

Quantum black holes and their lepton signatures at the LHC

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Credit: NASA/JPL-Caltech

**Black Holes and their Analogues: 100 years of General Relativity
in honor of Kirill Bronnikov's 70th birthday**

Ubu, Brazil, April 16th, 2015

OUTLINE

Based on arXiv:1412.2661, AB and Xavier Calmet

- Introduction -
QBH as a link between GR and QFT
- QBH at the LHC
 - Our approach
 - Model implementation
 - LHC potential to probe QBH in XD models
- Conclusions

Planckian (Quantum) Black Holes: QBH

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 - In three-dimensional gravity, the minimum mass of microscopic black hole is $M_{\text{PL}} = 10^{19} \text{ GeV}$

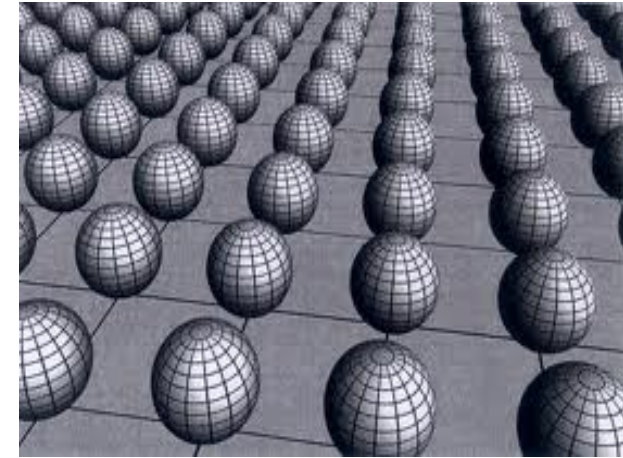
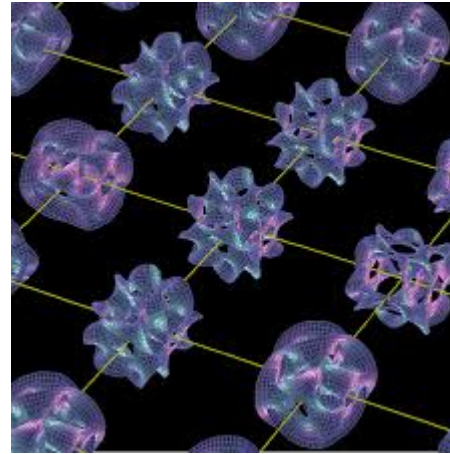
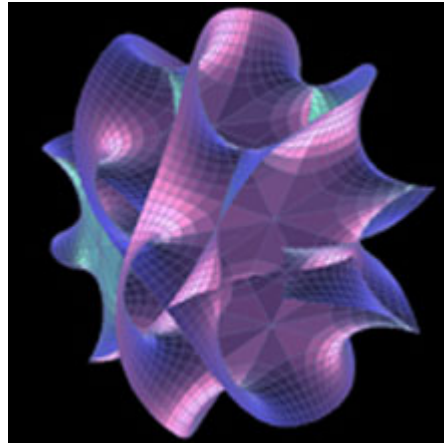
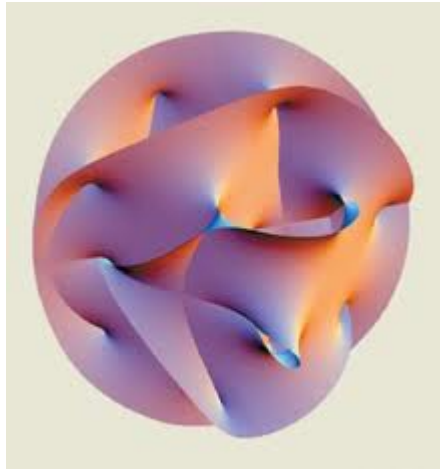
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 - To reach Planck length and Planck Mass with current technologies one should build collider with **about 1000 light years** in diameter!
- Luckily, in some scenarios involving (large) extra dimensions of space, the Planck mass can be as low as the TeV range
In this case QBH will be accessible at the LHC - this talk

eXtra Dimensions (XD)



Motivations

- String theory, the best candidate to unify gravity & gauge interactions, **is only consistent in 10 D space-time**
- Extending symmetries: **Internal symmetries - GUTs, technicolour...; Fermionic spacetime- SUSY; Bosonic spacetime - Extra dimensions**
- The presence of XD could have an impact on scales $\ll M_{\text{planck}}$ (started with ADD)

The question is what is the size and the shape of XD ?!

New perspectives of XD

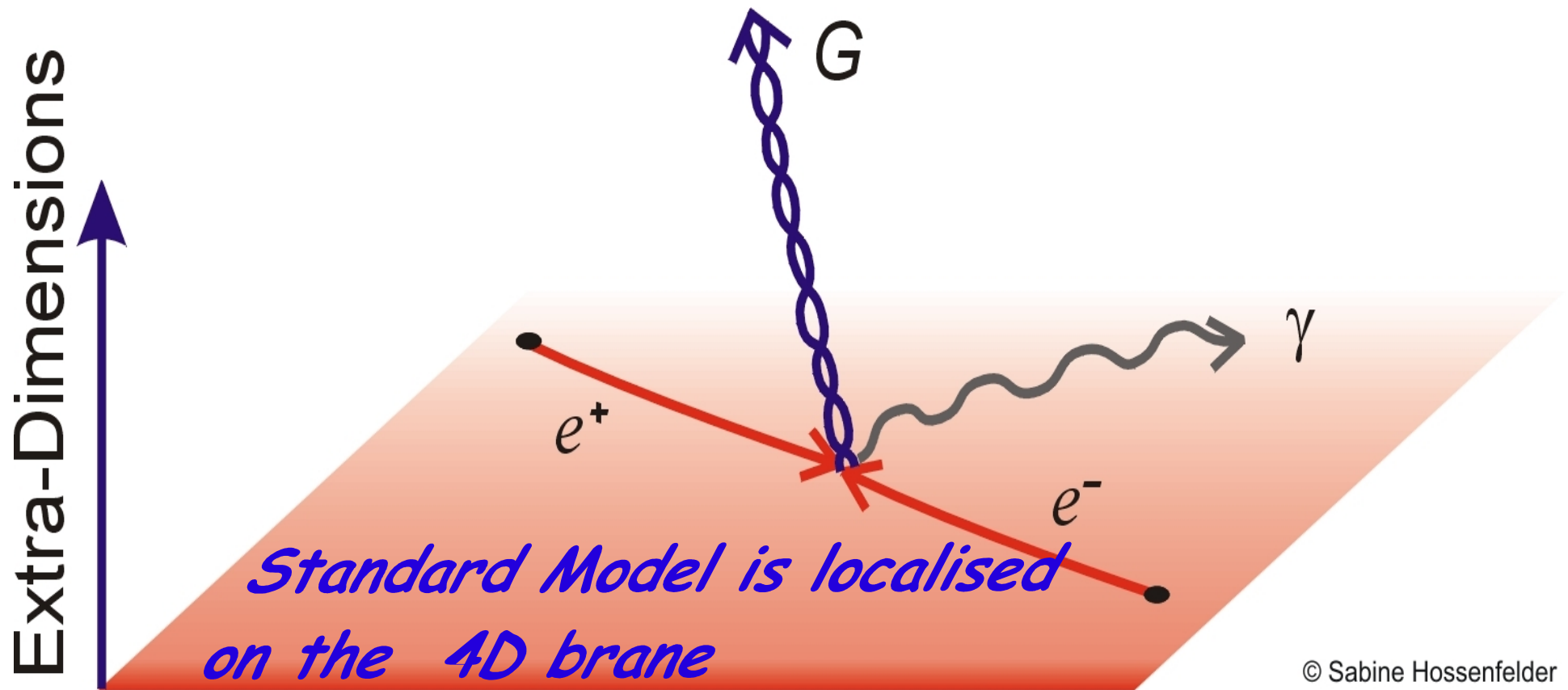
- The nature of electroweak symmetry breaking
- The origin of fermion mass hierarchies
- The supersymmetry breaking mechanism
- The description of strongly interacting sectors (provide a way to model them)
-

Brief History

- 1914: Nordstrom tried to unify gravity and electromagnetism in 5D
($A_\mu \rightarrow A_M$, where $M = 0,1,2,3,4$)
- 1920's: Kaluza and Klein tried using Einstein's equations in 5D ($g^{\mu\nu} \rightarrow g^{MN} \sim g^{\mu\nu}, g^{\mu 4}, g^{44}$)
- 1970's: Development of superstring theory and supergravity required extra dimensions
- 1998: Arkani-Hamed, Dimopoulos, and Dvali propose **Large Extra Dimensions** (ADD) as a solution to the Hierarchy /Fine tuning problem of the Standard Model

The idea of ADD

- The Standard Model has been tested to $r \sim 10^{-16}$ mm, Gravity has been tested to $r \sim 1$ mm only



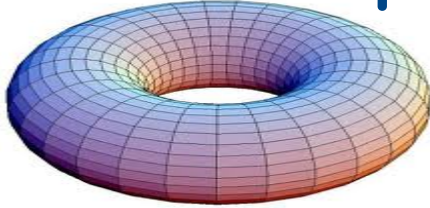
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- $4D \rightarrow (4 + n)D$

The effective $D = 4$ action is

$$\frac{M_f^{2+n}}{2} \int d^4x \int_0^{2\pi R} d^n Z \sqrt{G} R_{4+n} \longrightarrow \frac{1}{2} M_f^{2+n} V_n \int d^4x \sqrt{g} R$$

In case of toroidal compactification of equal radii, R

$$V_n = (2\pi R)^n$$

$$M_P^2 = M_f^{2+n} V_n$$

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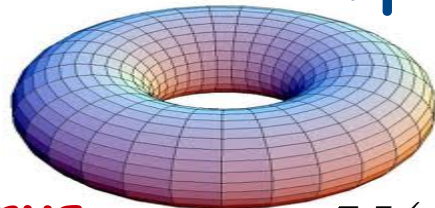
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effectively disappear**

$$V(r) = -G_N \frac{m_1 m_2}{r} = -\frac{m_1 m_2}{M_P^2 r}$$

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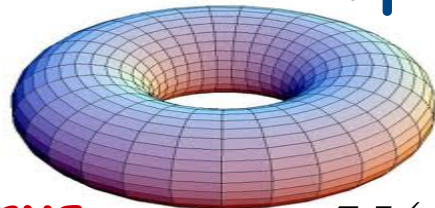
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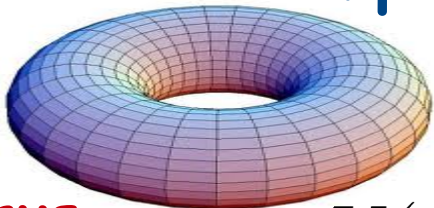
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**Fundamental quantum
gravity scale**



$$M_f^{2+n} r^{1+n}$$

The current status of ADD

So, $M_P^2 = M_f^{n+2} (2\pi R)^n$ and respectively,

$$R = \frac{1}{2\pi} \frac{1}{M_f} \left(\frac{M_P}{M_f} \right)^{\frac{2}{n}} [\text{GeV}^{-1}] \times 0.197 [\text{GeV m}]$$

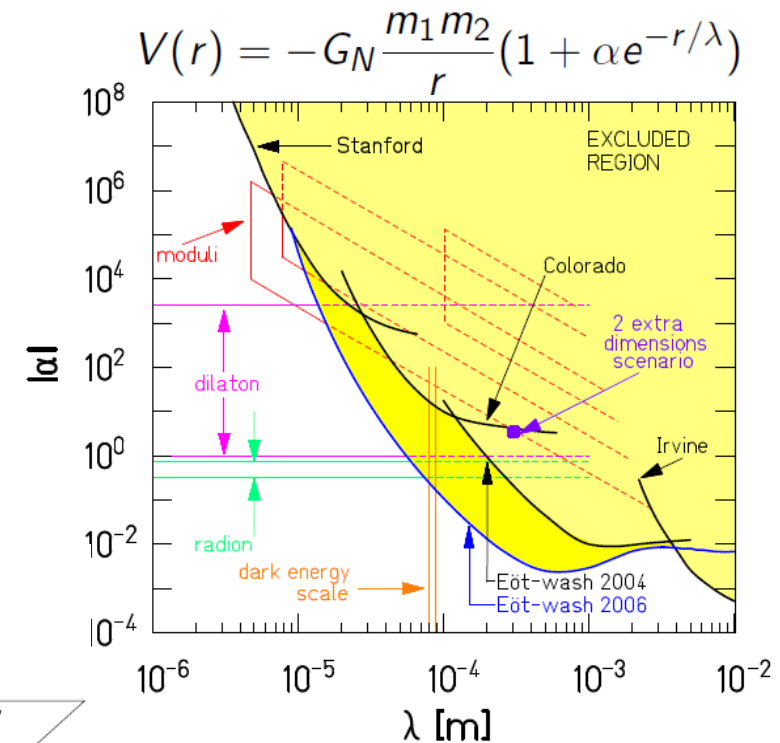
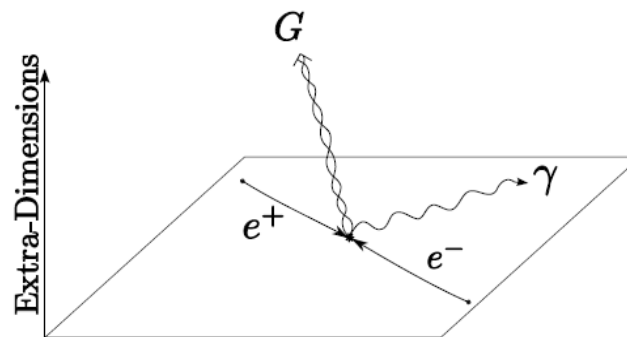
How big are these dimensions are?

Let us assume $M_f \sim 1 \text{ TeV}$, then

$$R \sim \begin{cases} 10^{15} \text{ mm} & n = 1 \quad \times \text{ Already} \\ 1 \text{ mm} & n = 2 \quad \times \text{ ruled out} \\ 10^{-6} \text{ mm} & n = 3 \\ \vdots \end{cases}$$

Collider signature:

$$pp \rightarrow \text{jet} + \cancel{E}_T$$



The current bound is $R < 37 \mu\text{m}$
For $n = 2$ this means that
 $M > 1.4 \text{ TeV}$

Summary of ADD ("flat scenario") [1998]

- gravity propagates in the compact & non-compact dimension
- the Einstein action

$$S_E = \frac{1}{8\pi G} \int d^D x \sqrt{-g} \frac{1}{2} \mathcal{R}$$

- implies the relation between the four-dimensional and D-dimensional Newton's constants

$$G_N = \frac{G_D}{V_{D-4}}$$

- The reduced fundamental Planck scale is generally related to the D-dimensional Newton's constant as

$$M_p^{D-2} = \frac{(2\pi)^{D-4}}{4\pi G_D}$$

Rundal Sundrum ("warp scenario") [1999]

- a warped geometry is described by a metric of the form

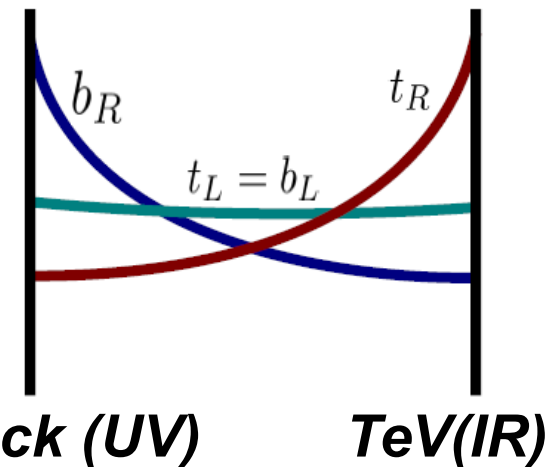
$$ds^2 = e^{2A(y)} dx_4^2 + g_{mn}(y) dy^m dy^n$$

- where $dx_4^2 = \eta_{\mu\nu} dx^\mu dx^\nu$ is the 4D Minkowski line element y^n parameterise the XD of space-time, with metric g_{mn}
- The function e^A is the warp factor, defining 4D mass scales
- gravity propagates in the compact & non-compact directions
- The 4D Newton's constant is related to the the D dimensional one by

$$G_N^{-1} = G_D^{-1} \int d^{D-4}y \sqrt{g} e^{2A}$$

- 4D Planck scale M_D exponentially suppressed by the warp factor

$$M_D = M_{PL} \exp(-\pi k R) \quad \text{where } k, R \text{ are related to } A(y)$$



QBH at colliders – features and assumptions

- high energy processes on the brane at scales of order or fundamental Planck Mass, $M_p \sim \text{TeV}$ in ADD/RS models can probe strong gravitational effects including QBH formation
- The lowest scale at which BH can be produced is M_p , which means that this BH will be essentially non-thermal, i.e. QBH
- Since QBH properties are not well understood, the best we can do is to use semi-classical limits:
 - Schwarzschild radius R_S in $(3+n)$ spacial dimensions:

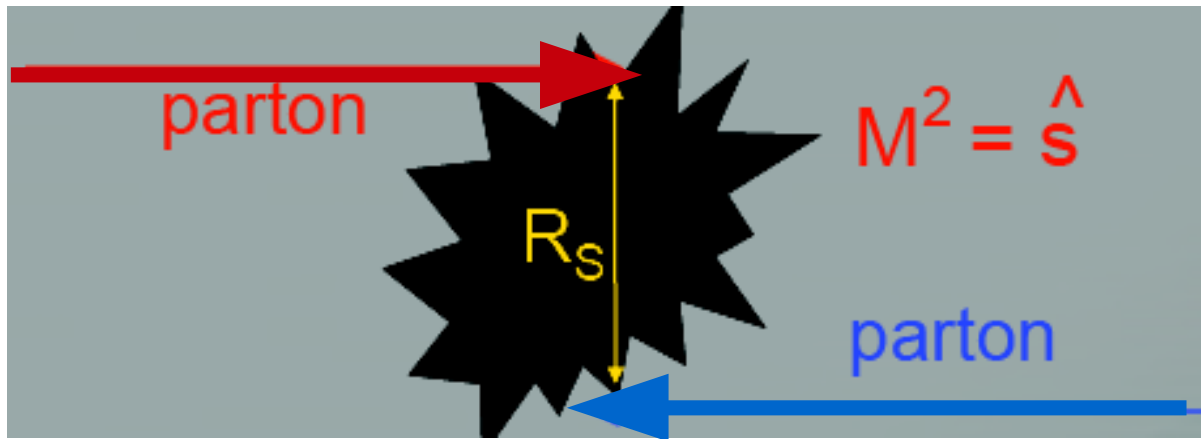
$$R_S(M_{\text{BH}}) = \frac{1}{M_D} \left[\frac{M_{\text{BH}}}{M_D} \Gamma \left(\frac{n+3}{2} \right) \frac{2^n \pi^{\frac{n-3}{2}}}{n+2} \right]^{\frac{1}{n+1}} \quad [\text{Meyers, Perry 1986}]$$

- Hawking temperature $T_H = \frac{n+1}{4\pi R_S} \sim 1 \text{ TeV}$

[Dimopolous, Landsberg 2001; Giddings, Thomas, 2002]

QBH at colliders - production

- we assume that the cross section for QBH's can be extrapolated from the classical and semi-classical cases, i.e. that it is given by the geometrical formula $\sigma = \pi r_s^2$ [Eath & Payne 1992; Eardley & Giddings 2002]
- Cross section, given by disk approximation



$$\sigma \sim \pi R_s^2 \sim 1 \text{ TeV}^{-2} \sim 10^{-38} \text{ m}^2 \sim 100 \text{ pb}$$

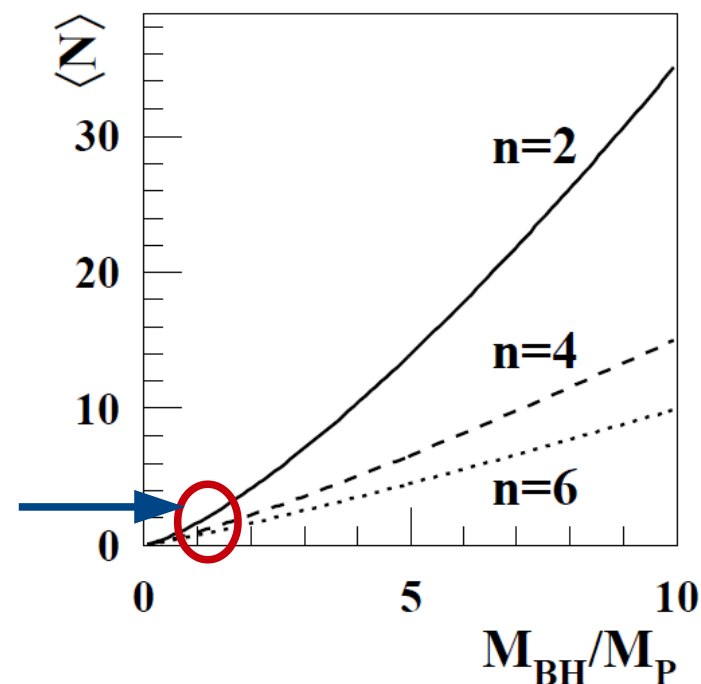
is comparable with that of the top-quark pair production!

QBH at colliders - evaporation (decay)

- Evaporation (decay) multiplicity of QBH

$$\langle N \rangle = \frac{2\sqrt{\pi}}{n+1} \left(\frac{M_{\text{BH}}}{M_P} \right)^{\frac{n+2}{n+1}} \left(\frac{8\Gamma\left(\frac{n+3}{2}\right)}{n+2} \right)^{\frac{1}{n+1}}$$

- In 4D $\langle N \rangle \approx \frac{M_{\text{BH}}}{2T_H}$
- So indeed, QBH will rather explode into just a few particles resembling strong gravitational re-scattering.
- It is thus tempting to treat QBH as particle, carrying the quantum numbers of the particles which created them.



[Dimopoulos & Landsberg 2001]

QBH at colliders - the approach (1)

- Calmet, Fragkakis and Gausmann (2011) proposed a field theoretical model to describe the interactions of QBHs
- It assumes that QBHs can be treated as quantum fields, classified according to representations of the Lorentz group.
- Furthermore, they are classified according to their transformations under the gauge groups of the SM. This fixes their interactions with matter.
- In a proton-proton collider, quantum black holes would be produced from the collisions of quarks and gluons

QBH at colliders - the approach (2)

- We are thus particularly interested in the QBHs carrying QCD and QED quantum numbers and with spins 0, 1/2 and 1 since these should be the lowest lying states.
- Generically speaking, quantum black holes form representations of $SU(3)_c$ and carry a QED charge.
- We consider process of two partons p_i and p_j forming a quantum black hole in the c representation of $SU(3)_c$ and charge q as: $p_i p_j \rightarrow QBH_c^q$

QBH at colliders - details of the calculations(1)

$$\sigma(pp \rightarrow QBH) = \sum_{i,j} \int dx_1 dx_2 f_i(x_1) f_j(x_2) \sigma(ij \rightarrow QBH)$$

- Continuous QBH mass = $\sqrt{s}$ invariant mass of two partons
- In 4D case (just for the sake of simplicity) [running of the gravitational coupling constant can be radically affected by a hidden sector with a large number of particles - analogous to KK modes - Calmet, Hsu, reeb, 2008]

we have

$$r_s(s, \overline{M}_{PL}) = \frac{\sqrt{s}}{4\pi \overline{M}_{PL}^2} \quad \text{and} \quad \sigma = \frac{1}{16\pi} \frac{s}{\overline{M}_{PL}^4} \Theta(\sqrt{s} - \overline{M}_{PL})$$

QBH at colliders - details of the calculations(2)

- Let us consider production of spinless QBH in the collision of two fermions (quarks for example)
- Recalling that $\sigma \sim s$, one can use **4-fermion contact interaction** to describe this cross section behaviour as a function of s :

$$\mathcal{L}_{cont} = \frac{g_c}{\Lambda^2} \bar{\psi}_i \psi_i \bar{\psi}_j \psi_j$$

- which gives $|M|^2 = \frac{g_c^2}{\Lambda^4}$ for $qq \rightarrow qq$ scattering and the respective total cross section (for massless fermions)

- $\sigma = \frac{1}{16\pi s} |M|^2 = \frac{g_c^2}{16\pi \Lambda^4} s$ which reproduces the QBH cross section

for $g_c = 1$ and $\Lambda = \overline{M}_{PL}$!

Implementation into CalcHEP (1)

- **CalcHEP – Calculator for High Energy Physics**
[AB,Pukhov,Christensen '04, created by Pukhov et al in '89 as a CompHEP]
- **Uses the model written in the form of Feynman Rules for any theory (including effective ones) and**
 - Evaluates Matrix element square Analytically
 - Evaluates total and differential CS numerically
 - Simulate events which can be used in further analysis and in further (detector simulations) by CMS, ATLAS and other experimental collaboration
 - One of the leading packages for the LHC physics
 - Very convenient/user-friendly GUI and online help

Implementation into CalcHEP (2)

- QBH model was implemented as contact interactions

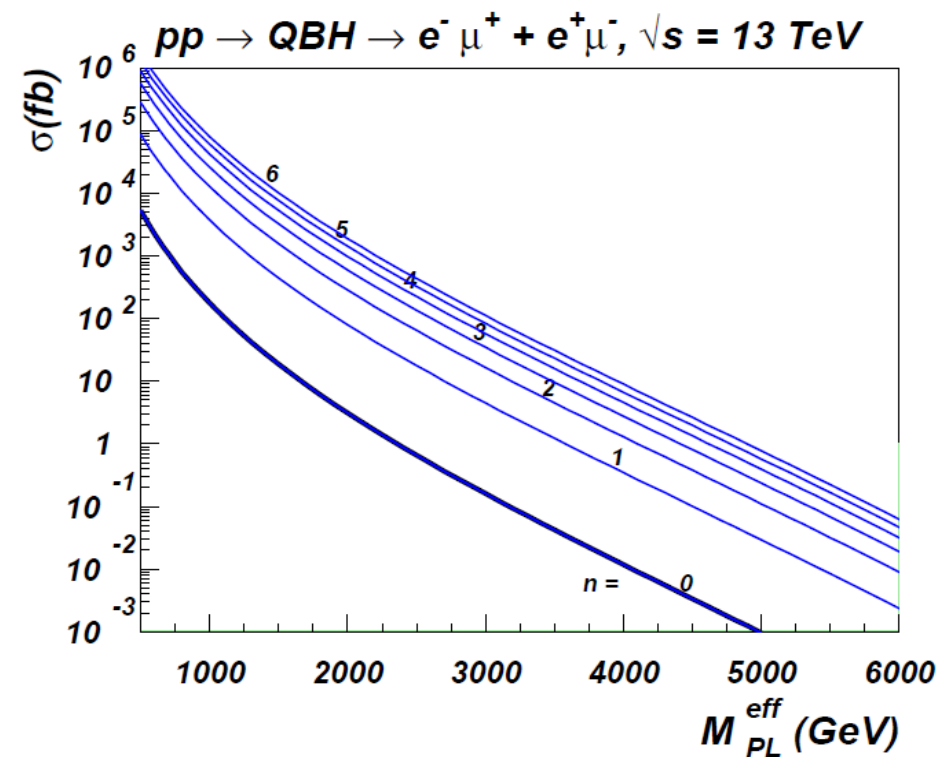
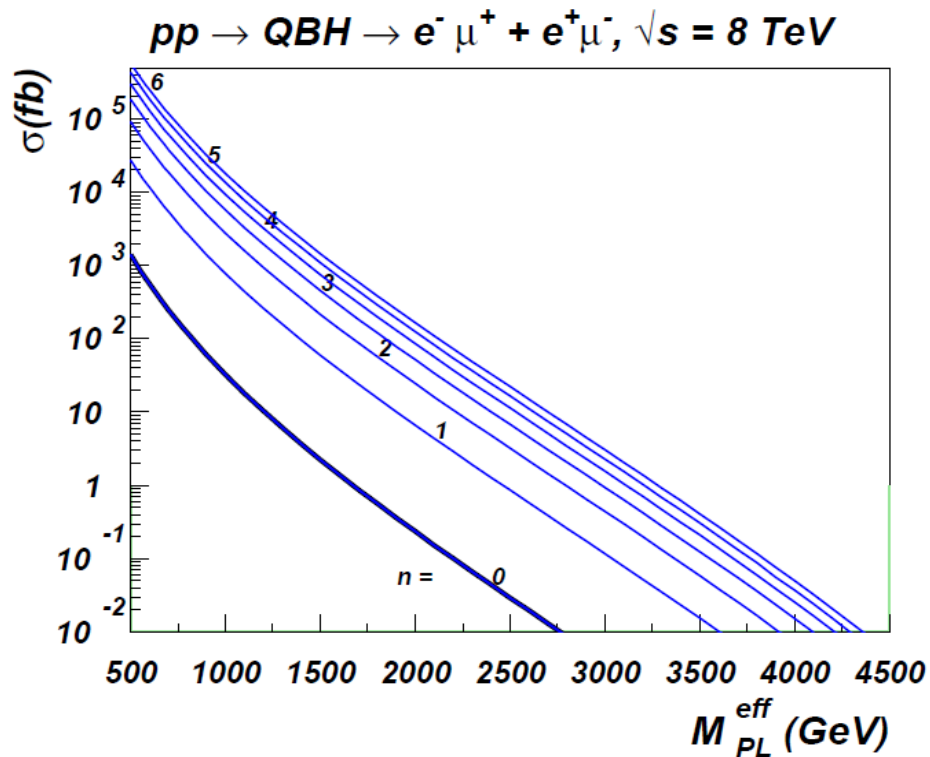
$$\mathcal{L}_{cont}^X = g_c \left(\sum_{leptons} \bar{\psi}_i^\ell \psi_j^\ell X + \sum_{quarks} \bar{\psi}_i^q \psi_j^q X \right)$$

to take into account all quark and lepton combinations, where i, j are lepton and quark flavour indices, propagator of **X "contact"** field is $\frac{i}{M_{PL}^2}$

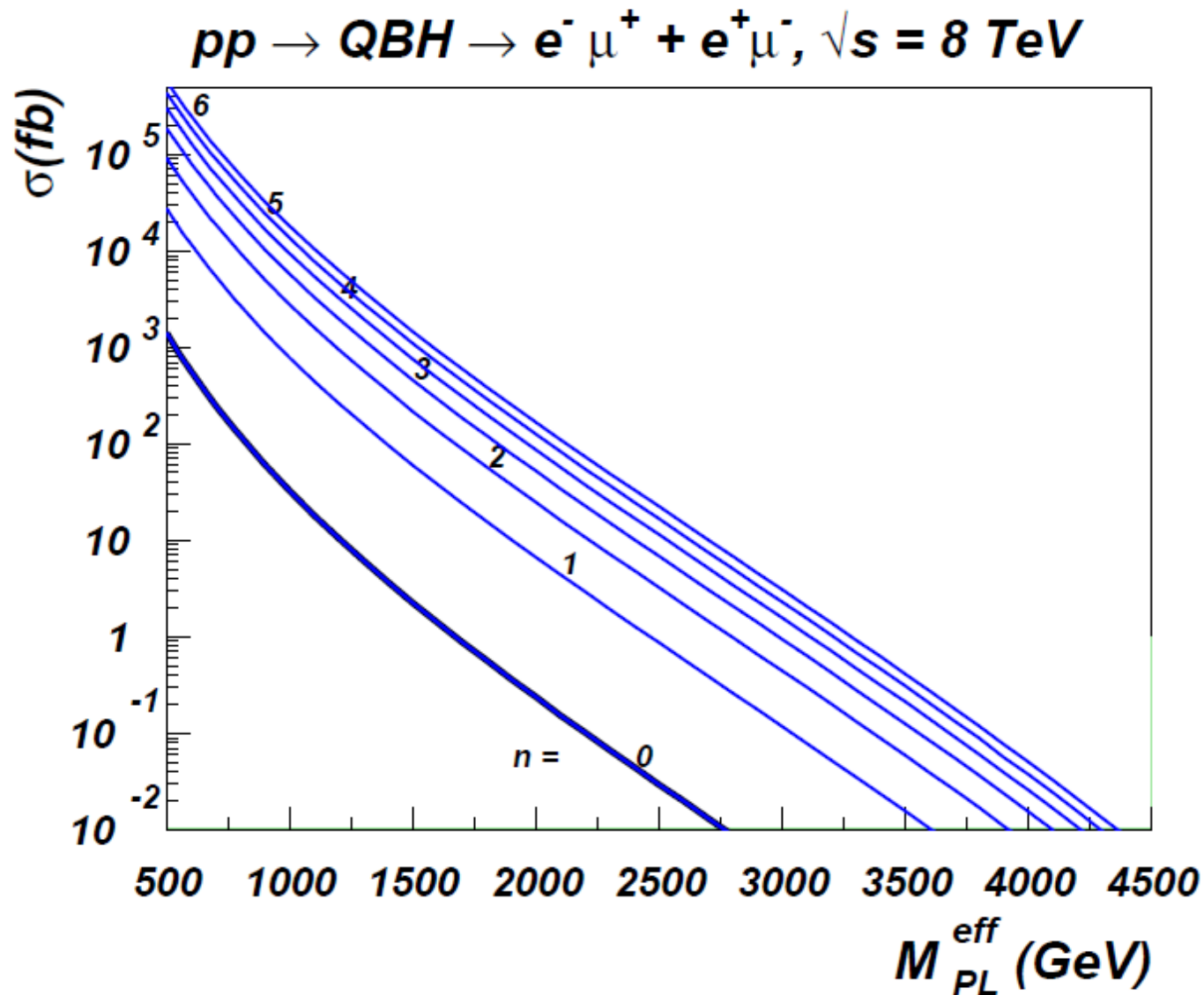
- $g_c = (n_l + 3n_q)^{-1/4}$ is the normalisation factor accounting number of lepton, $n_l = 9$ and quark, $n_q = 18$ combinations, including the quark colour factor
- This Lagrangian exactly reproduces the cross section of the spinless neutral QBH production and decay in the $2 \rightarrow 2$ fermion process as a function of s

QBH Production rate at the LHC for di-lepton signature

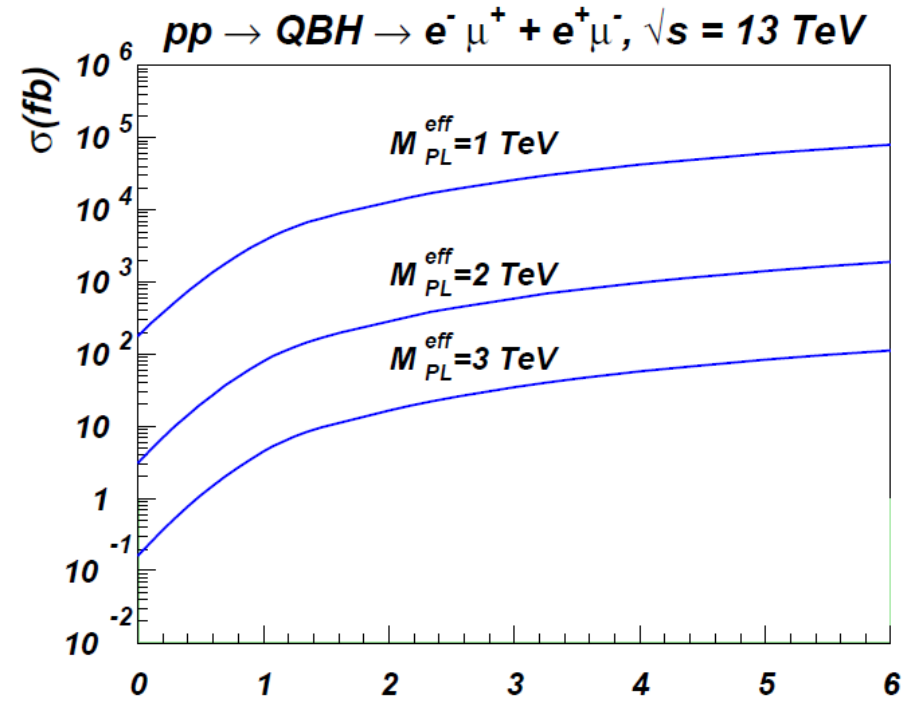
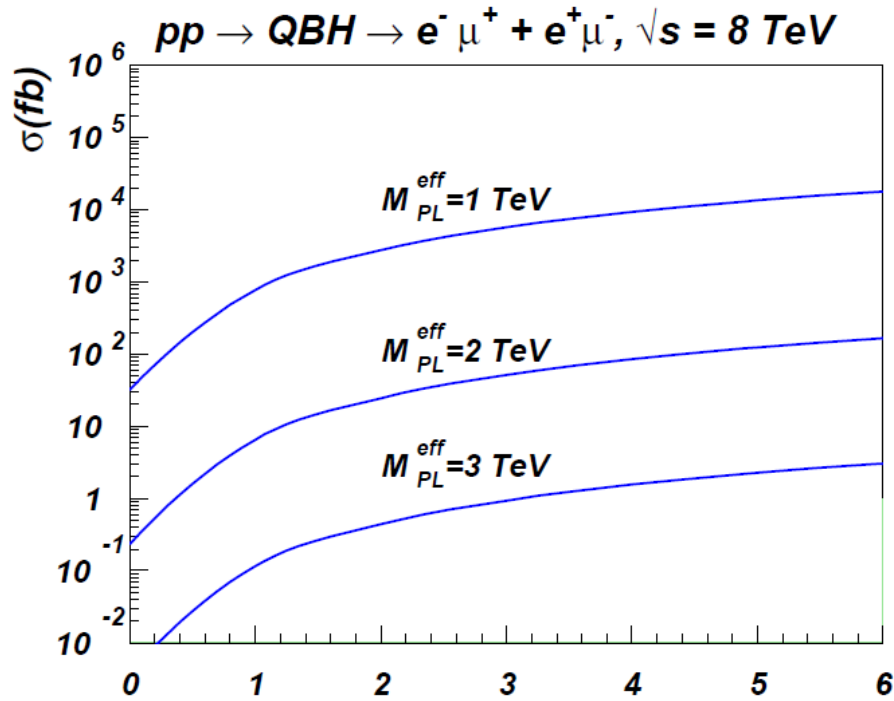
- Why di-lepton signature?
 - very “clean” (low background) , especially for $e\mu$ signature
 - Best energy resolution
 - High ($\sim 90\%$) lepton ID efficiency



QBH Production rate at the LHC for di-lepton signature



QBH Production rate at the LHC



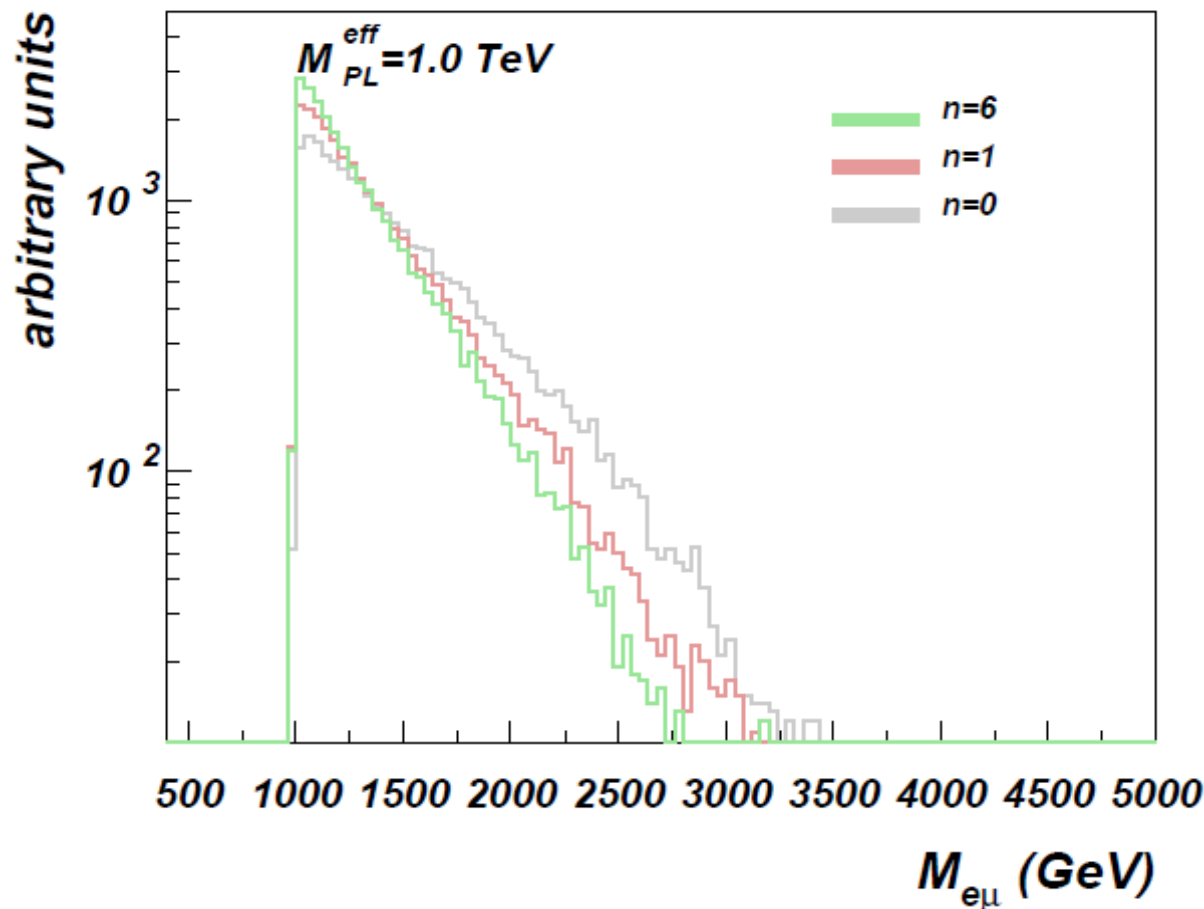
cross section of $pp \rightarrow QBH \rightarrow e\mu$ process in fb

	$\bar{M}_{PL} \text{ (TeV)}$	n=0	n=1	n=2	n=3	n=4	n=5	n=6
LHC@8TeV	1.0	32.3	782.	2760.	5730.	9370.	13500.	18100.
	2.0	0.235	6.60	24.5	51.8	85.7	124.	166.
	3.0	0.00388	0.116	0.439	0.939	1.56	2.28	3.06
LHC@13TeV	1.0	177.	373.	12800.	26000.	42200.	60500.	80400.
	2.0	3.11	79.7	286.	596.	980.	1420.	1890.
	3.0	0.161	4.48	16.5	34.8	57.4	83.5	112.

Di-lepton invariant mass distribution

- It does not look like a resonance(!), the fall of parton densities is (much) more steep with the energy than the increase of the parton-level cross section(!)

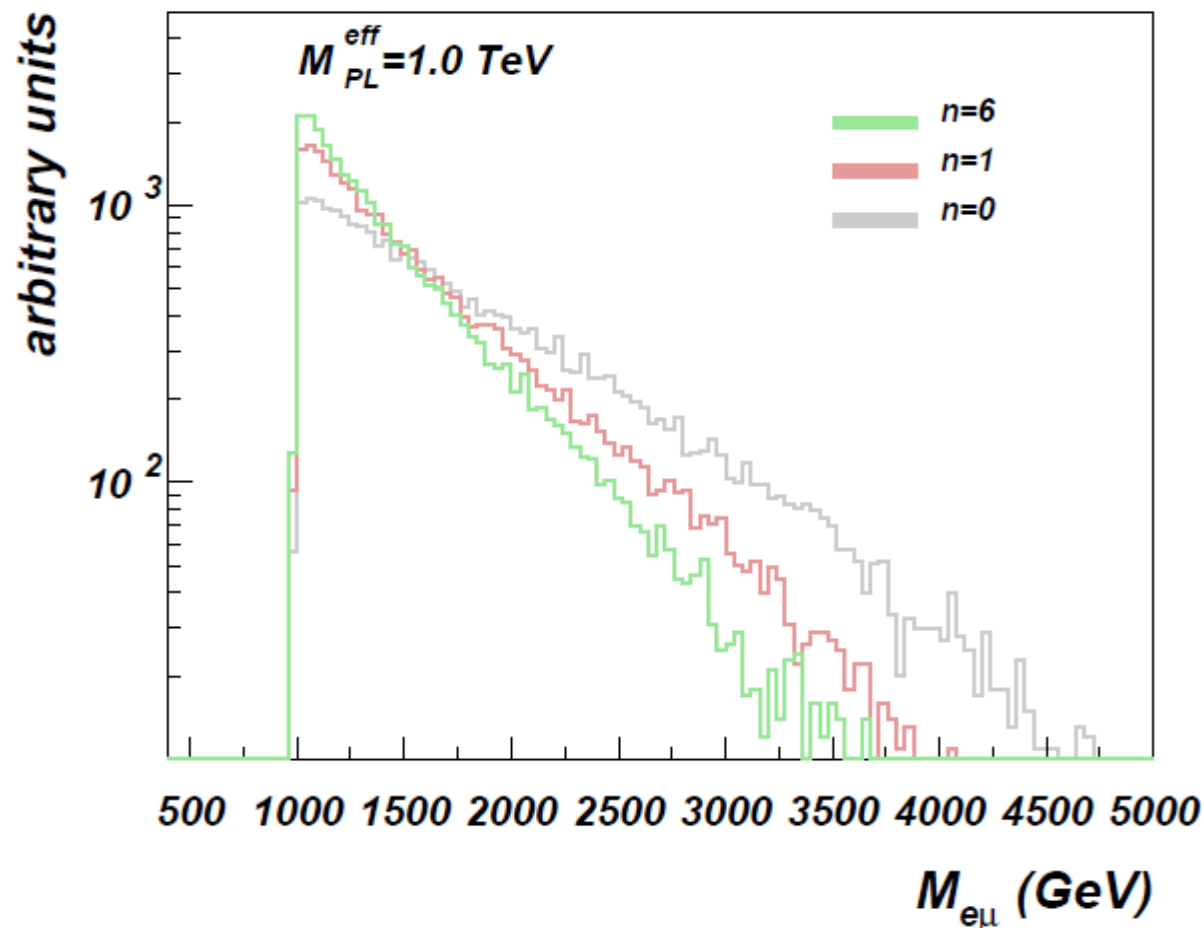
$pp \rightarrow QBH \rightarrow e^+ \mu^-$ @LHC, $\sqrt{s} = 8$ TeV



Di-lepton invariant mass distribution

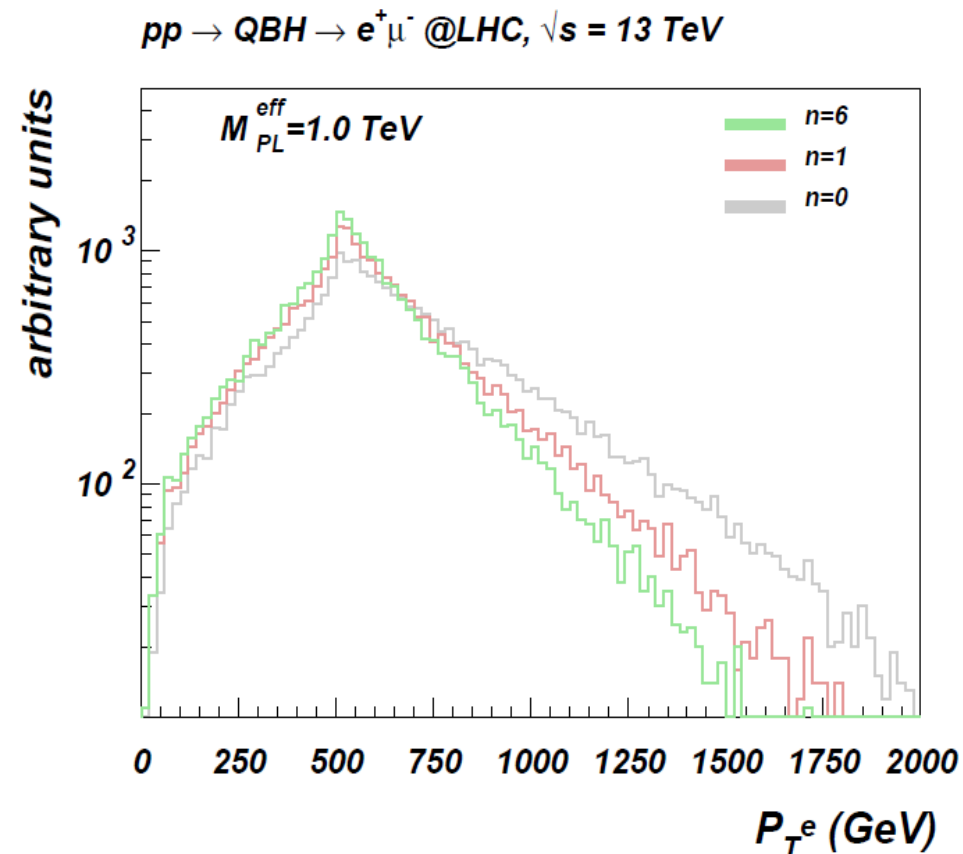
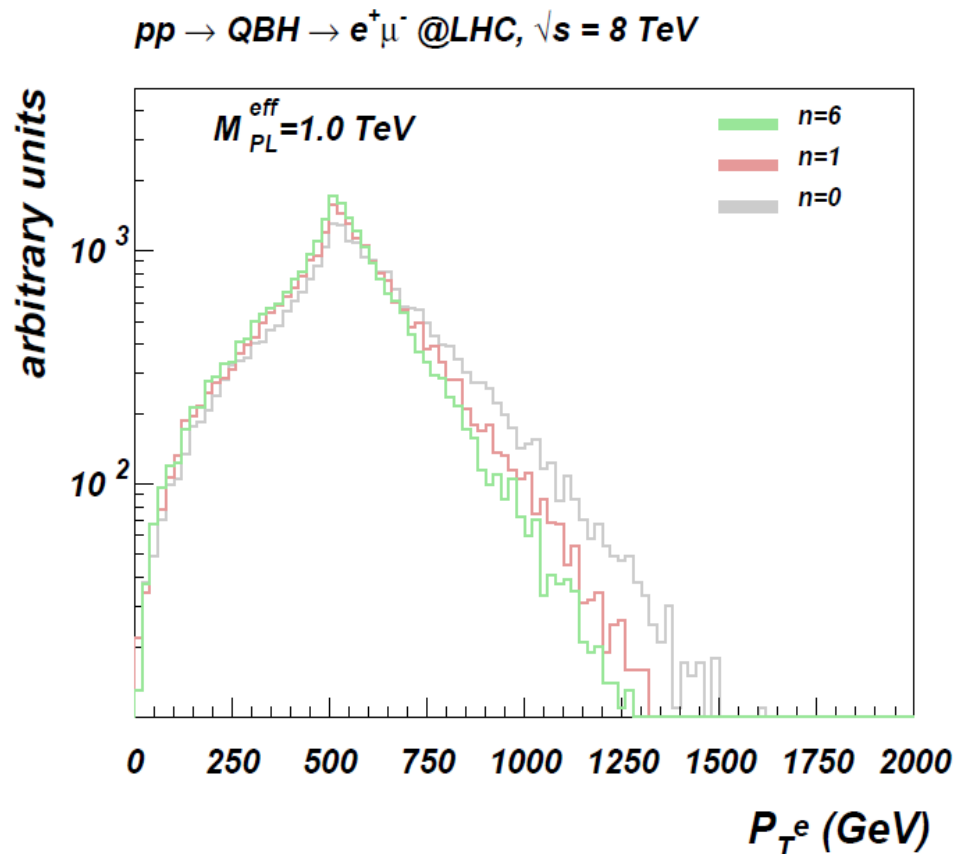
- As far as more phase space is available at 13 TeV LHC, one can see the difference in the tail for different # of XD reflecting non-trivial R_s energy scaling

$pp \rightarrow QBH \rightarrow e^+ \mu^-$ @LHC, $\sqrt{s} = 13$ TeV



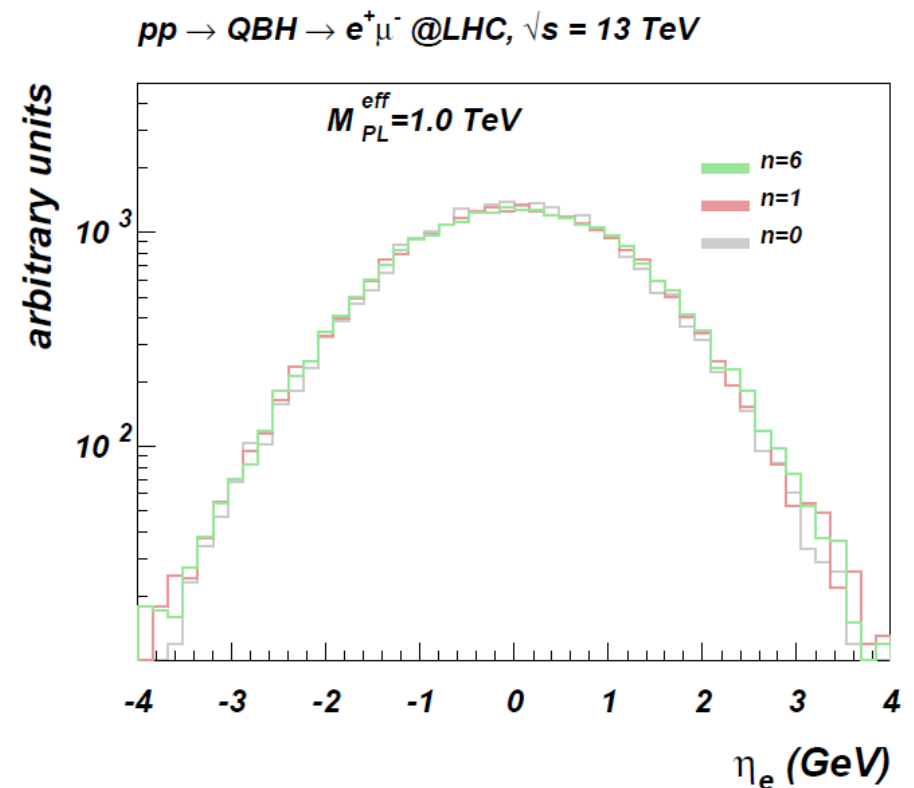
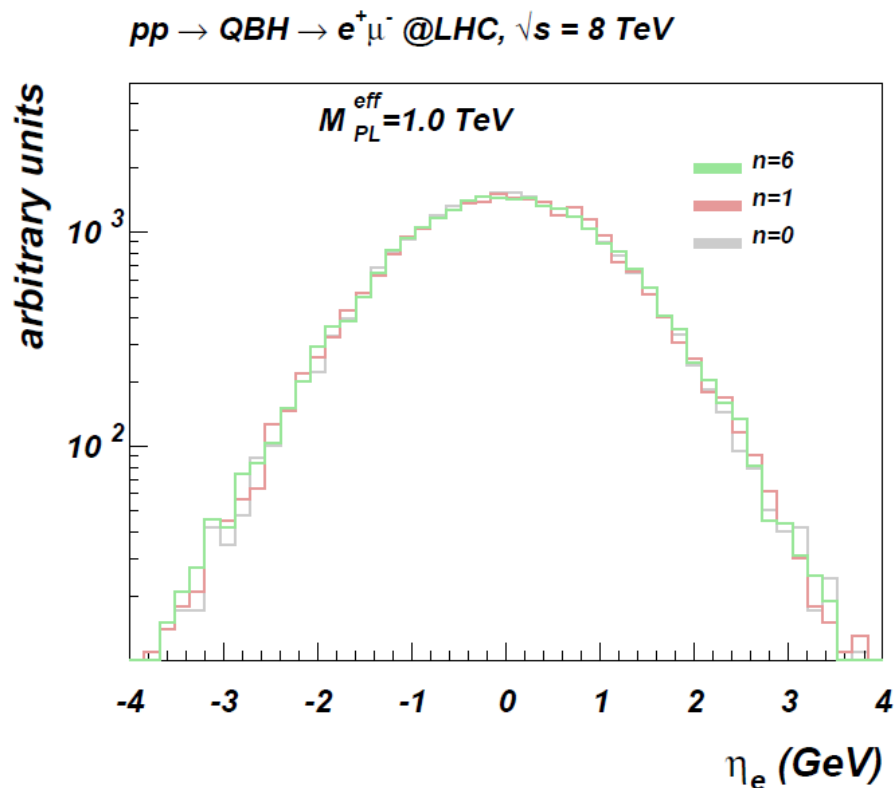
Lepton transverse momentum

- High enough to fall into the acceptance region, $P_T > 20$ GeV

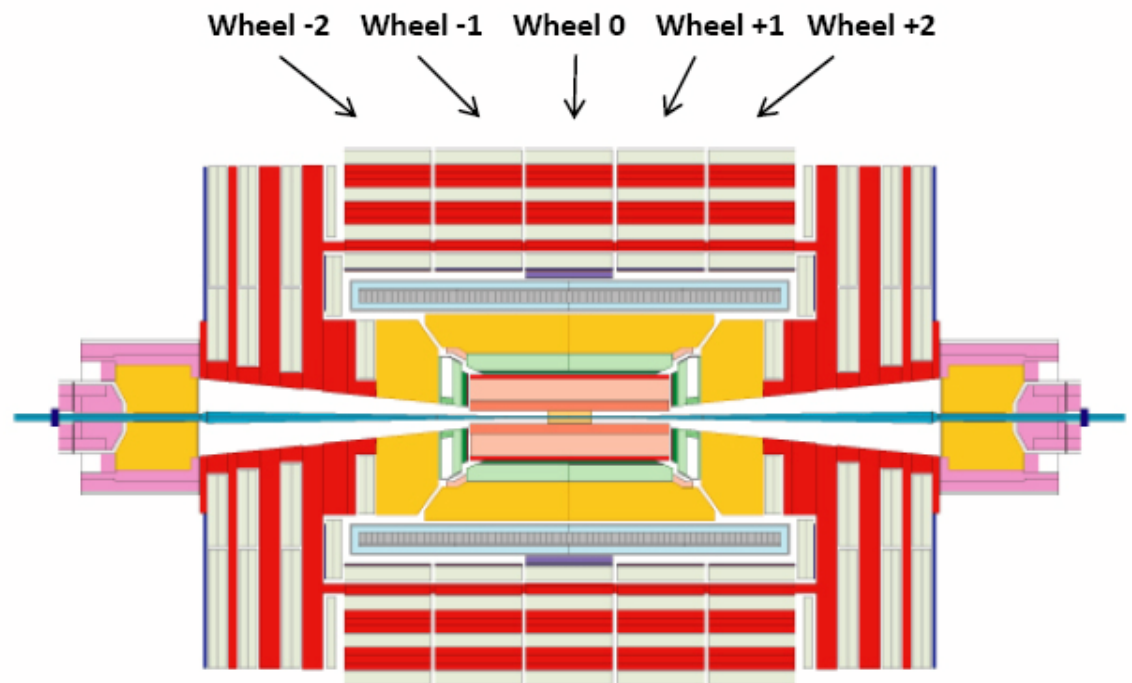
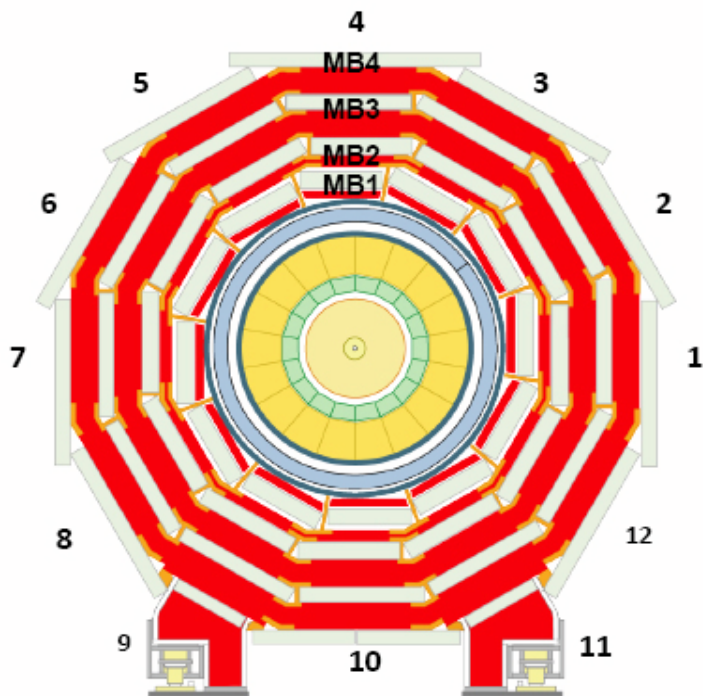


Lepton pseudo-rapidity

- central enough to fall mostly into the acceptance region, $|\eta_\tau| < 2.5$



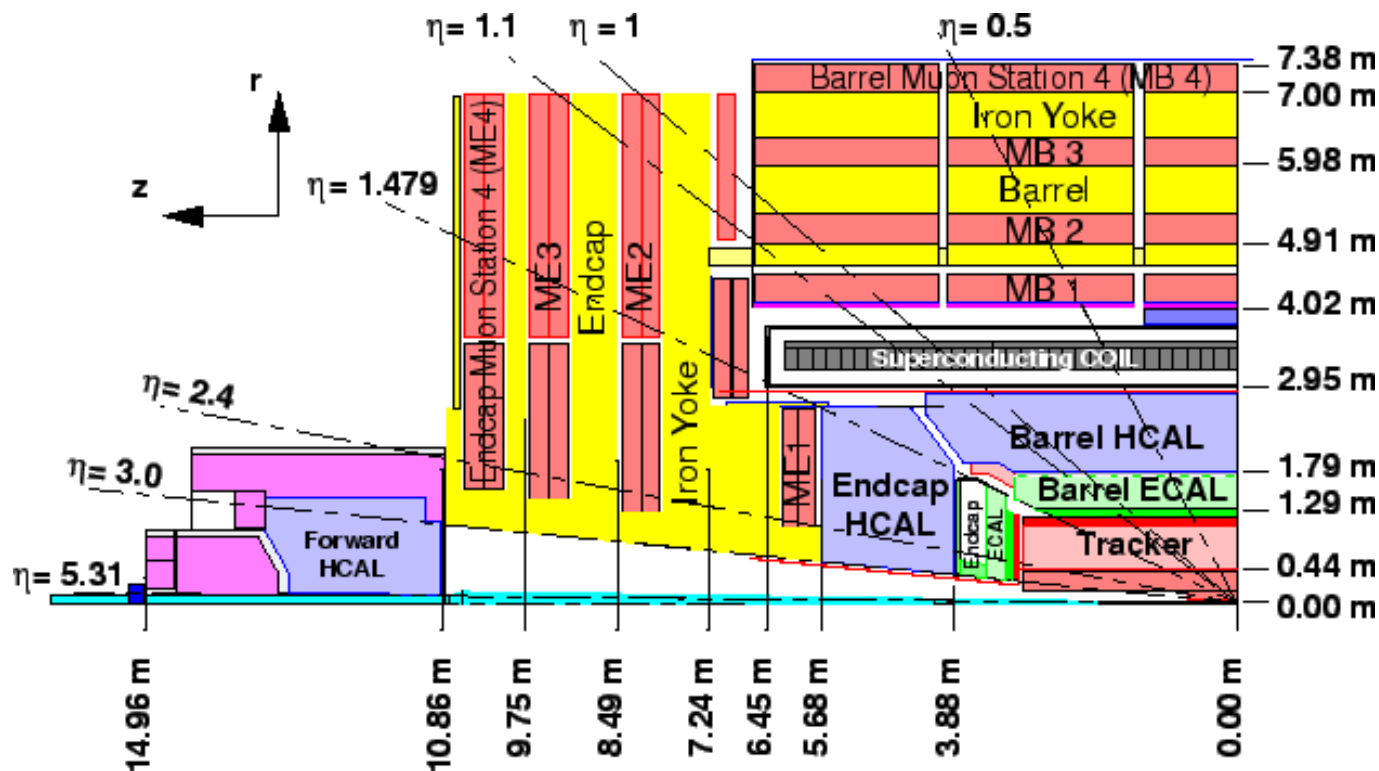
CMS detector



Pseudo-rapidity

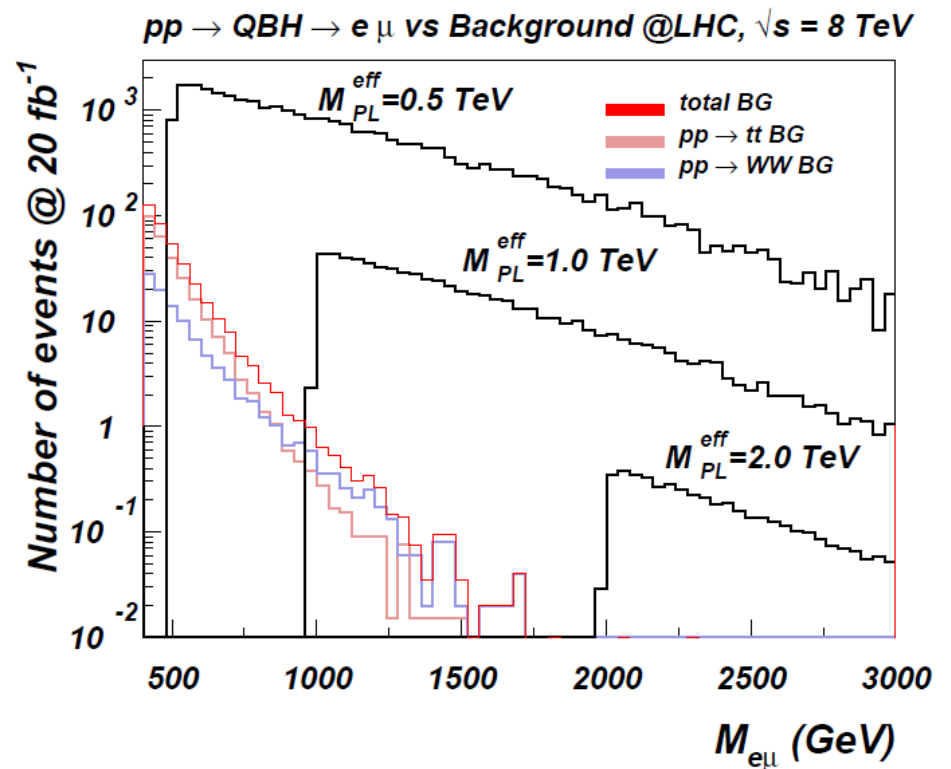
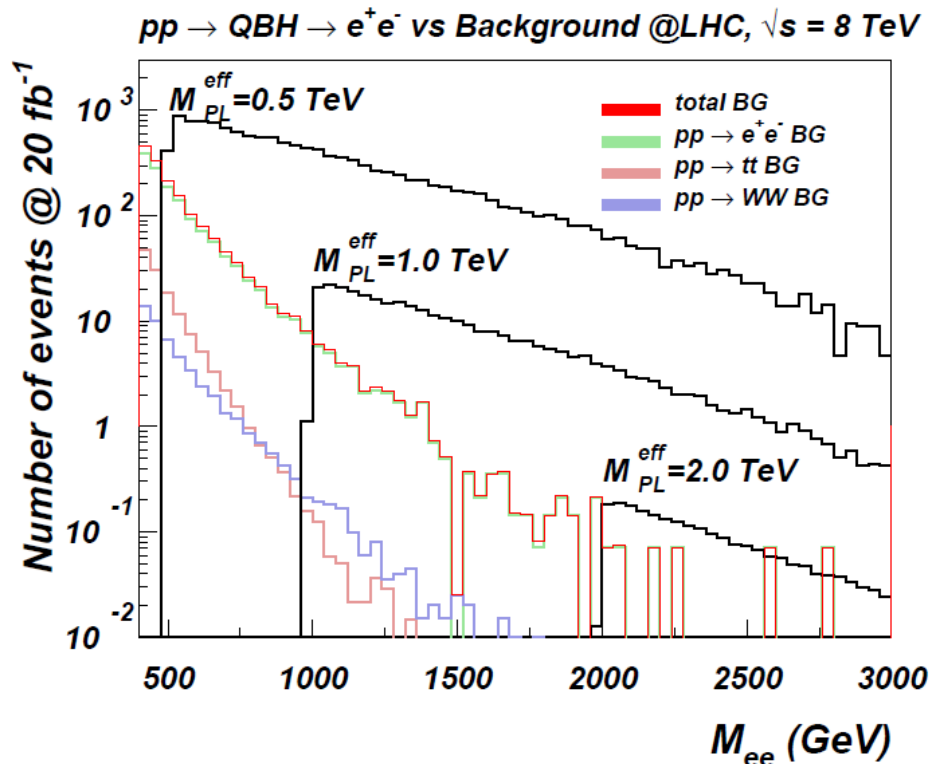
CMS detector - Longitudinal view

$$\eta = \frac{1}{2} \ln \left(\frac{|\mathbf{p}| + p_L}{|\mathbf{p}| - p_L} \right) = \operatorname{artanh} \left(\frac{p_L}{|\mathbf{p}|} \right)$$



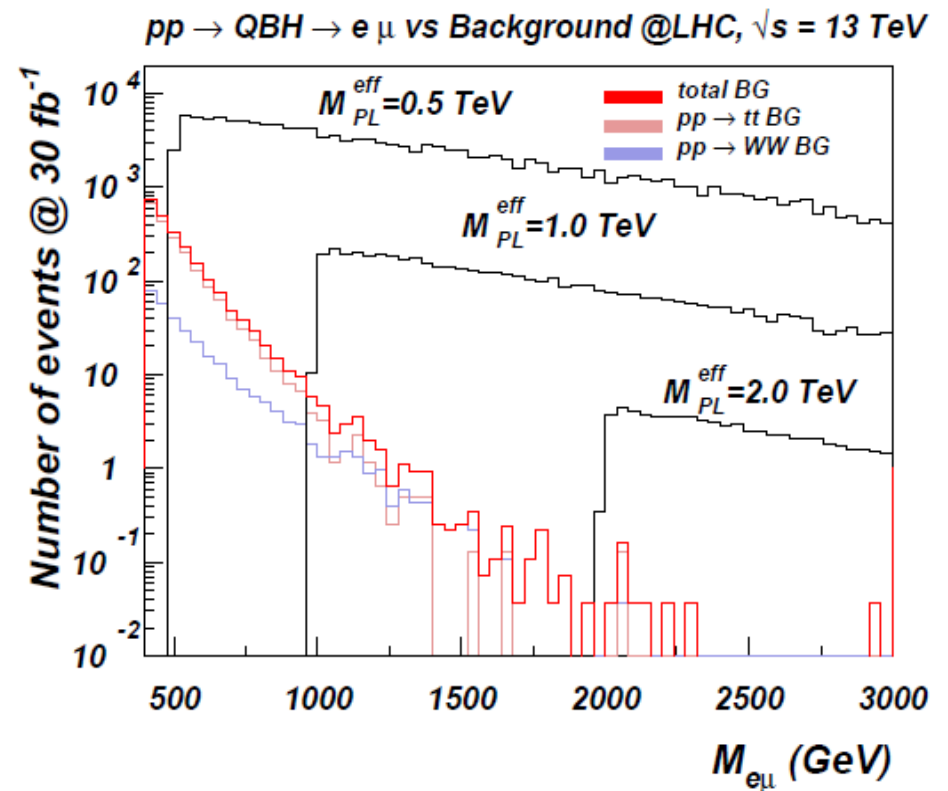
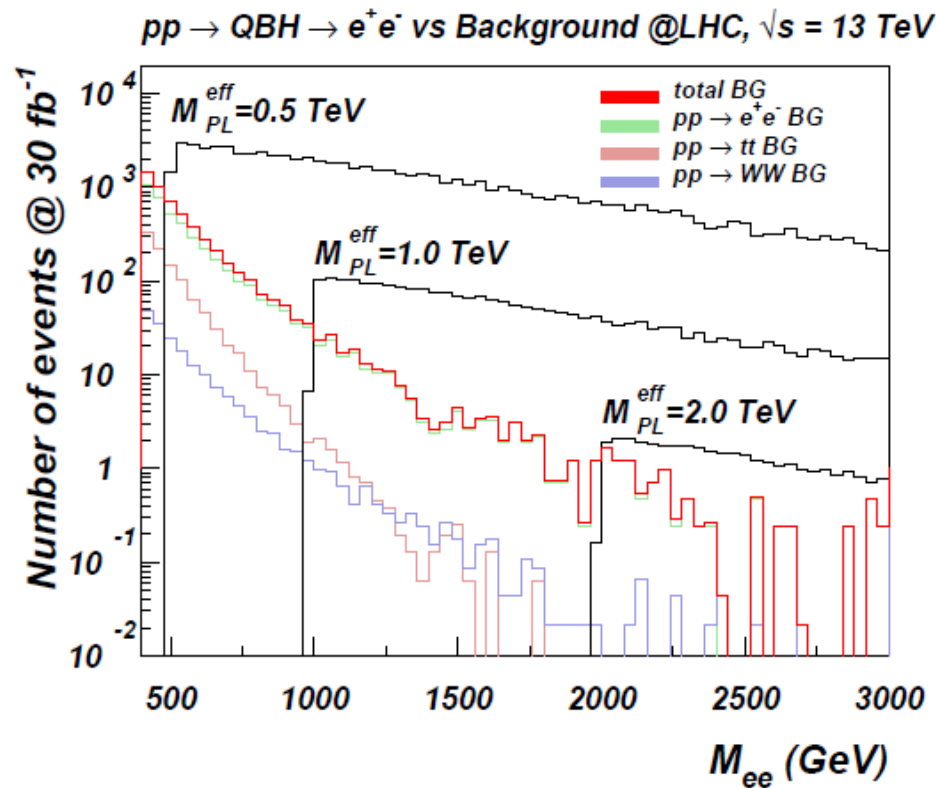
Signal vs Background @ 8 TeV

- For ee or $\mu\mu$ signatures - Drell-Yan process is dominant BG
- For $e\mu$ signature - the leading BG comes from tt production, but it is negligible!



Signal vs Background @ 13 TeV

- For ee or $\mu\mu$ signatures - Drell-Yan process is dominant BG
- For $e\mu$ signature - the leading BG comes from $t\bar{t}$ production, but it is negligible!

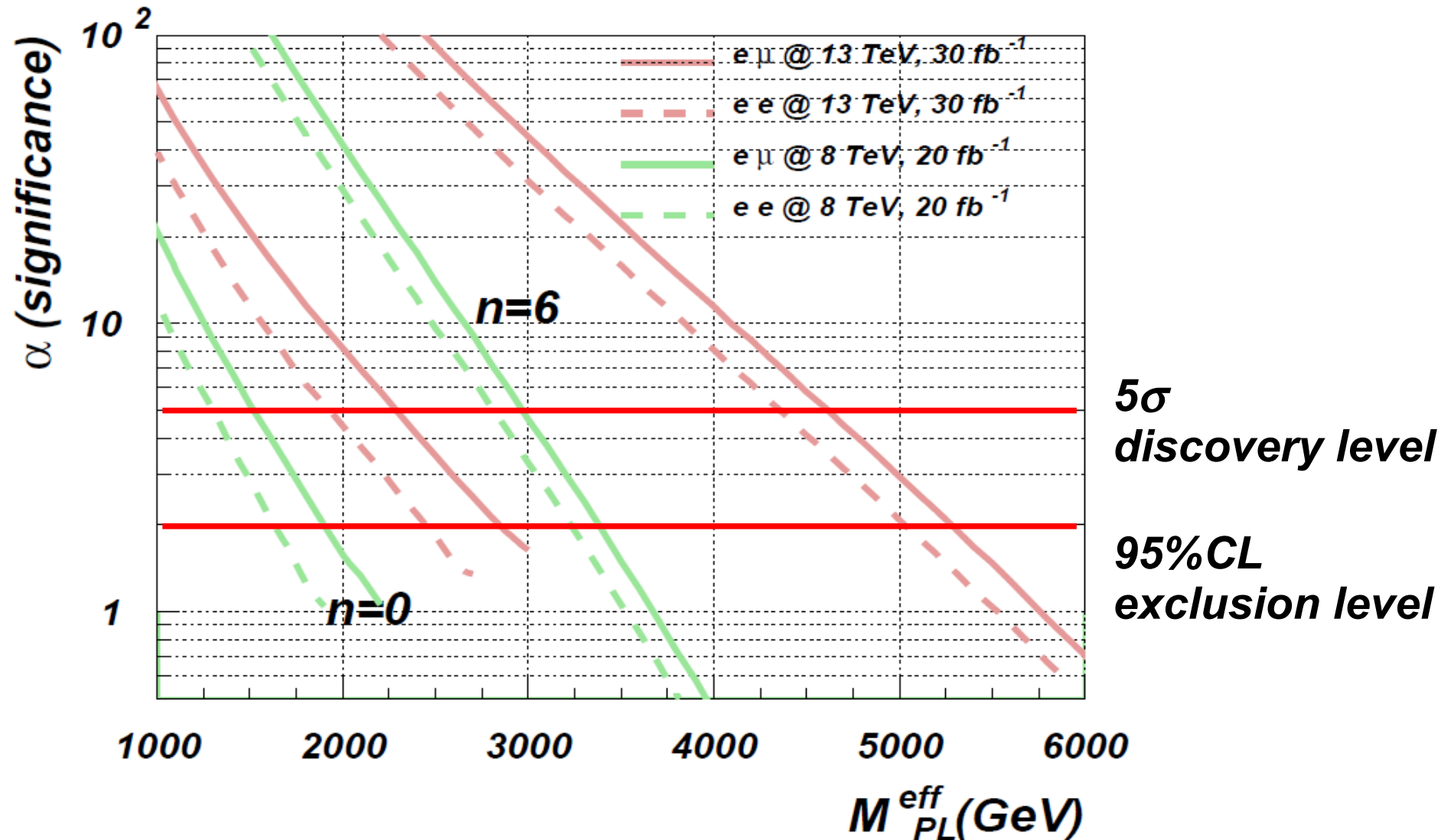


Simulation and selection details

- parton-level analysis, but we take into account realistic electromagnetic energy resolution, using a value of $0.15/\sqrt{E(\text{GeV})}$ typical for the ATLAS and CMS detectors
- We also suggest the simple analysis cut $M_{e\mu}(M_{ee}) > 1.1 \times \overline{M}_{PL}$ noting that the acceptance efficiency will be very similar for different n models
- For both criteria, exclusion and discovery, we use the following formula for statistical signal significance

$$\alpha = 2(\sqrt{N_S + N_B} - \sqrt{N_B})$$

LHC sensitivity@8TeV & 13 TeV



LHC sensitivity@8TeV & 13 TeV

- With the signature suggested M_p in XD models can be tested up to ~ 6 TeV via QBH production!

		LHC@8TeV		LHC@13TeV	
CL	n	e^+e^-	$e\mu$	e^+e^-	$e\mu$
95%CL	0	1920 GeV	2240 GeV	3360 GeV	3780 GeV
5σ	0	1550 GeV	1810 GeV	2790 GeV	3180 GeV
95%CL	6	3540 GeV	3680 GeV	5510 GeV	5750 GeV
5σ	6	3140 GeV	3300 GeV	4850 GeV	5100 GeV

Pheno vs Experimental results

- Results are consistent with very recent CMS limits, which confirms reliability of the simulations and projections for the LHC@13 TeV

OUR RESULTS

		LHC@8TeV	
CL	n	e^+e^-	$e\mu$
95%CL	0	1920 GeV	2240 GeV
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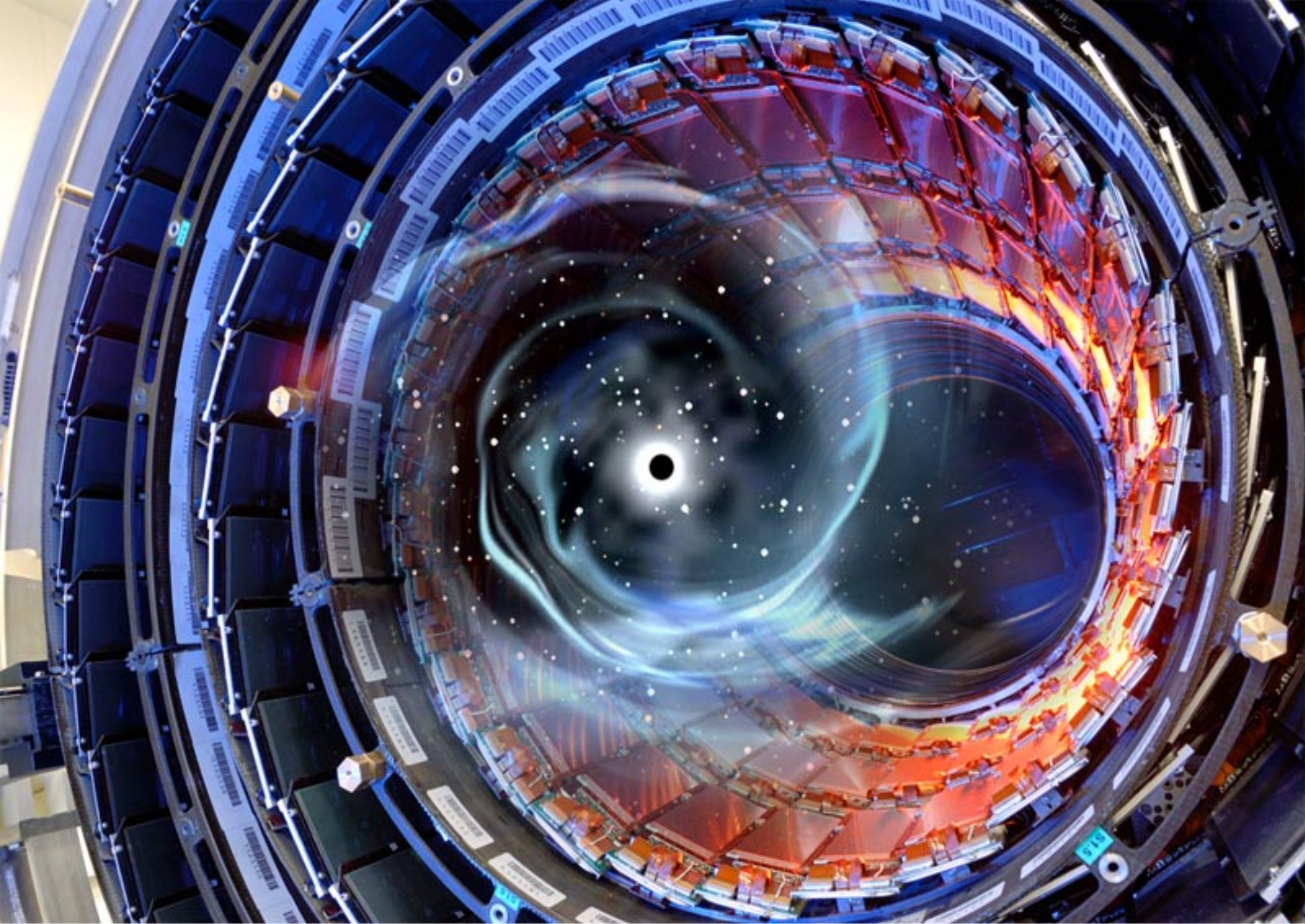


CMS PAS EXO-13-002

Theoretical Model	Mass Limit (TeV)
QBH n=0	1.99
QBH n=1 (RS)	2.36
QBH n=1 (PDG)	2.81
QBH n=2	3.15
QBH n=3	3.34
QBH n=4	3.46
QBH n=5	3.55
QBH n=6	3.63

Conclusions

- QBH – could provide a crucial link between GR and QFT if we could explore them at the LHC
- We have proposed an effective Lagrangian to describe QBH production at the LHC
- The model is implemented into CalcHEP and is publicly available at HEPMDB – High Energy Physics Model Database
<http://hepmdb.soton.ac.uk/hepmdb:1113.0146>
- Di-lepton signature has been suggested and explored
- Model/generator as been already used by CMS for LHC@8TeV data
- LHC@13 has an impressive potential to test M_p in XD models upto ~ 6 TeV





THANK YOU!