondi Mass and Angular Momentum



- Maximum:

$$\dot{j}_B \approx 0.95$$

- For $\Delta\sigma_{\mathcal{H}} \gg \Delta\sigma_{\mathcal{H}}^{II}$

$$M_B \rightarrow M$$

$$\dot{j}_B \rightarrow 0.5$$

Black hole inequalities

- Assuming weak cosmic censorship to hold, Penrose conjectured

$$A_{\mathcal{H}} \leq 16\pi M_{\text{ADM}}^2.$$

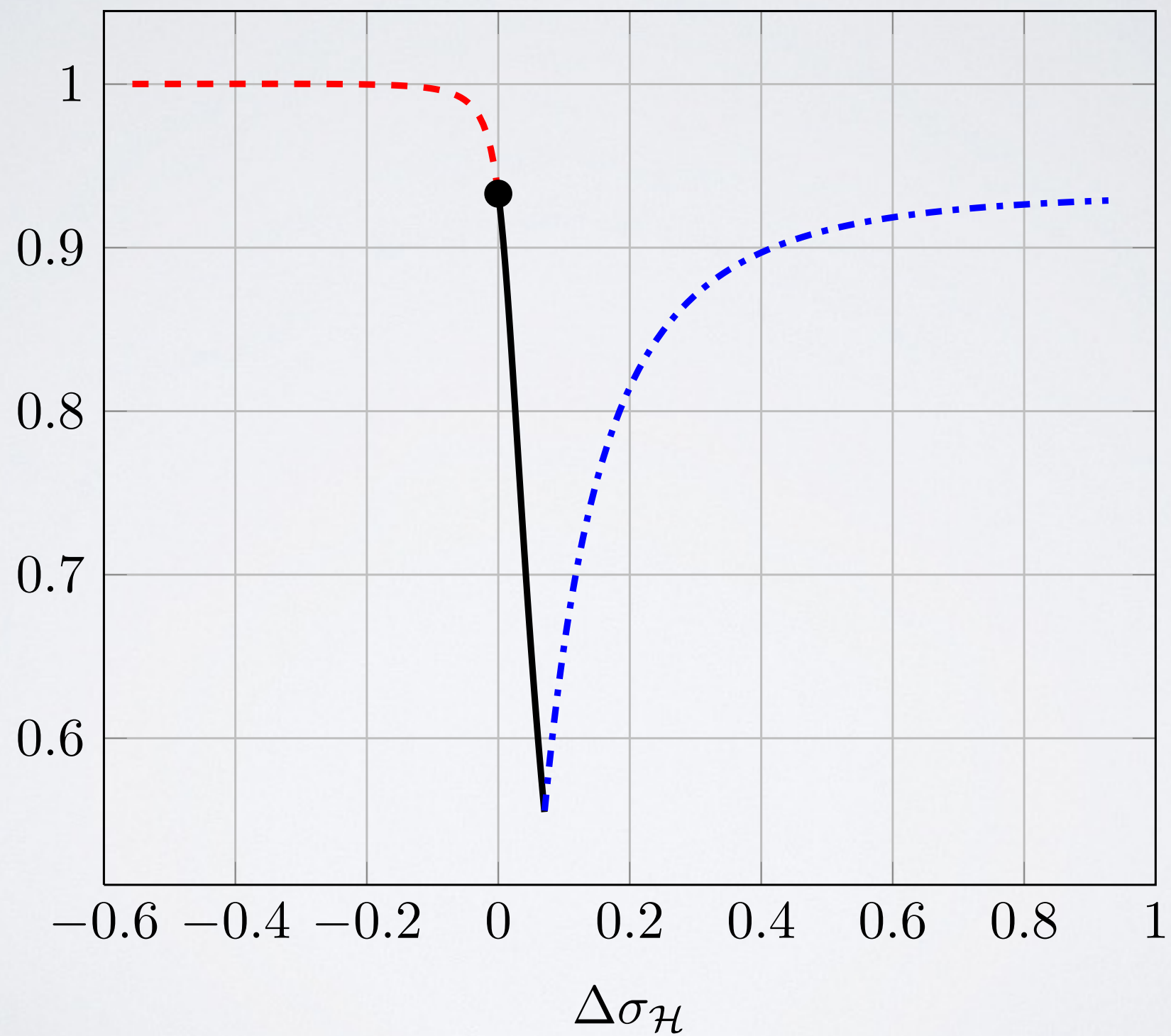
- Reasoning to the inequality passes through an intermediate stage

$$A_{\mathcal{H}} \leq 16\pi M_{\text{B}}^2,$$

- Stronger inequality since $M_{\text{B}} \leq M_{\text{ADM}}$
- Rigidity: equality holds only for Schwarzschild

Black hole inequalities

$$\mathcal{A}_{\mathcal{H}}/(16\pi M_B^2)$$



$(j = 0.5)$

Black hole inequalities

- For Kerr, Dain *et. al.* ('02) introduced a refined inequality

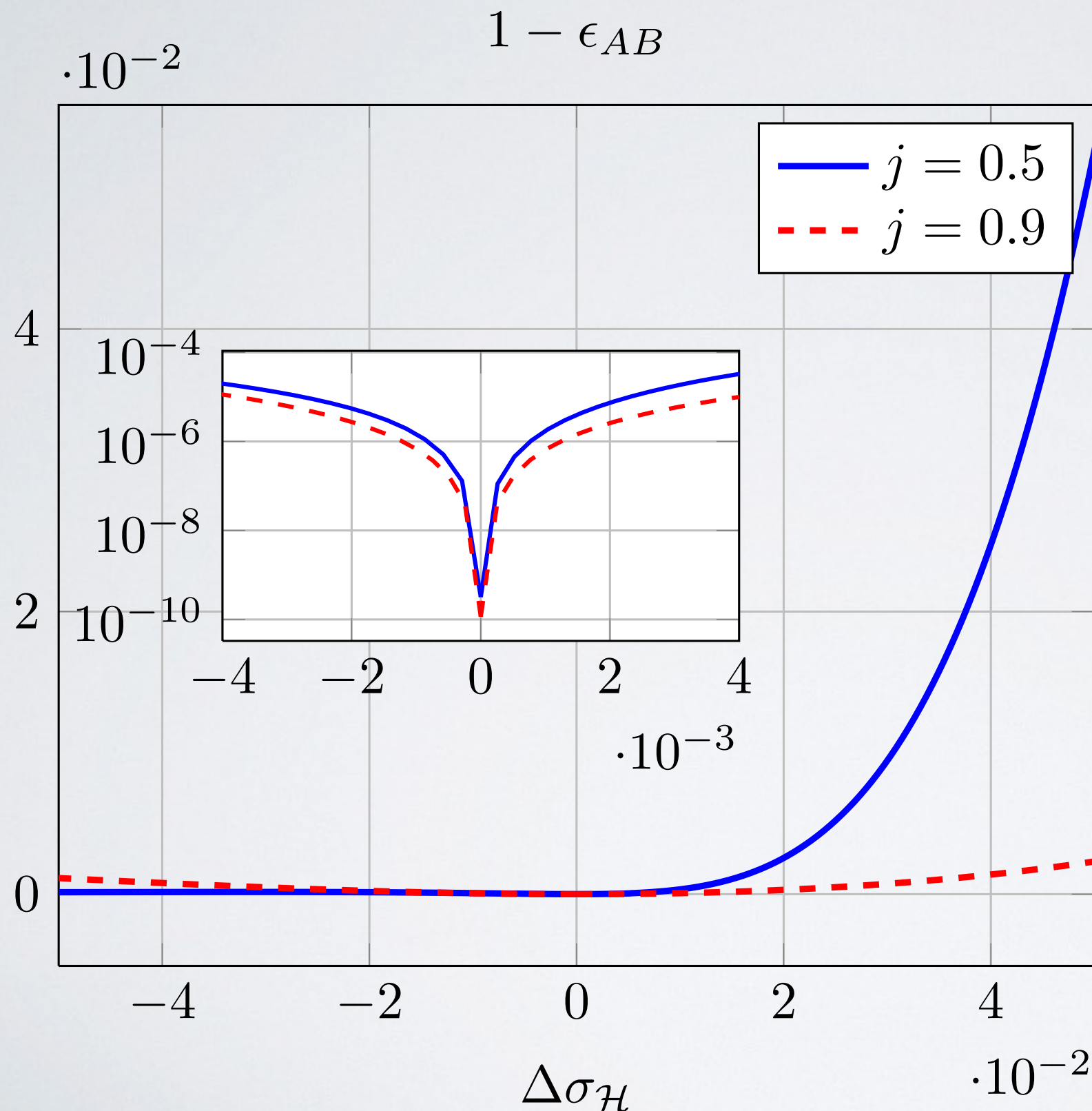
$$\epsilon_A := \frac{A_{\mathcal{H}}}{8\pi(M_{\text{ADM}}^2 + \sqrt{M_{\text{ADM}}^4 - J^2})} \leq 1,$$

- Following the arguments for the Penrose inequality, we (thanks to J.L. Jaramillo) propose a stronger relation in terms of a **Dain-Bondi number**

$$\epsilon_{AB} = \frac{A_{\mathcal{H}}}{8\pi(M_{\text{B}}^2 + \sqrt{M_{\text{B}}^4 - J^2})} \leq 1$$

- Rigidity: equality holds only for Kerr

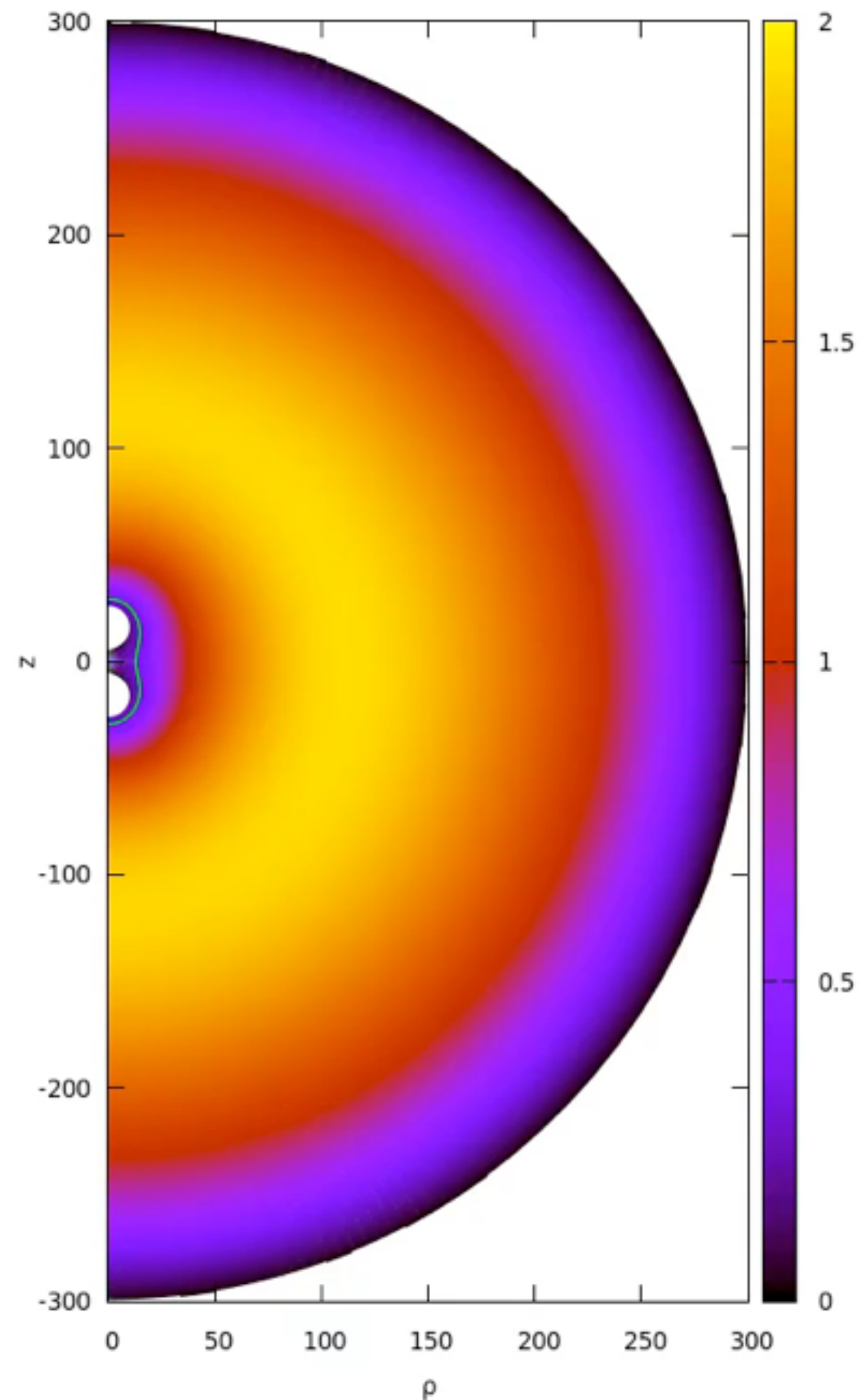
Black hole inequalities



- Results confirms the inequality (within numerical precision)
- Spectral methods are essential to attest the validity of such inequalities

Future Work

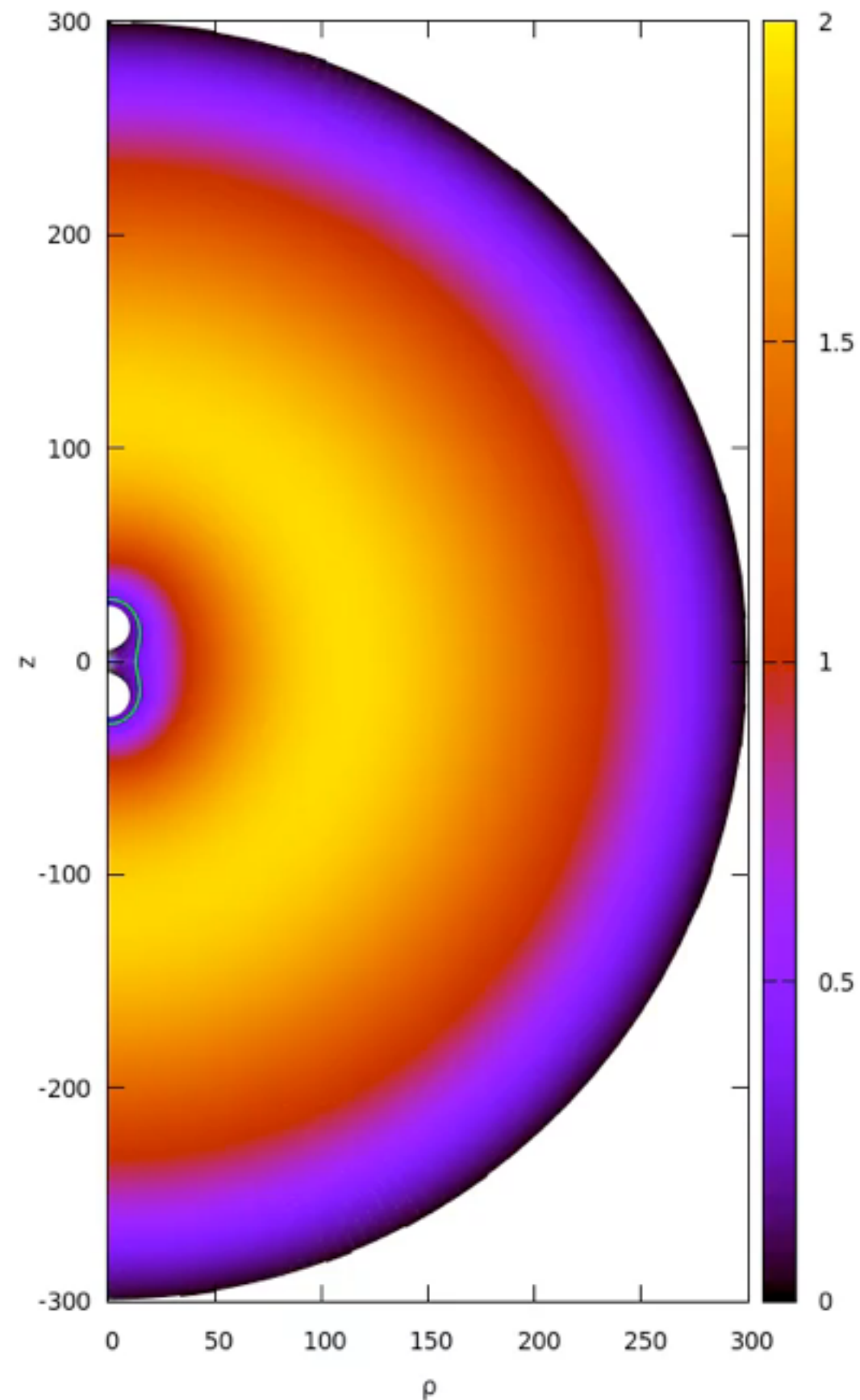
- Initial data for axis-symmetric binary black holes (head-on) in hyperboloidal slices



Credit: David Schinkel

Future Work

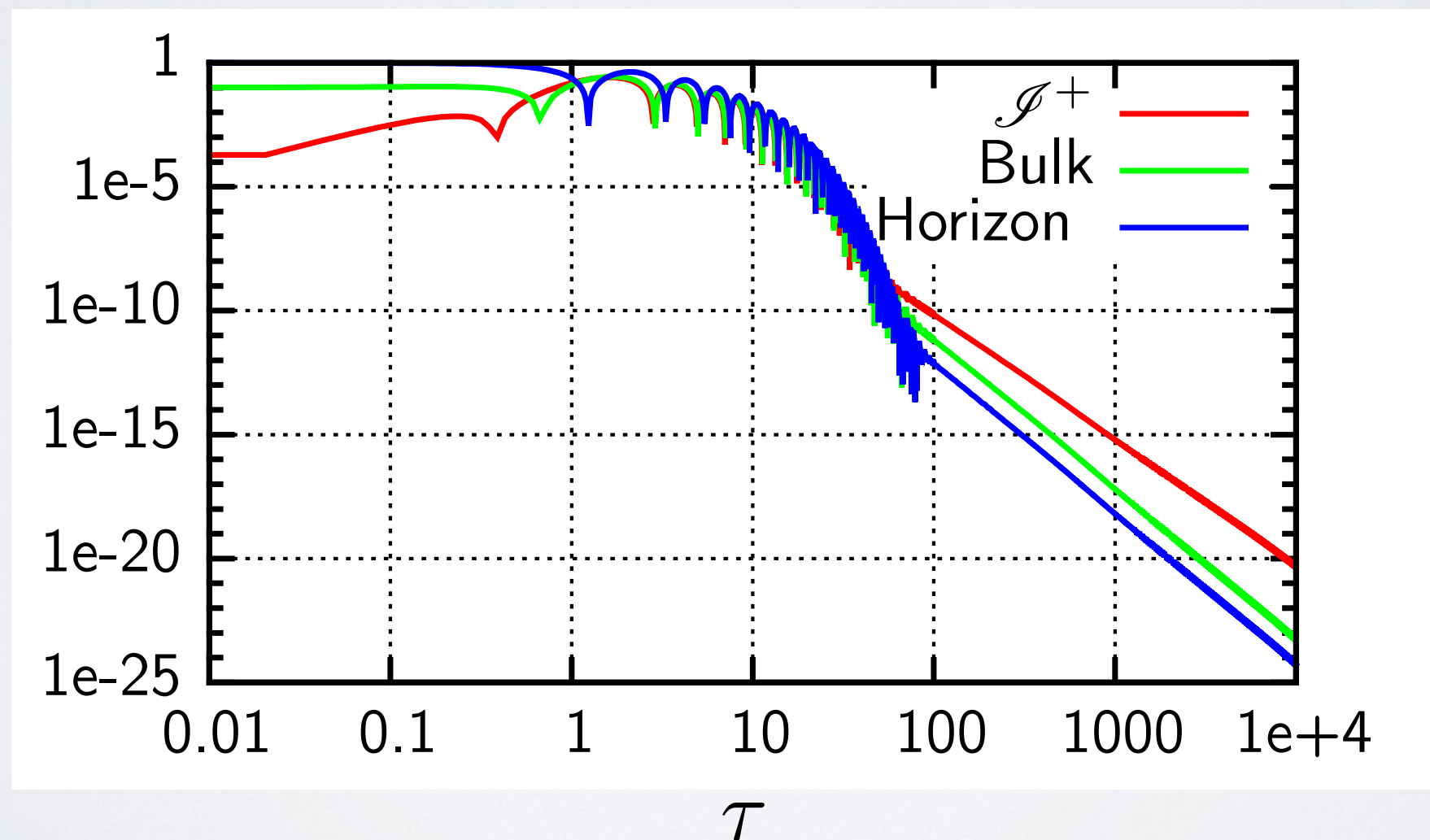
- Initial data for axis-symmetric binary black holes (head-on) in hyperboloidal slices



Credit: David Schinkel

Future Work

- Time evolution of Einstein's equations
- Fully-spectral Numerical Code: J. Comp. Phys **276** 357 (2014)



Linear Regime (Teukolsky Equation)

Conclusion

- Initial data for perturbed rotating black hole on hyperboloidal slices
- Trapped surfaces showed a rich structure for strong perturbations
 - New black and white hole regions
- Fast rotating perturbed black holes $j_B \approx 0.95$
free parameters ($j = 0.5$)
- Dain-Bondi number as a relevant quantity in the standard picture of gravitational collapse

$$\epsilon_{AB} = \frac{A_{\mathcal{H}}}{8\pi(M_B^2 + \sqrt{M_B^4 - J^2})}$$

Mass and angular multipoles

- Does the solution describes a perturbed black hole or is it just Kerr in a different slice (with new parameters)?

Mass and angular multipoles

- Does the solution describes a perturbed black hole or is it just Kerr in a different slice (with new parameters)?
- Ashtekar's mass and angular multipoles:

$$M_n = M_{\mathcal{H}} \frac{R_{\mathcal{H}}^n}{8\pi} \oint_{\mathcal{H}} {}^{(2)}R(\mu) P_n(\hat{\mu}(\mu)) dA$$

$$J_n = \frac{R_{\mathcal{H}}^{n-1}}{8\pi} \oint_{\mathcal{H}} K_{\Phi_S}(\mu) P'_n(\hat{\mu}(\mu)) dA$$

Mass and angular multipoles

- Does the solution describes a perturbed black hole or is it just Kerr in a different slice (with new parameters)?
- Ashtekar's mass and angular multipoles:

$$M_n = M_{\mathcal{H}} \frac{R_{\mathcal{H}}^n}{8\pi} \oint_{\mathcal{H}} {}^{(2)}R(\mu) P_n(\hat{\mu}(\mu)) dA$$

$$J_n = \frac{R_{\mathcal{H}}^{n-1}}{8\pi} \oint_{\mathcal{H}} K_{\Phi_S}(\mu) P'_n(\hat{\mu}(\mu)) dA$$

- Compute M_0 and J_1 from the solution

Mass and angular multipoles

- Does the solution describes a perturbed black hole or is it just Kerr in a different slice (with new parameters)?
- Ashtekar's mass and angular multipoles:

$$M_n = M_{\mathcal{H}} \frac{R_{\mathcal{H}}^n}{8\pi} \oint_{\mathcal{H}} {}^{(2)}R(\mu) P_n(\hat{\mu}(\mu)) dA$$

$$J_n = \frac{R_{\mathcal{H}}^{n-1}}{8\pi} \oint_{\mathcal{H}} K_{\Phi_S}(\mu) P'_n(\hat{\mu}(\mu)) dA$$

- Compute M_0 and J_1 from the solution
- Assume Kerr space-time with $M = M_0$ and $a = \frac{J_1}{M_0}$

Mass and angular multipoles

- Does the solution describes a perturbed black hole or is it just Kerr in a different slice (with new parameters)?
- Ashtekar's mass and angular multipoles:

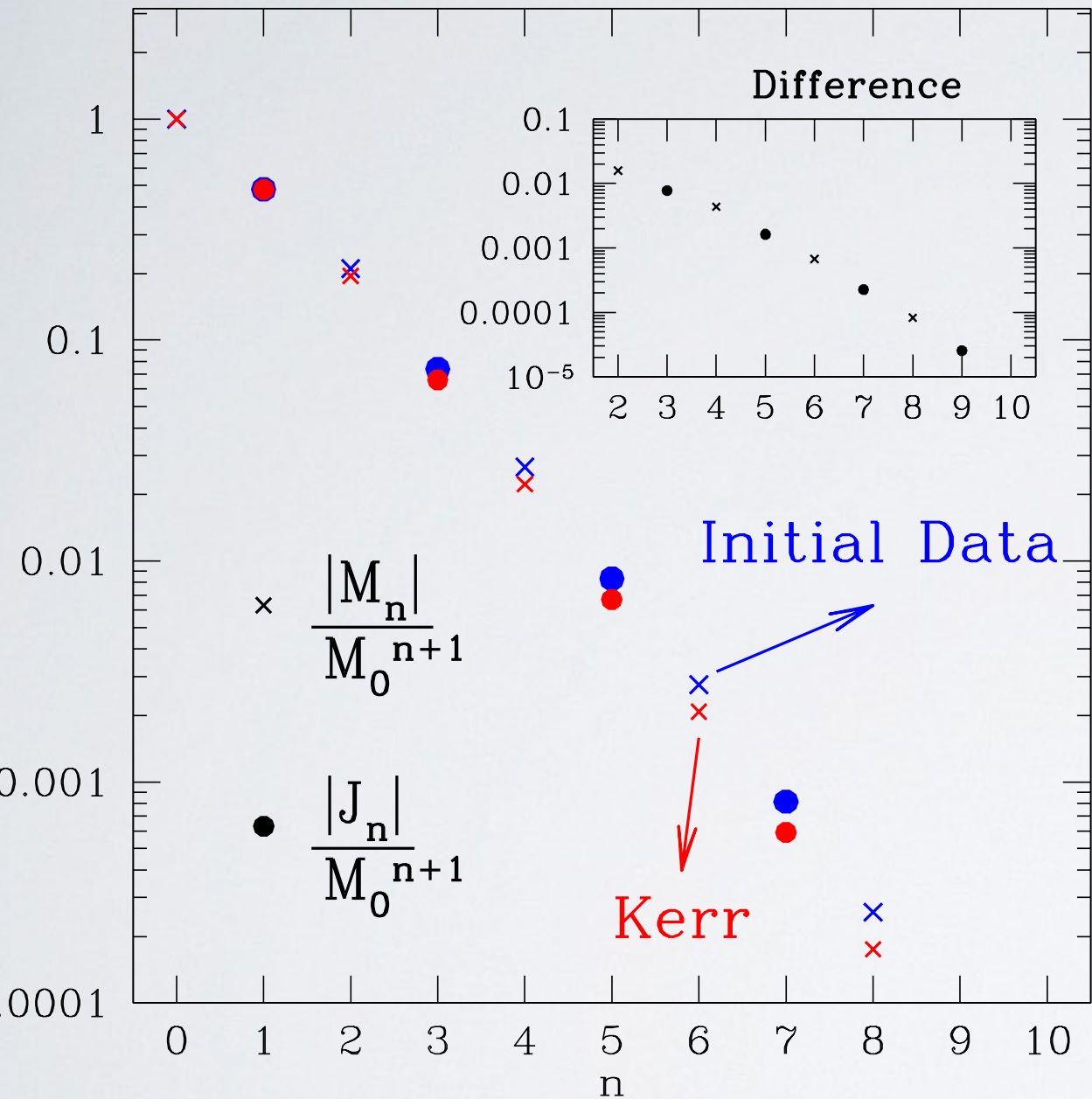
$$M_n = M_{\mathcal{H}} \frac{R_{\mathcal{H}}^n}{8\pi} \oint_{\mathcal{H}} {}^{(2)}R(\mu) P_n(\hat{\mu}(\mu)) dA$$

$$J_n = \frac{R_{\mathcal{H}}^{n-1}}{8\pi} \oint_{\mathcal{H}} K_{\Phi_S}(\mu) P'_n(\hat{\mu}(\mu)) dA$$

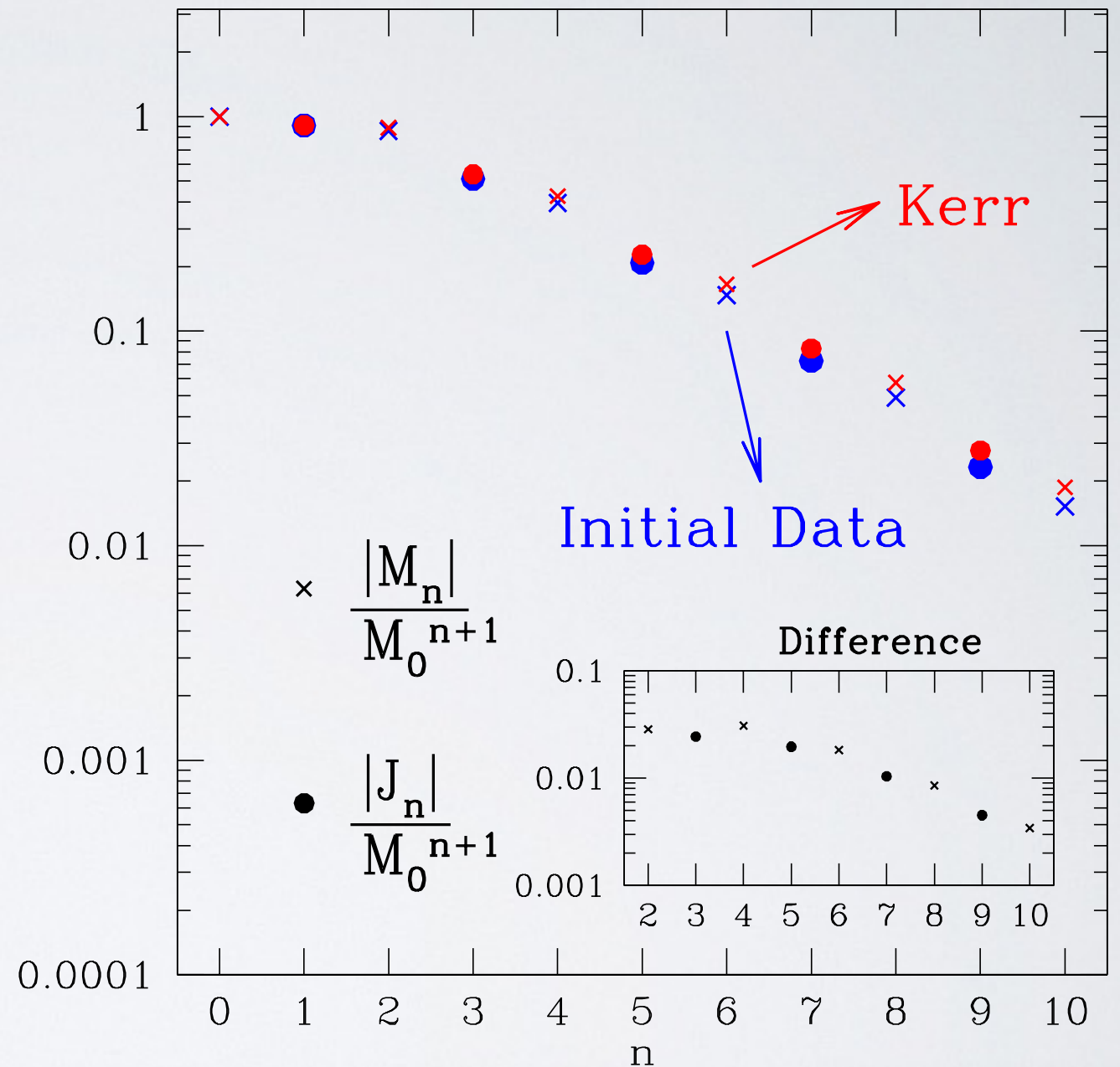
- Compute M_0 and J_1 from the solution
- Assume Kerr space-time with $M = M_0$ and $a = \frac{J_1}{M_0}$
- Compare higher multipoles

Mass and angular multipoles

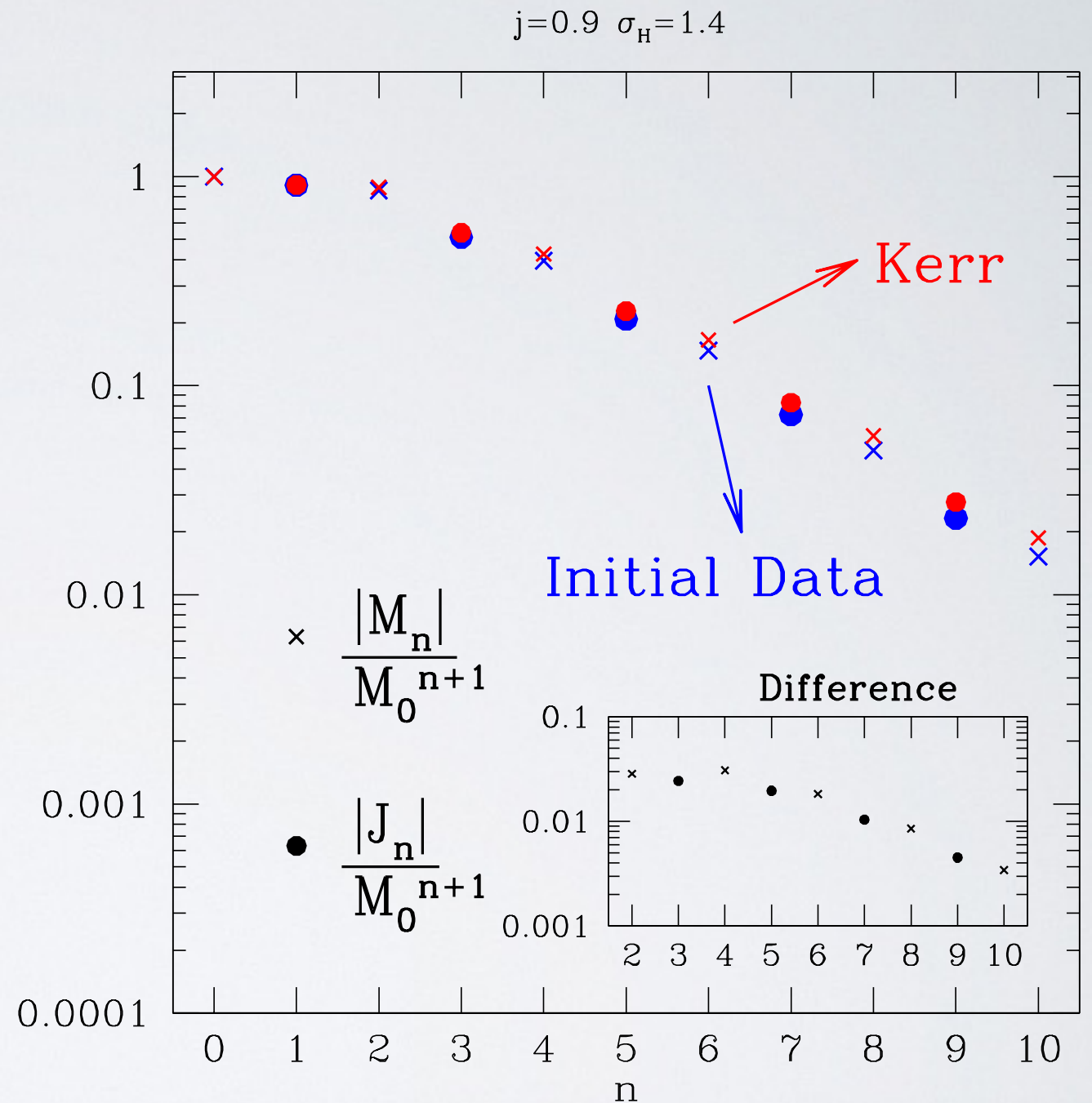
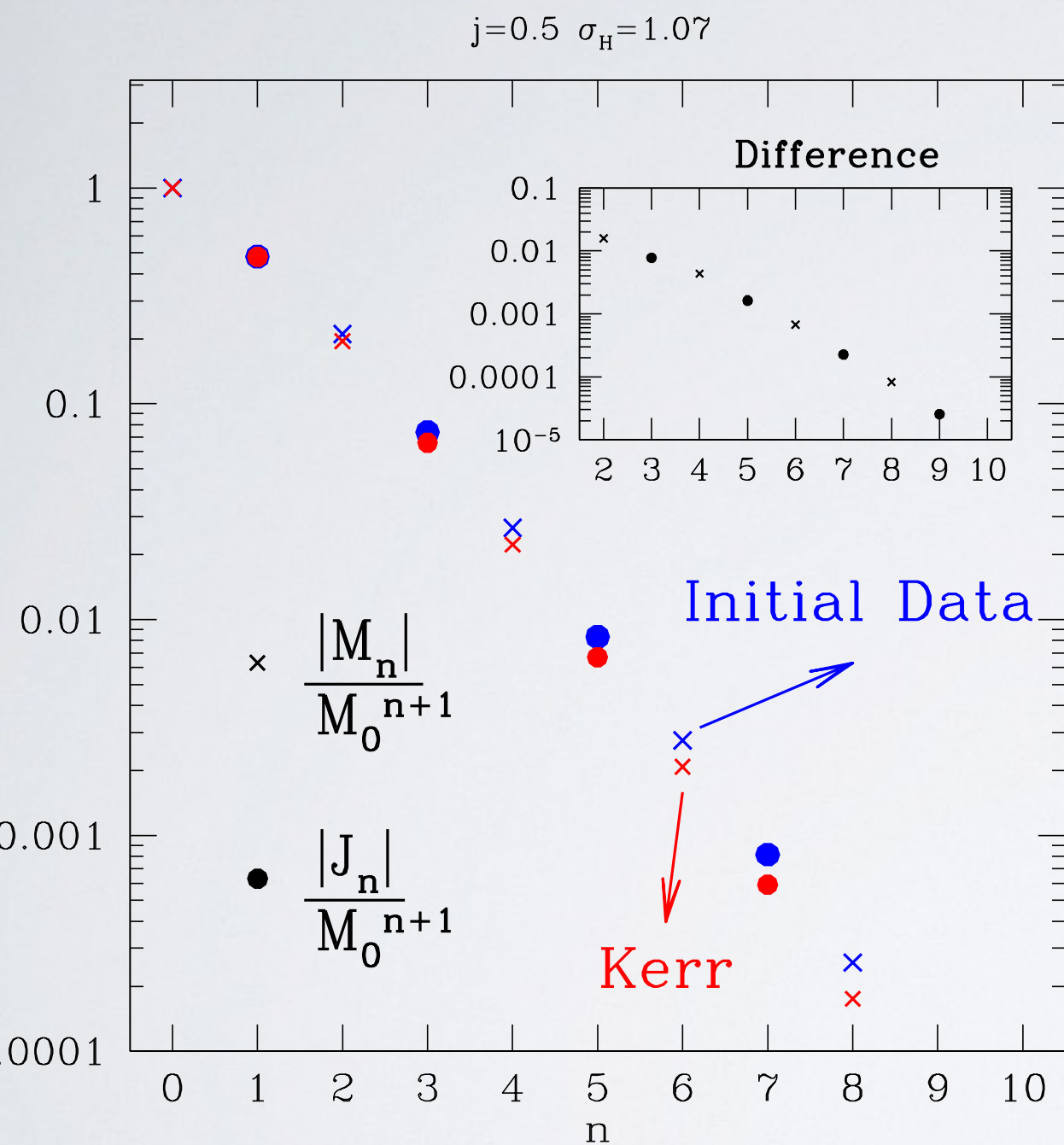
$j=0.5 \quad \sigma_H=1.07$



$j=0.9 \quad \sigma_H=1.4$

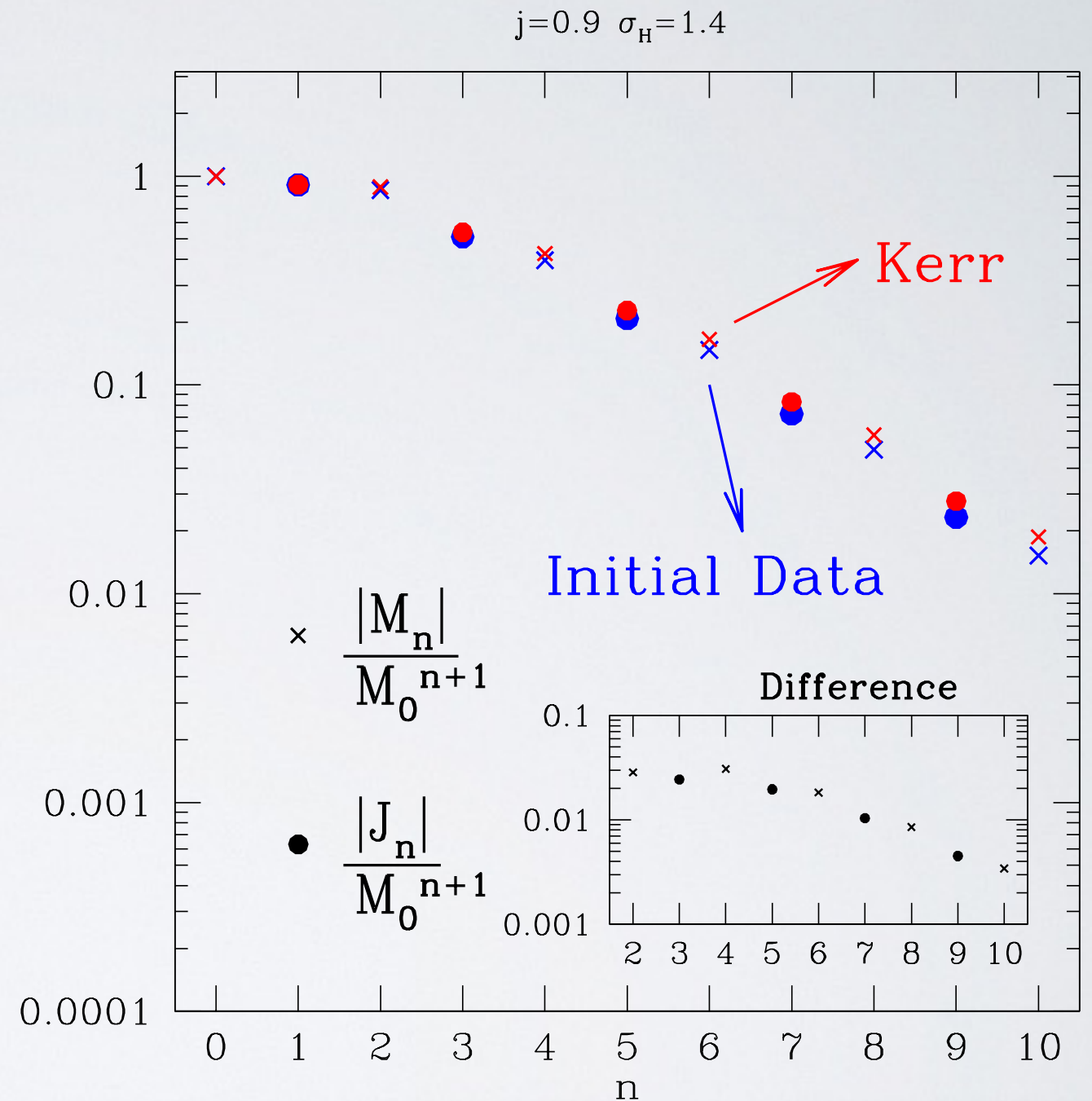
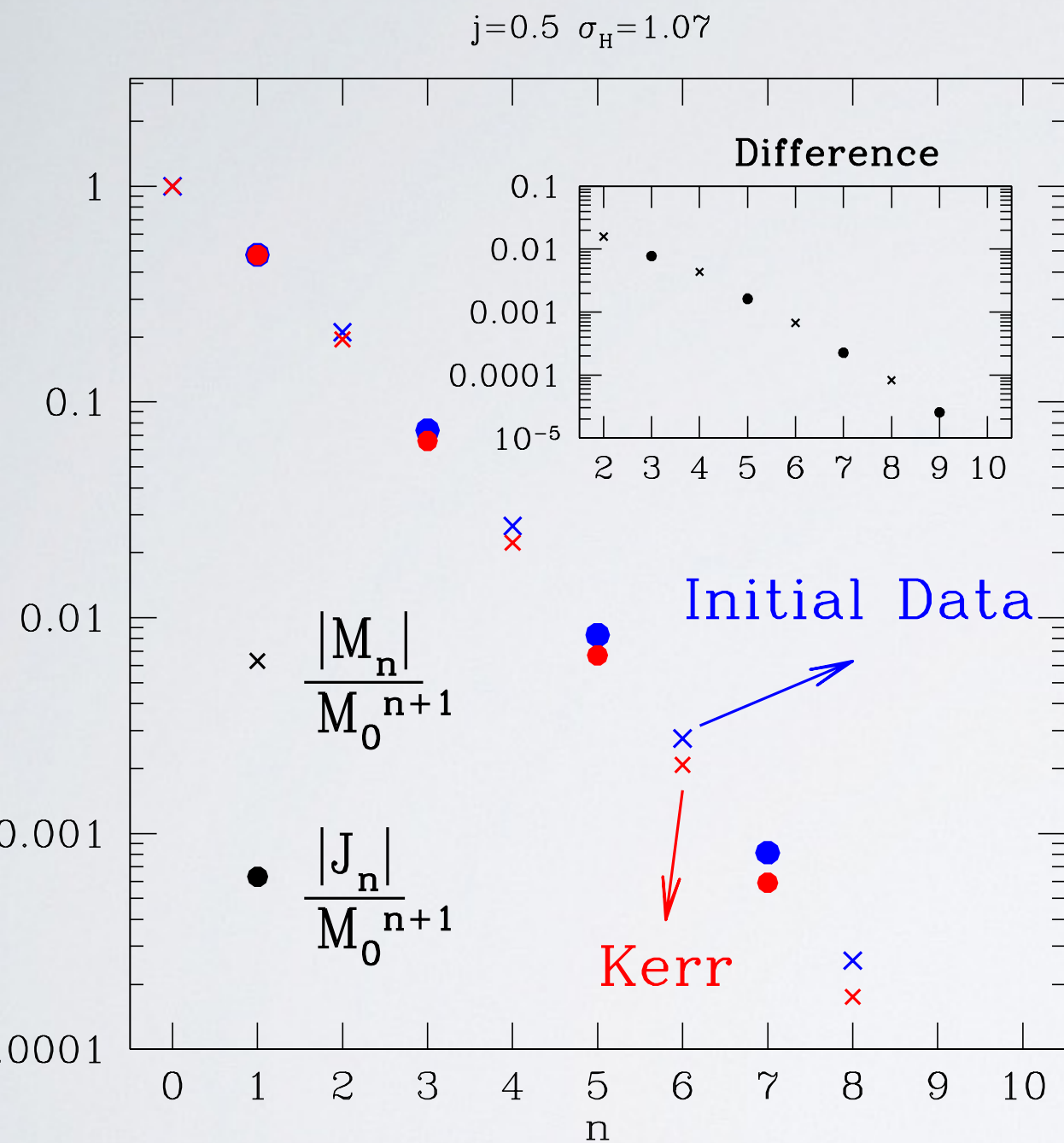


Mass and angular multipoles



- Small perturbation regime (no new trapped surfaces)

Mass and angular multipoles



- Small perturbation regime (no new trapped surfaces)
- Solution is a perturbed rotating black hole

Black hole inequalities

- Rigidity:

- For $\Delta\sigma_{\mathcal{H}} \gg \Delta\sigma_{\mathcal{H}}^{II}$

“Strong perturbation”
with the same
parameters

$$M_B \rightarrow M$$

$$j_B \rightarrow j$$

$$\epsilon_{AB} \rightarrow 1$$

- Similar behaviour also in J.L. Jaramillo, N.Vasset and M. Ansorg ('08) : in terms of ADM mass

- Is it still a perturbed black hole or is it Kerr?

- Does $\epsilon_{AB} = 1 \Rightarrow \text{Kerr}$?

- Possible way to answer this question is to compute the multipoles for “strong perturbations” (future work)